

Theory of electromagnetic interactions: from few- to many-body systems

Sonia Bacca

October 3rd, 2017

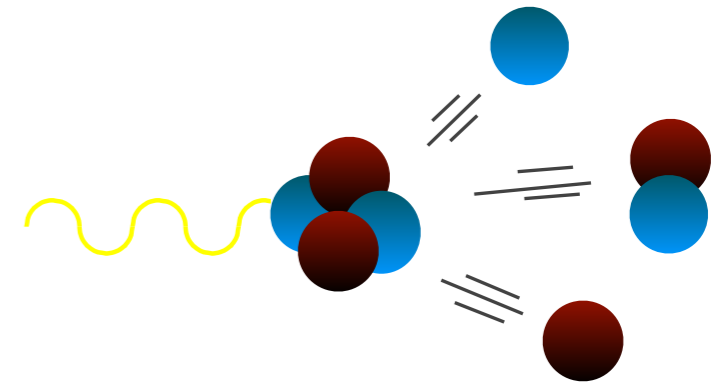
2nd Workshop of the SFB 1245
Schloss Waldhausen, Budenheim

- Electromagnetic probes (coupling constant $\ll 1$)

"With the electromagnetic probe, we can immediately relate the cross section to the transition matrix element of the current operator, thus to the structure of the target itself"

[De Forest-Walecka, Ann. Phys. 1966]

$$\sigma \propto |\langle \Psi_f | J^\mu | \Psi_0 \rangle|^2$$

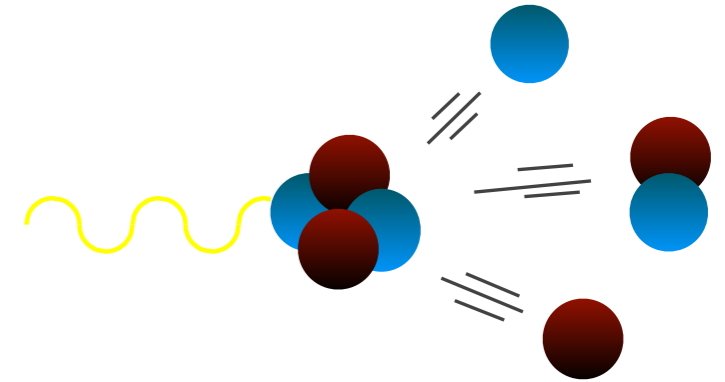


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- With ab-initio calculations we can study these observables and provide predictions with error bars, first for light nuclei and then for heavier nuclei

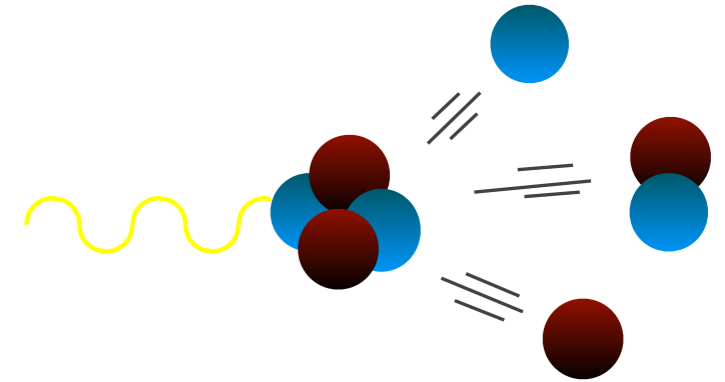
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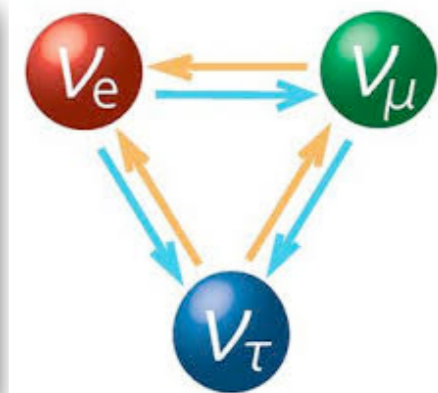


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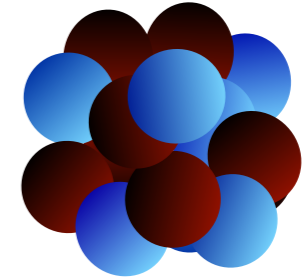
“from few- to many-body systems”

- Provide important informations in other fields of physics, where nuclear physics plays a crucial role:

- Astrophysics:
- Atomic physics
- Particle physics



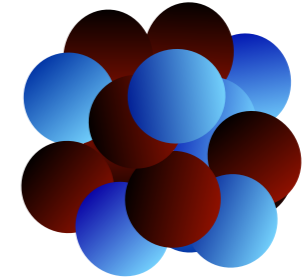
- Start from neutrons and protons as building blocks (centre of mass coordinates, spins, isospins)



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- Solve the non-relativistic quantum mechanical problem of A -interacting nucleons

$$H|\psi_i\rangle = E_i|\psi_i\rangle$$

$$H = T + V_{NN}(\Lambda) + V_{3N}(\Lambda) + \dots$$

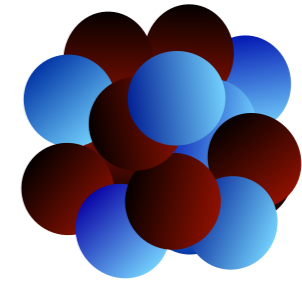


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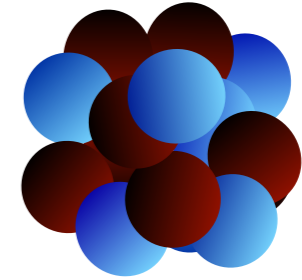
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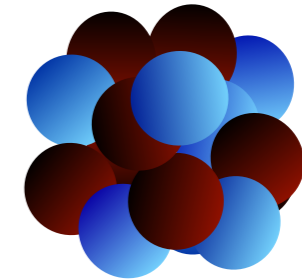
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Chiral Effective Field Theory

