

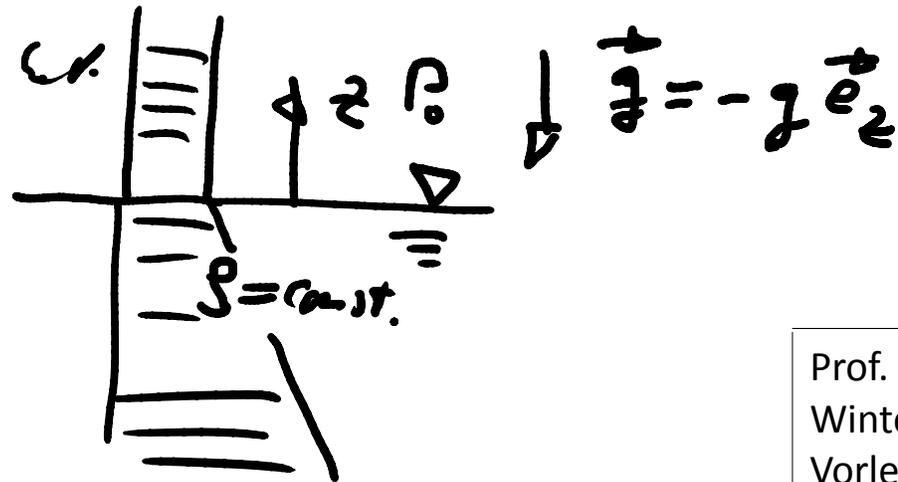
Weitere Spezialisierung

Hydrostatik $\vec{p} \equiv \sigma$, da keine Relativbewegung
zwischen Flüssigkeitsteilchen.

$$\rho \frac{D\vec{u}}{Dt} = \sigma$$

$\vec{p} = \rho \vec{g}$ Hydrostatische Grundgleichung.

z.B. Druckbestimmung, schwere Flüssigkeit



$$\frac{dp}{dz} \vec{e}_z = -\rho \vec{g}$$

$$p(z) = -\rho g z + p_0$$

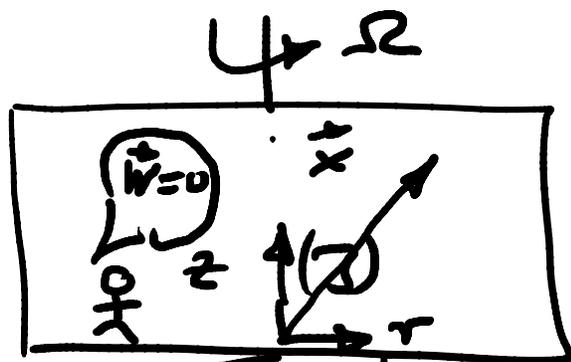


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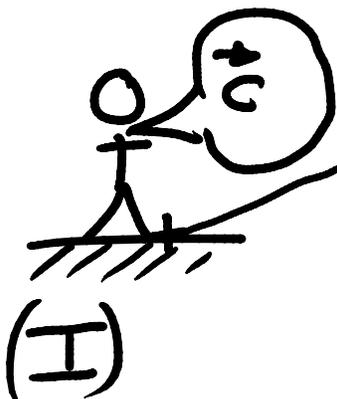
Technische
Fluidsysteme

Prof. Dr.-Ing. Peter Pelz
Wintersemester 2012/13
Vorlesung 5 F 57



$$\vec{c} = \vec{w} + \vec{u} + \vec{v}$$

$$\vec{v} = \left[\frac{D\vec{r}}{Dt} \right]_{\mathbf{I}} ; \vec{w} = \left[\frac{D\vec{x}}{Dt} \right]_{\mathbf{B}}$$



mit bewegtes System
↳ keine Relativgeschw.

$$\vec{p} \equiv 0$$

$$\rho \frac{D\vec{w}}{Dt} \equiv 0$$

\vec{w} Relativgeschw.

$$\vec{u} \text{ Umfahrgeschw. } \vec{u} = \vec{\Omega} \times \vec{r}$$

\vec{v} Führunggeschw.