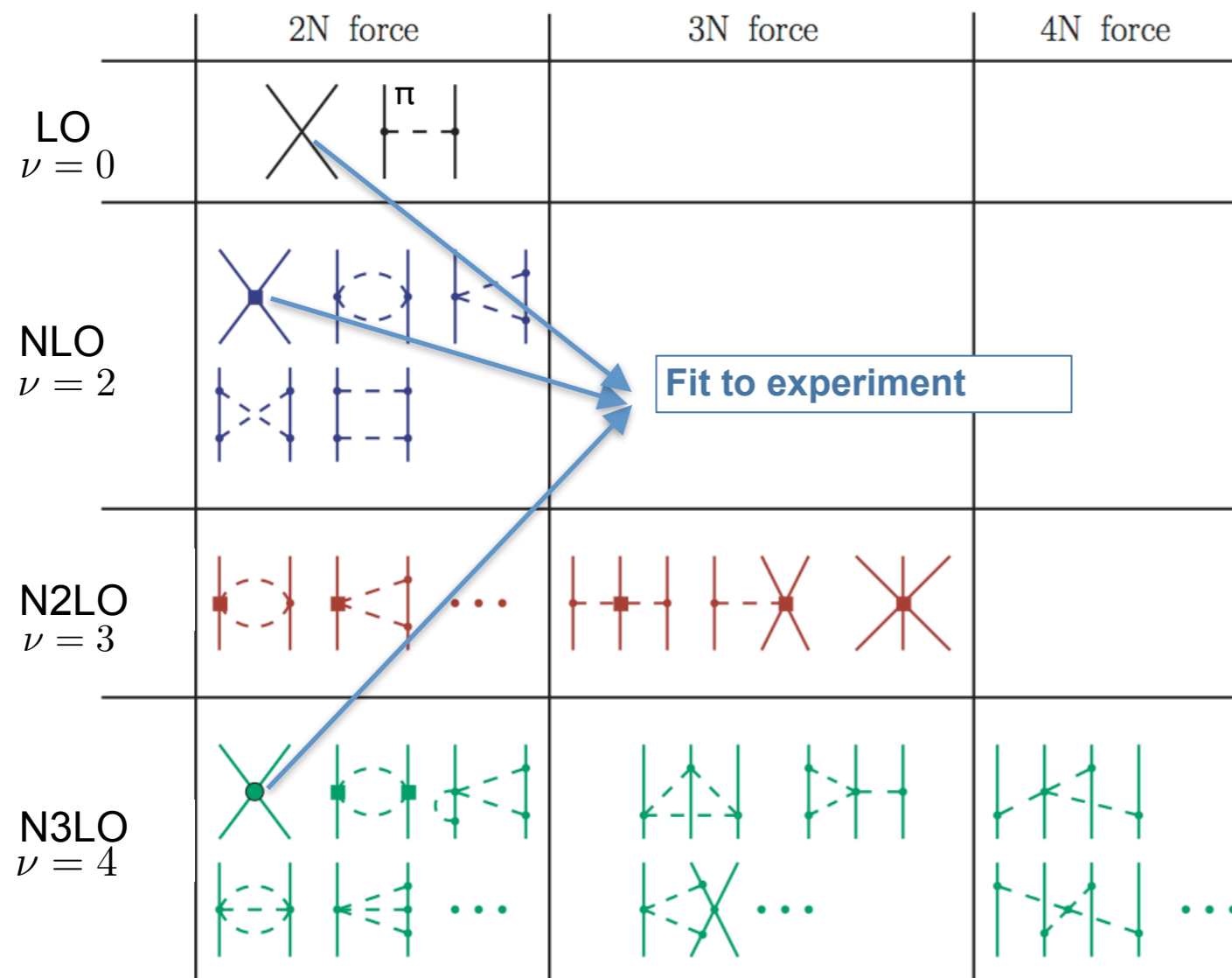
Weinberg, van Kolck, Epelbaum, Meissner, Machleidt

	2N force	3N force	4N force
LO $\nu = 0$			
NLO $\nu = 2$			
N2LO $\nu = 3$			
N3LO $\nu = 4$			

Systematic expansion $\mathcal{L} = \sum_{\nu} c_{\nu} \left(\frac{Q}{\Lambda_b} \right)^{\nu}$

Chiral Effective Field Theory

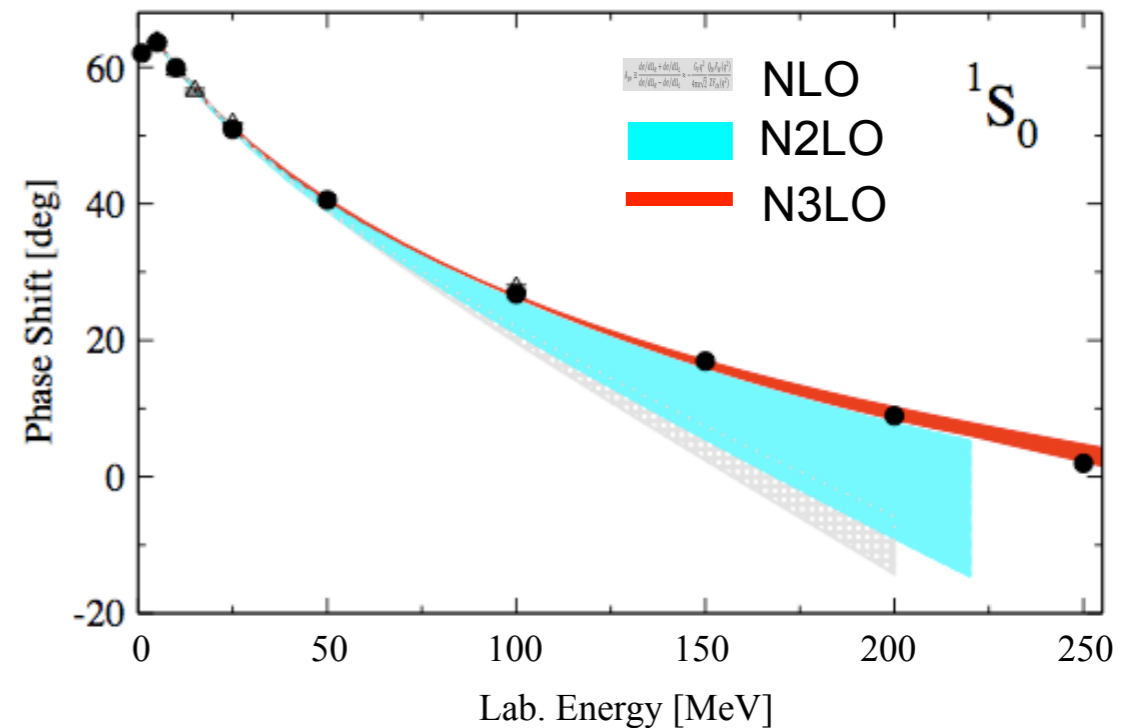
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Systematic expansion $\mathcal{L} = \sum_{\nu} c_{\nu} \left(\frac{Q}{\Lambda_b} \right)^{\nu}$

Details of short distance physics not resolved, but captured in **low energy constants (LEC)**

LEC fit to experiment - NN sector -

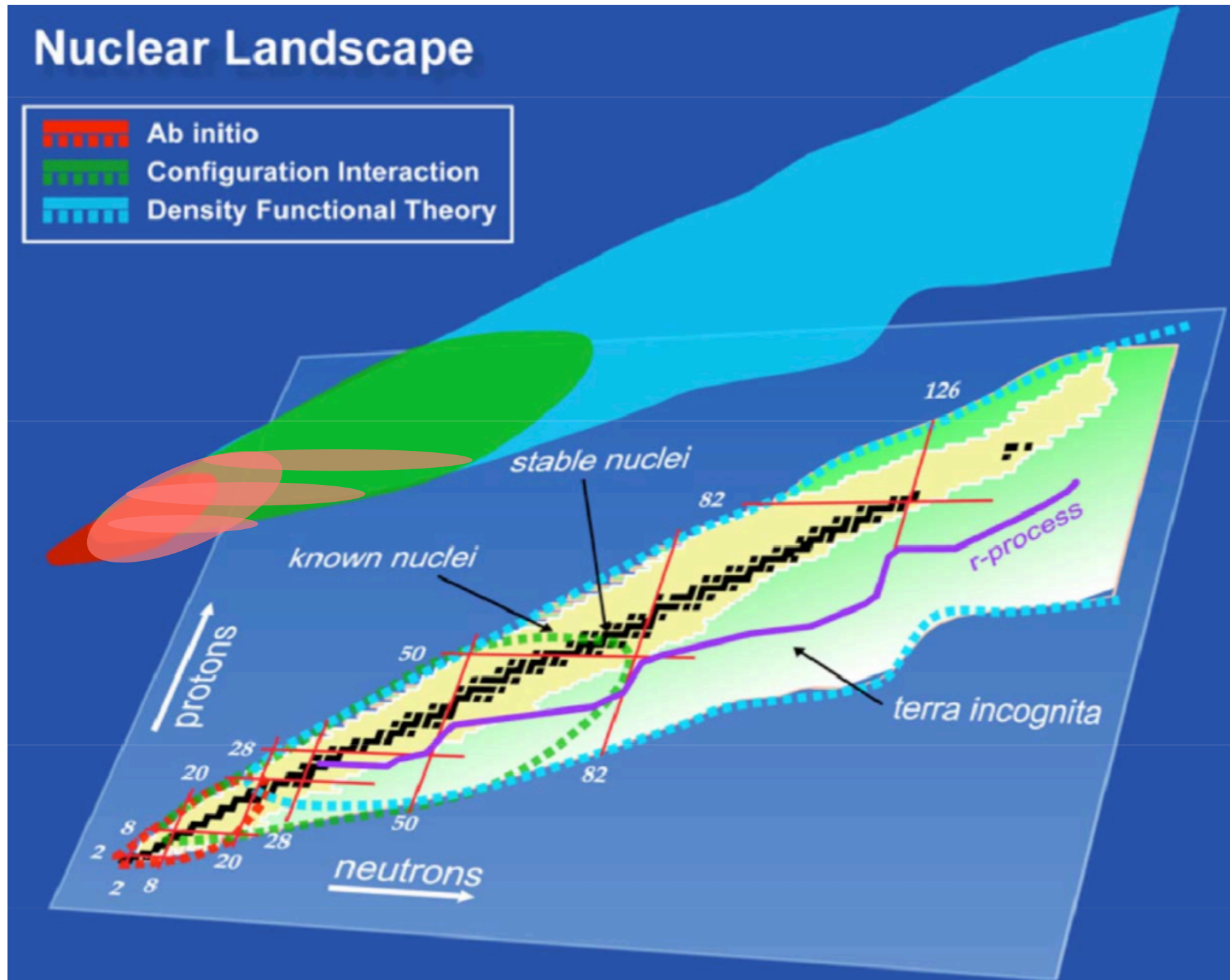


Epelbaum *et al.* (2009)