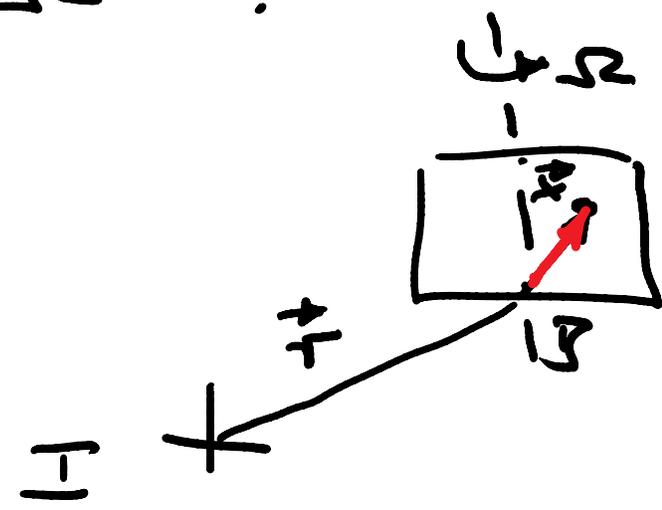
\vec{c} Absolutgeschw.



$$\vec{c} = \left[\frac{D \vec{x}}{Dt} \right]_{\mathcal{B}} + \left[\frac{D \vec{r}}{Dt} \right]_{\mathcal{I}} + \vec{\Omega} \times \vec{r}$$

\mathcal{B} : Beschränktes System

\mathcal{I} : Inertialsystem



$$\vec{c} = \left[\frac{D}{Dt} (\vec{x} + \vec{r}) \right]_{\mathcal{I}} = \left[\frac{D \vec{x}}{Dt} \right]_{\mathcal{I}} + \left[\frac{D \vec{r}}{Dt} \right]_{\mathcal{I}}$$

$$\left[\frac{D \vec{x}}{Dt} \right]_{\mathcal{I}} = \left[\frac{D \vec{x}}{Dt} \right]_{\mathcal{B}} + \vec{\Omega} \times \vec{x}$$

gilt allgemein für eine beliebige Vektor \vec{b} z.B.