LEC fit to experiment - 3N sector -

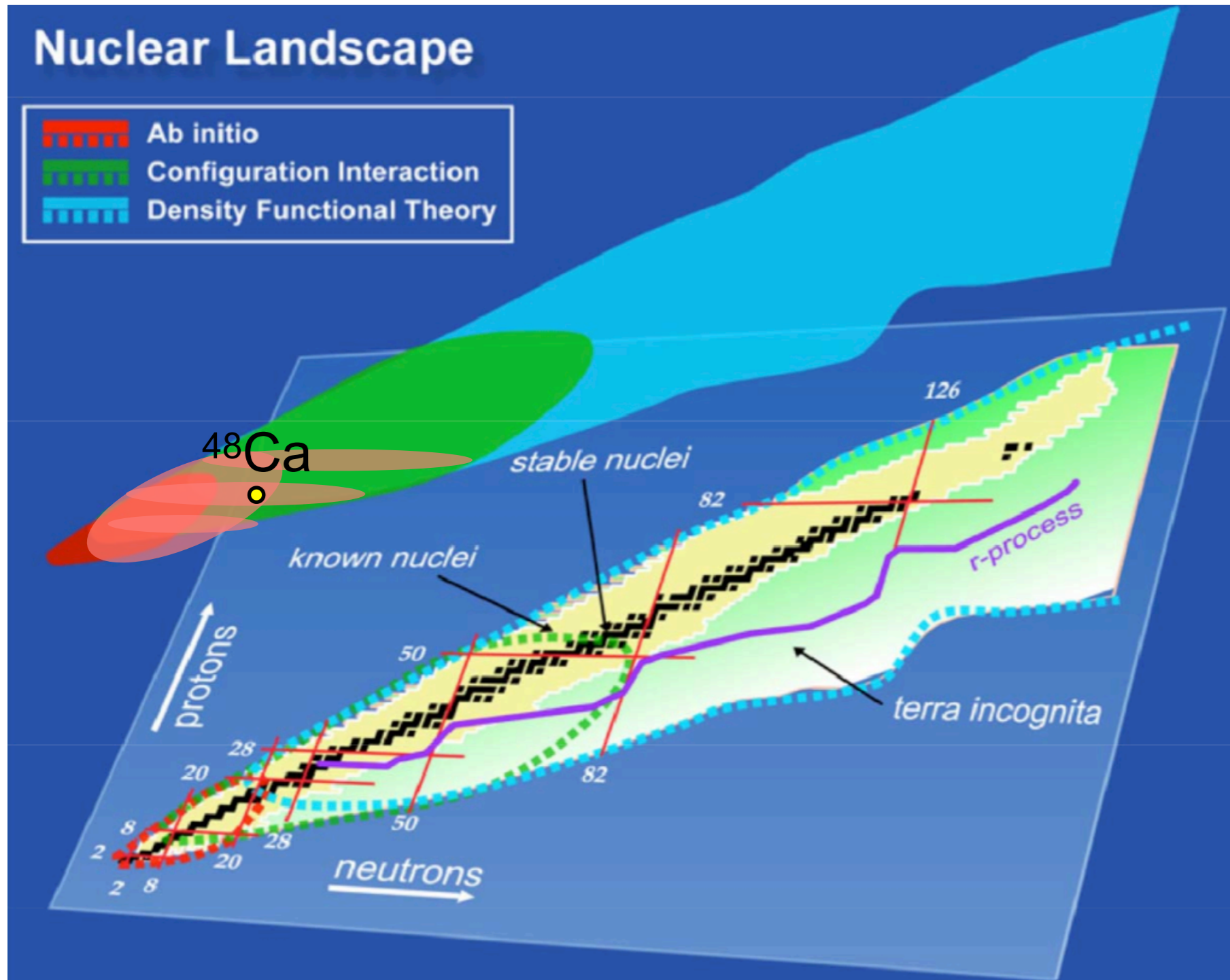
Using A=3 data or including other A>3 nuclei