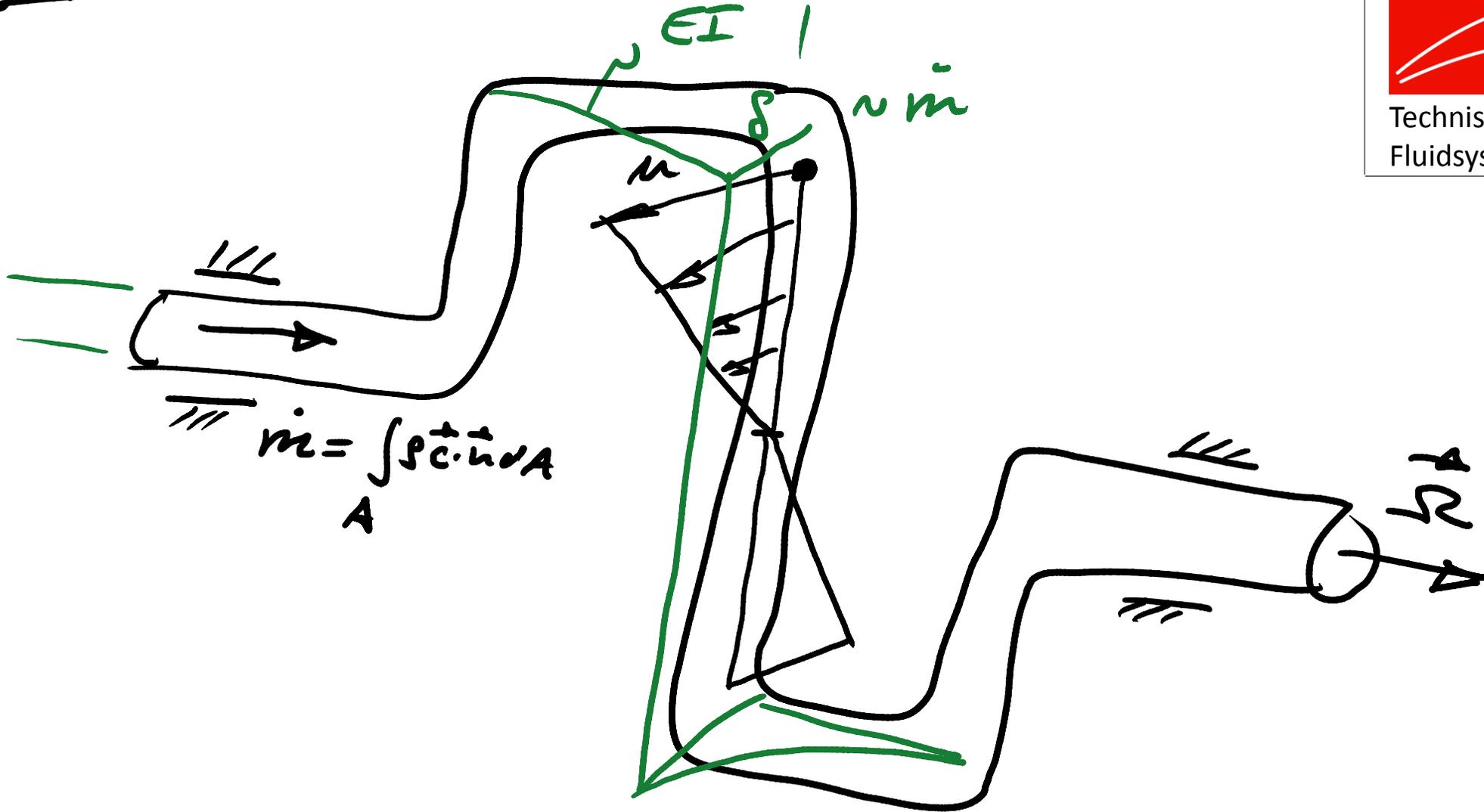
$$\vec{I} = \int \rho \vec{c} dV$$

$$\vec{\Omega} = \int \rho \vec{x} \times \vec{c} dV$$

$$\left[\frac{D}{Dt} \right]_{\vec{I}} = \left[\frac{D}{Dt} \right]_{\vec{I}} + \vec{\Omega} \times \vec{x} \dots$$