Various methods to solve the many-body problem



<http://unedf.org>

Various methods to solve the many-body problem

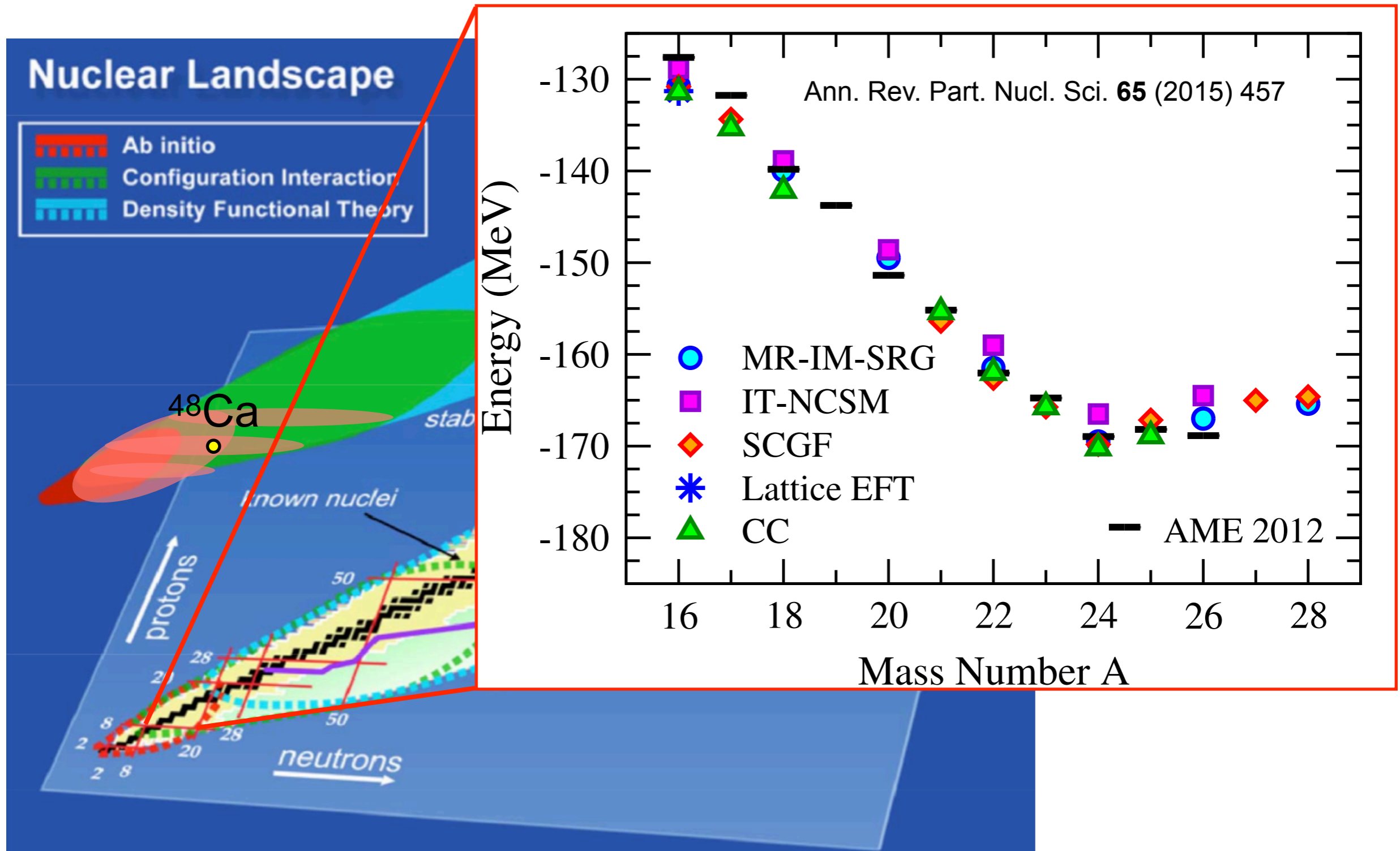


<http://unedf.org>

Nuclear Structure Theory

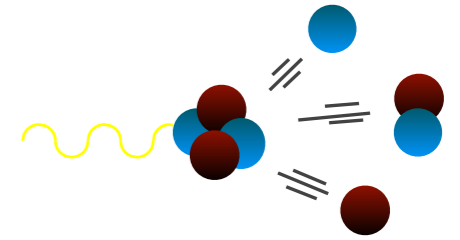
Various methods to solve the many-body problem

Oxygen chain



<http://unedf.org>

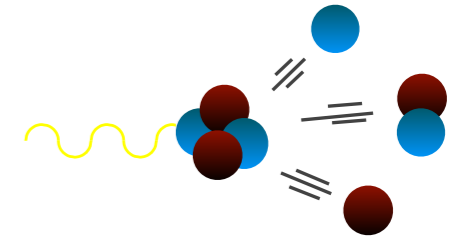
How does the nucleus respond to external electromagnetic excitations?



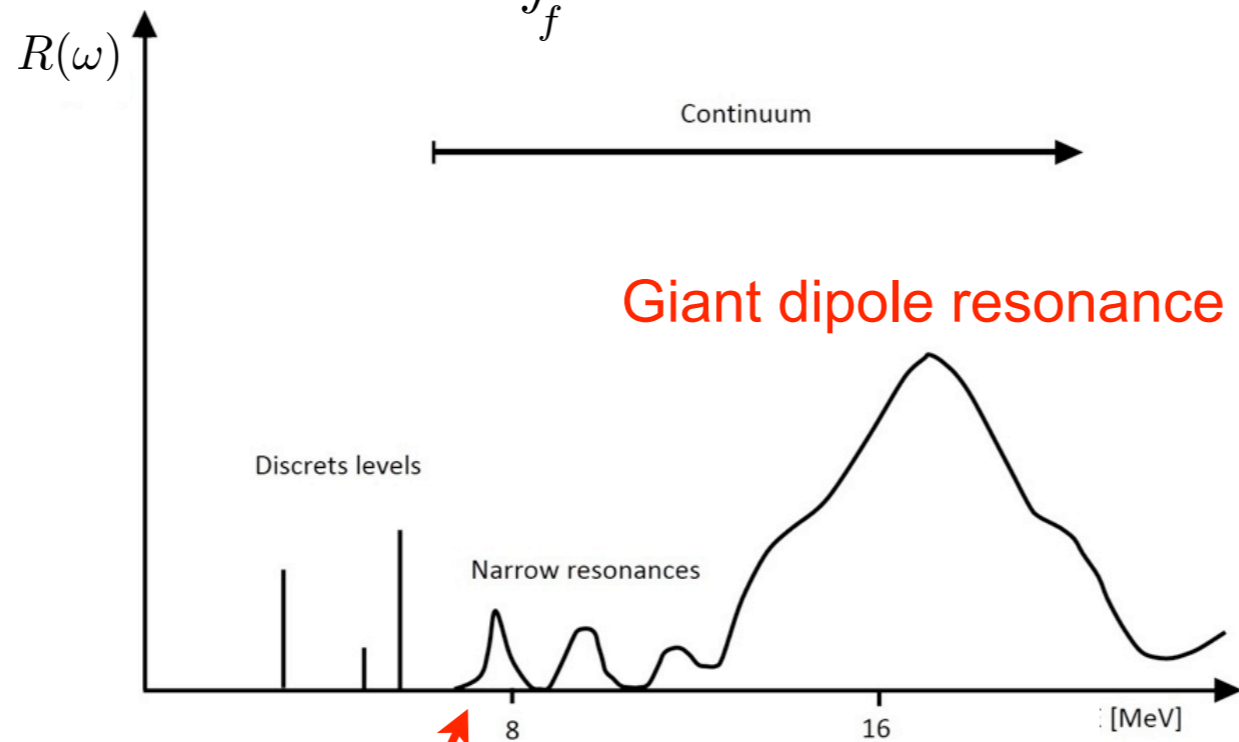
$$R(\omega) = \sum_f |\langle \Psi_f | \mathcal{O} | \Psi_0 \rangle|^2 \delta(\omega - E_f + E_0)$$

Nuclear Reactions

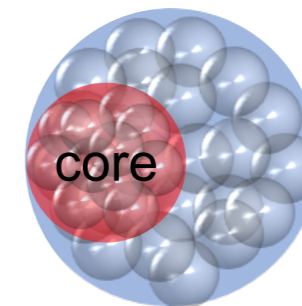
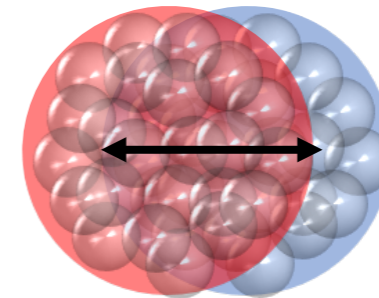
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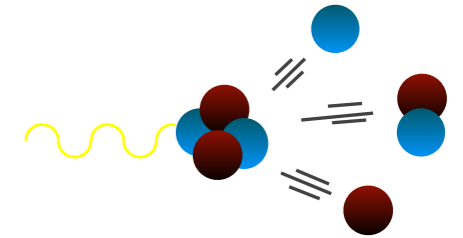


Pigmy dipole resonance in neutron-rich nuclei

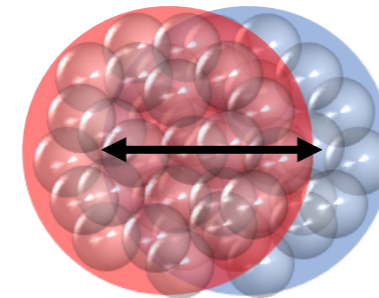
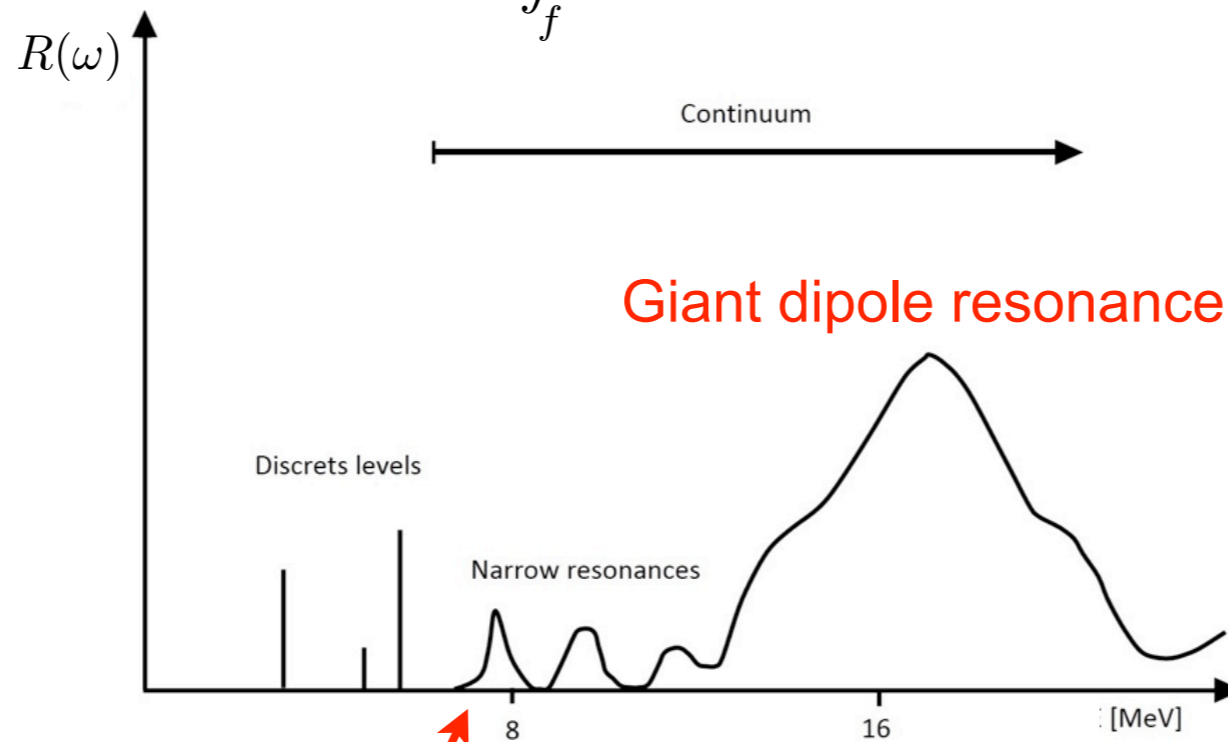