Beispiel Controlvolumen

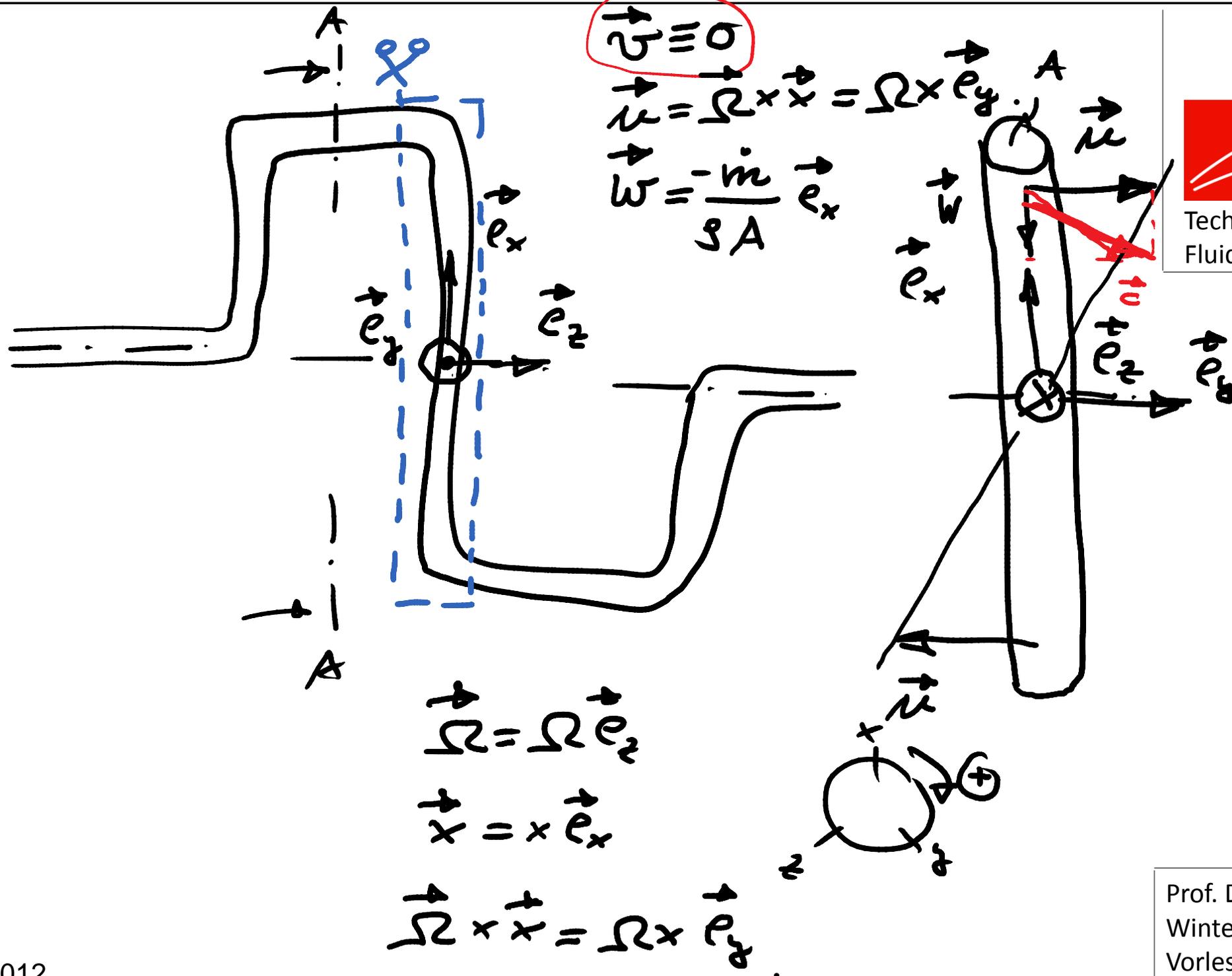


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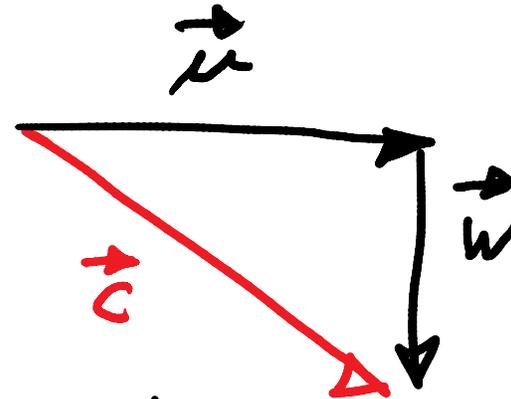
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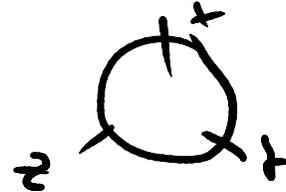


$$\left[\frac{D\vec{I}}{Dt} \right]_I = \vec{F}$$

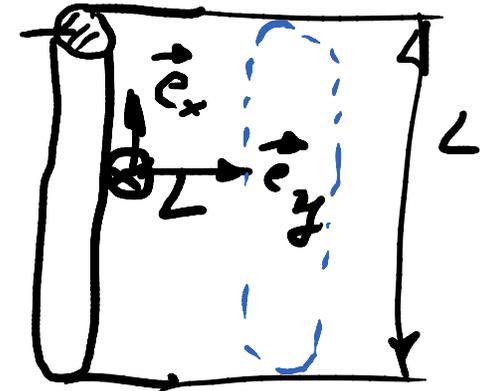


Geschwindigkeitssdreieck.

$$\vec{I} = \int_V \rho \vec{c} dV = \int_V \rho (\vec{u} + \vec{w}) dV$$



$$\left[\frac{D\vec{I}}{Dt} \right]_I = \underbrace{\left[\frac{D\vec{I}}{Dt} \right]_I}_{\equiv 0} + \underbrace{\vec{\Omega} \times \vec{I}}_{\text{Coriolis}} = \vec{F} = \Omega m L \vec{e}_2 \times \vec{e}_x = -\Omega m L \vec{e}_y \quad m = \rho V = \rho w s A$$



$$= \underbrace{\vec{\Omega} \times \int_V \rho (\vec{\Omega} \times \vec{x}) dV}_{\equiv 0} + \underbrace{\vec{\Omega} \times \int_V \rho \vec{w} dV}_{\text{Coriolis}} = \vec{\Omega} \times \rho \vec{w} A L$$



$$\rho \left[\frac{D\vec{c}}{Dt} \right]_I = \nabla \cdot \vec{T} + \vec{f}$$

$$\rho \left[\frac{D\vec{c}}{Dt} \right]_B + \rho \vec{\Omega} \times \vec{c} = \nabla \cdot \vec{T} + \vec{f}$$

$$\vec{c} = \vec{w} + \vec{u} + \vec{\Omega} \times \vec{x}$$

Scheinkräfte

$$\rho \left[\frac{D\vec{w}}{Dt} \right]_B = \nabla \cdot \vec{T} + \vec{f} - \rho \left(\vec{a} + 2 \vec{\Omega} \times \vec{w} + \left(\vec{\Omega} \times (\vec{\Omega} \times \vec{x}) \right) + \rho \vec{\Omega} \times \dot{\vec{x}} \right)$$



$$\nabla p = \vec{f}$$

Hydrostatische
Gleichg.

$$\vec{f} = -\nabla \psi^*$$

ψ^* ist das Potential
des Volumenkrp.

z.B.