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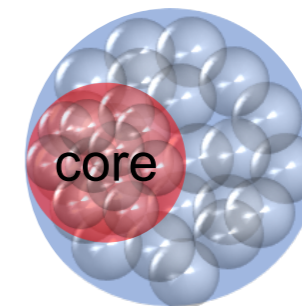
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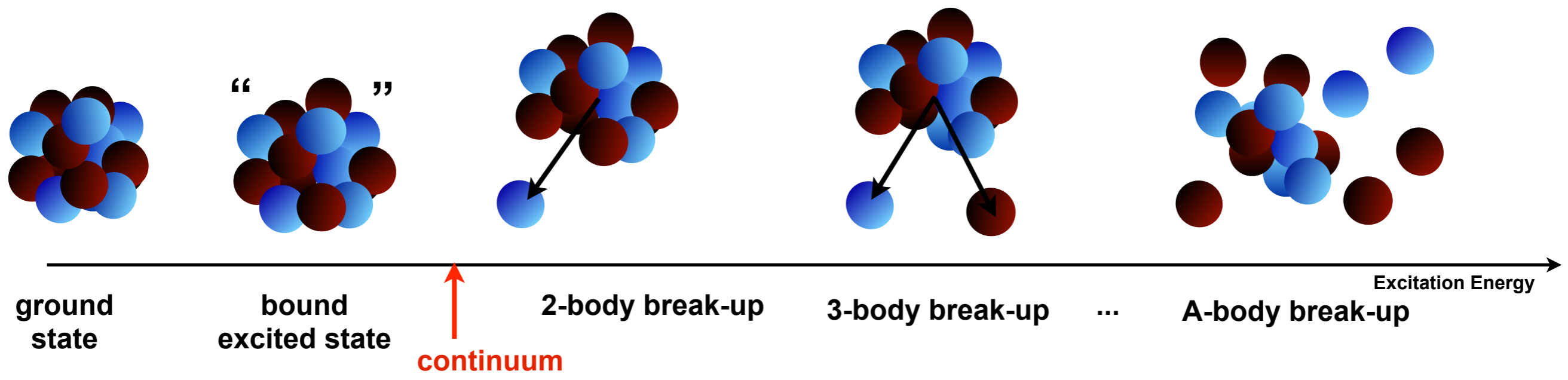
Pigmy dipole resonance in neutron-rich nuclei



$$\alpha_D = 2\alpha \int_{\omega_{th}}^{\infty} d\omega \frac{R(\omega)}{\omega}$$

→ Low-energy part of strength dominates

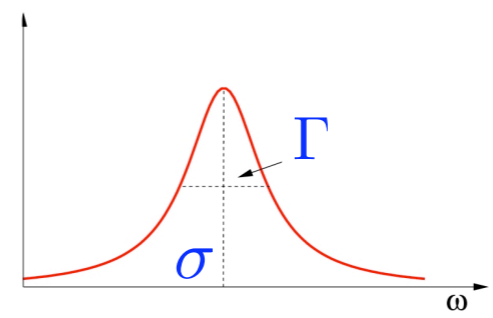
The Continuum Problem



$$R(\omega) \propto |\langle \Psi_f | \mathcal{O} | \Psi_0 \rangle|^2$$

Exact knowledge limited in energy and mass number

$$L(\sigma, \Gamma) = \int d\omega \frac{R(\omega)}{(\omega - \sigma)^2 + \Gamma^2} = \langle \tilde{\psi} | \tilde{\psi} \rangle < \infty$$



Lorentz Integral Transform → Reduce the continuum problem to a bound-state-like equation

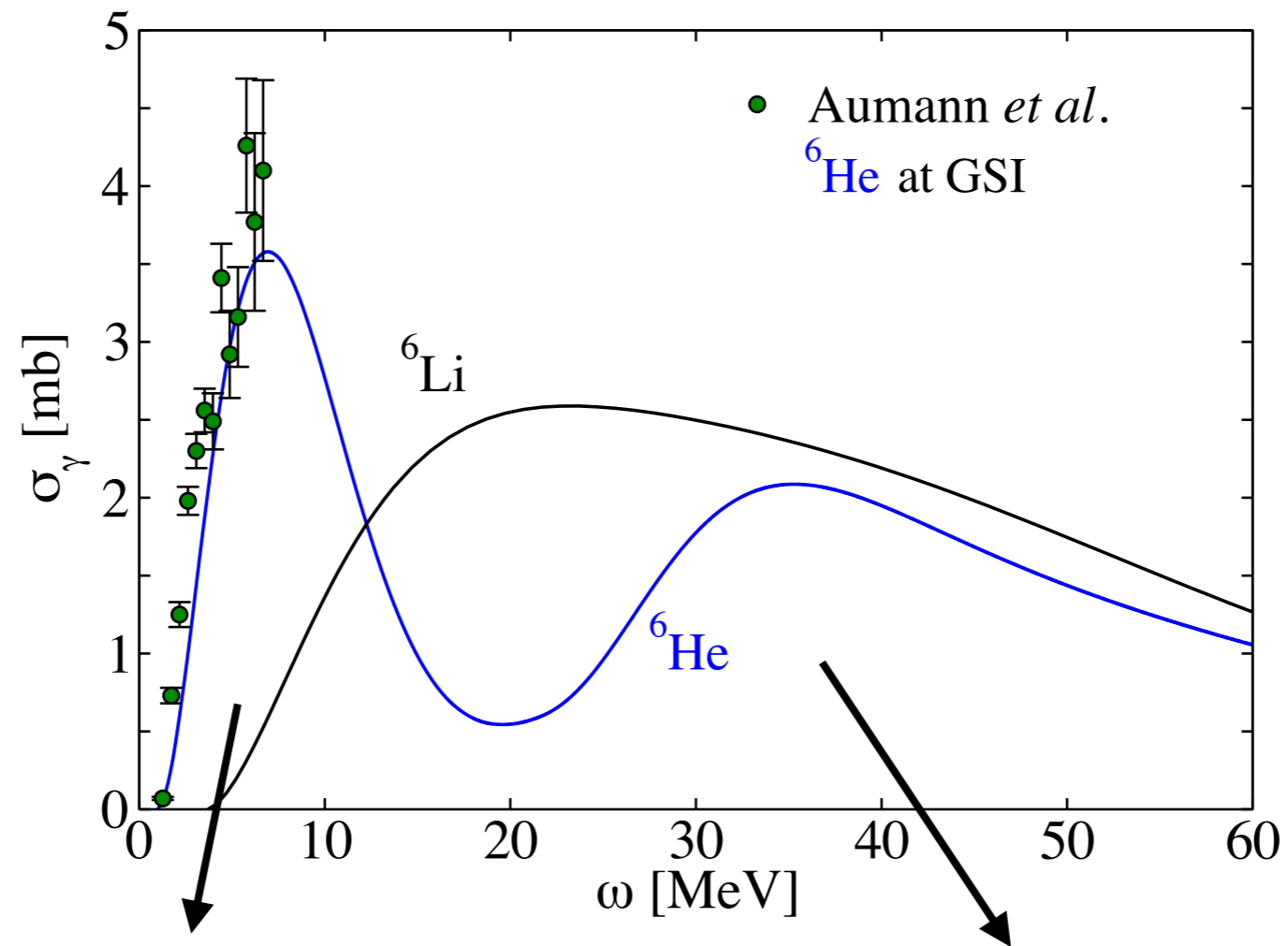
Efros, *et al.*, JGP.: Nucl.Part.Phys. **34** (2007) R459

$$(H - E_0 - \sigma + i\Gamma) | \tilde{\psi} \rangle = \hat{O} | \psi_0 \rangle$$