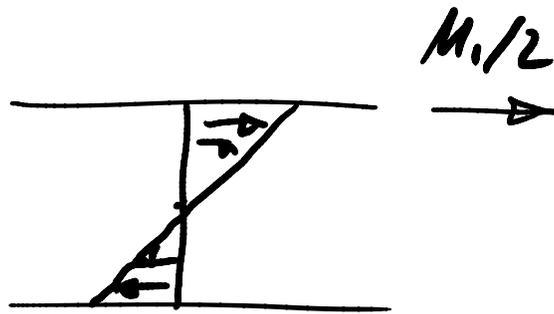
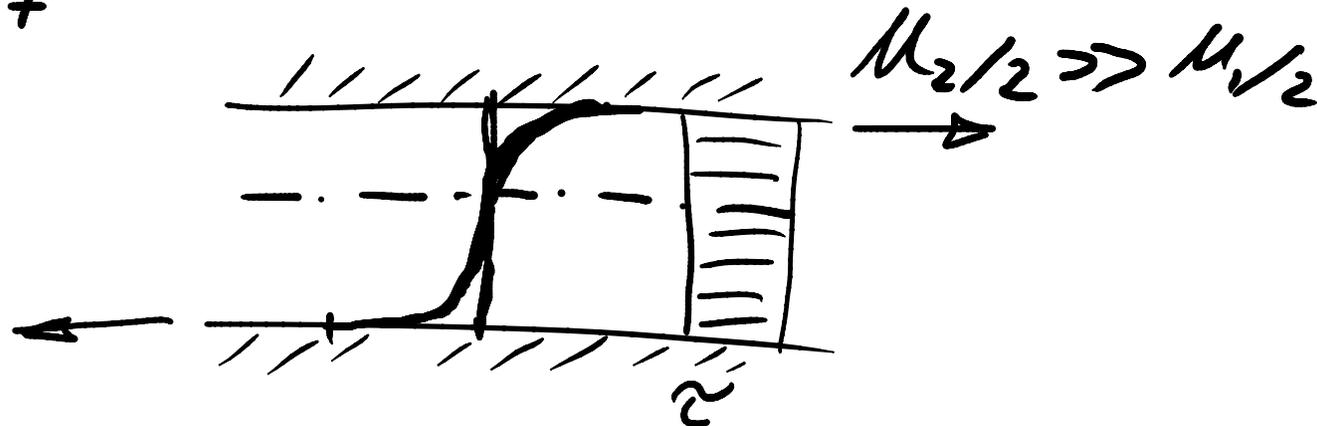
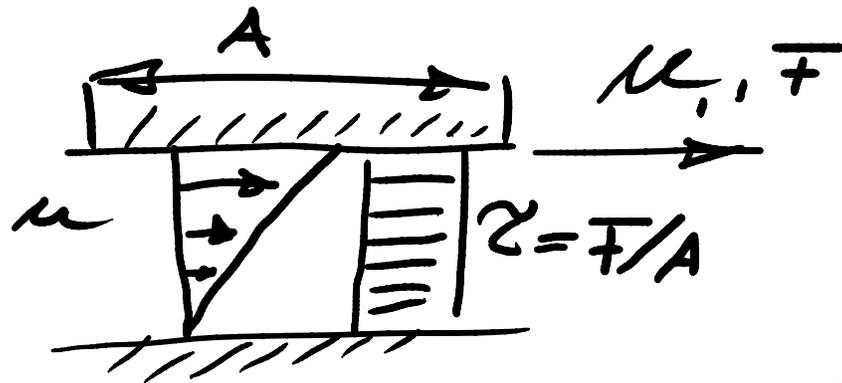
$$\psi^* = \rho g z + \text{const.} \quad \text{Potential des Schw.}$$

$$\psi^* = \frac{1}{2} \rho \Omega^2 r^2 + \text{const} \quad \text{Potential des zentrifugalen Krp.}$$

$$\underline{\underline{\rho + \psi^* = \text{const}}}$$



Geschwindigkeitsverteilung im Spalten
infolge schlepp Vöndel und infolge
Druckverl. ρ .

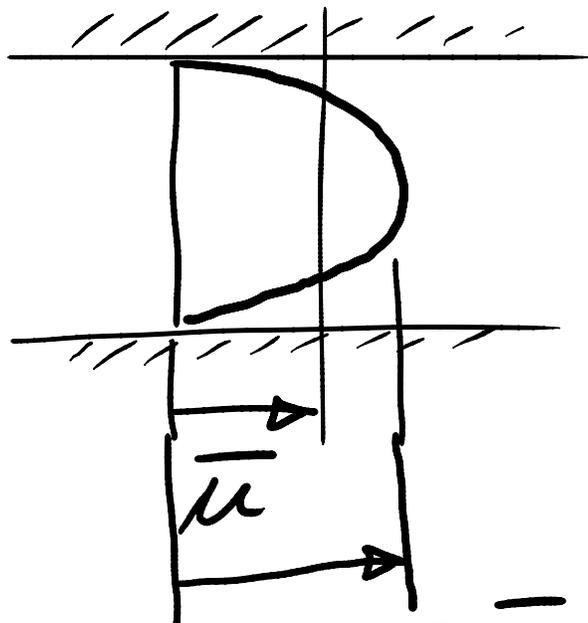


$-u_1/2$ laminare Couette Str.

turbulente Couette Str.



2D



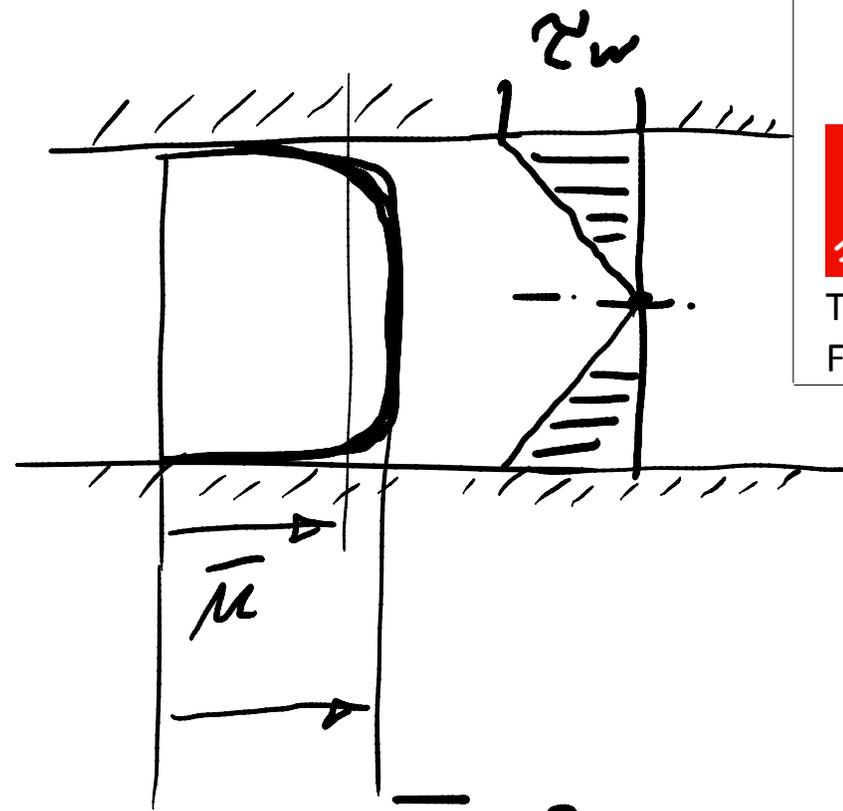
$$U_{max} = \frac{3}{2} \bar{U}$$

Anm. im Rohr ist

$$U_{max} = 2 \bar{U}$$

Poiseuille-Strömung

od Druckström.



$$U_{max} = \bar{U} + \frac{3}{2} \frac{1}{\kappa} U_*$$

$\kappa = 0.4$ Kármán'sche Konstante

$$U_* = \sqrt{\tau_w / \rho} \text{ Schubspannungsgeschw.}$$