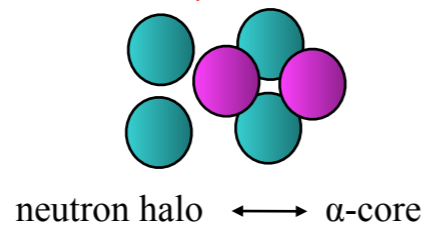
***Few-body nuclei
with hyper-spherical harmonic expansions***

Dipole Response Function

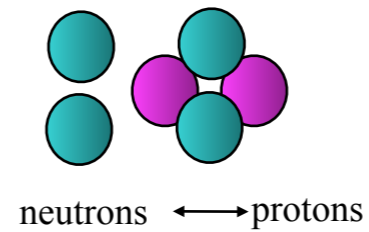
S.Bacca et al, PRL **89** 052502 (2002)



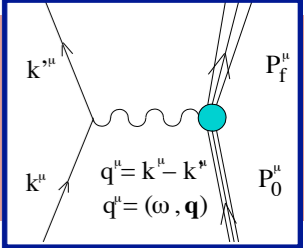
Soft-dipole Mode



Giant Dipole Mode

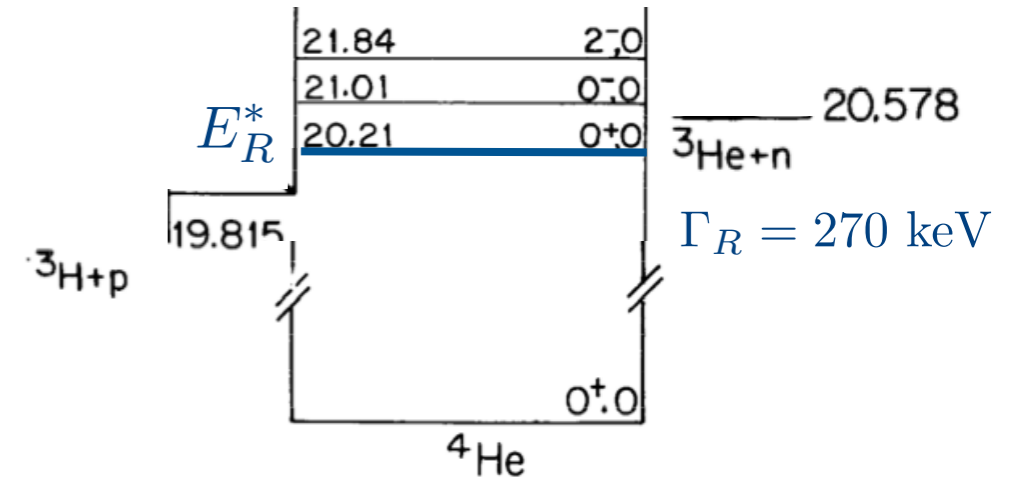


Monopole Resonance ${}^4\text{He}(e,e')0^+$

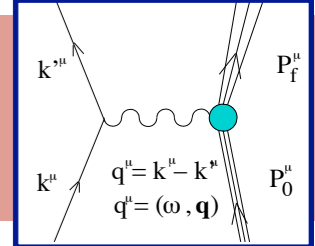


Resonant Transition Form Factor
 $0_1^+ \longrightarrow 0_2^+$

$$|F_{\mathcal{M}}(q)|^2 = \frac{1}{Z^2} \int d\omega R_{\mathcal{M}}^{\text{res}}(q, \omega)$$



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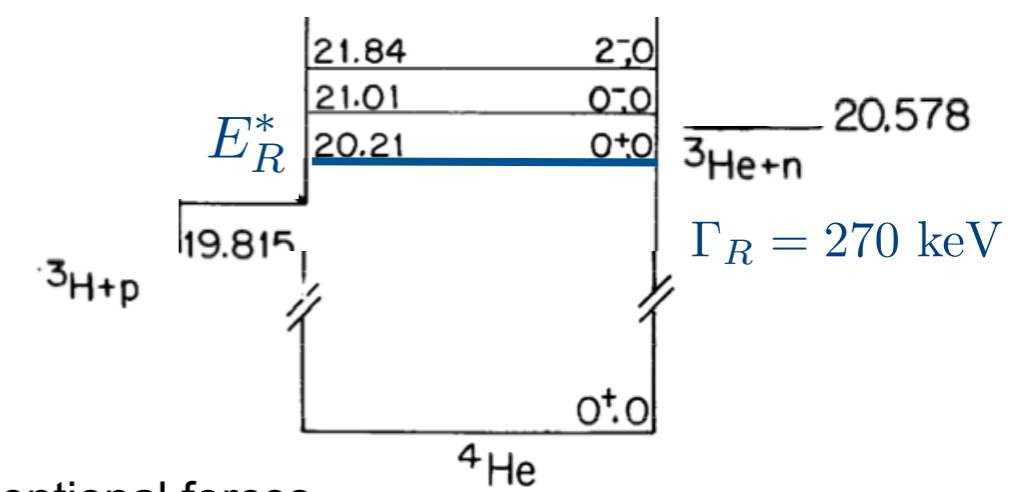
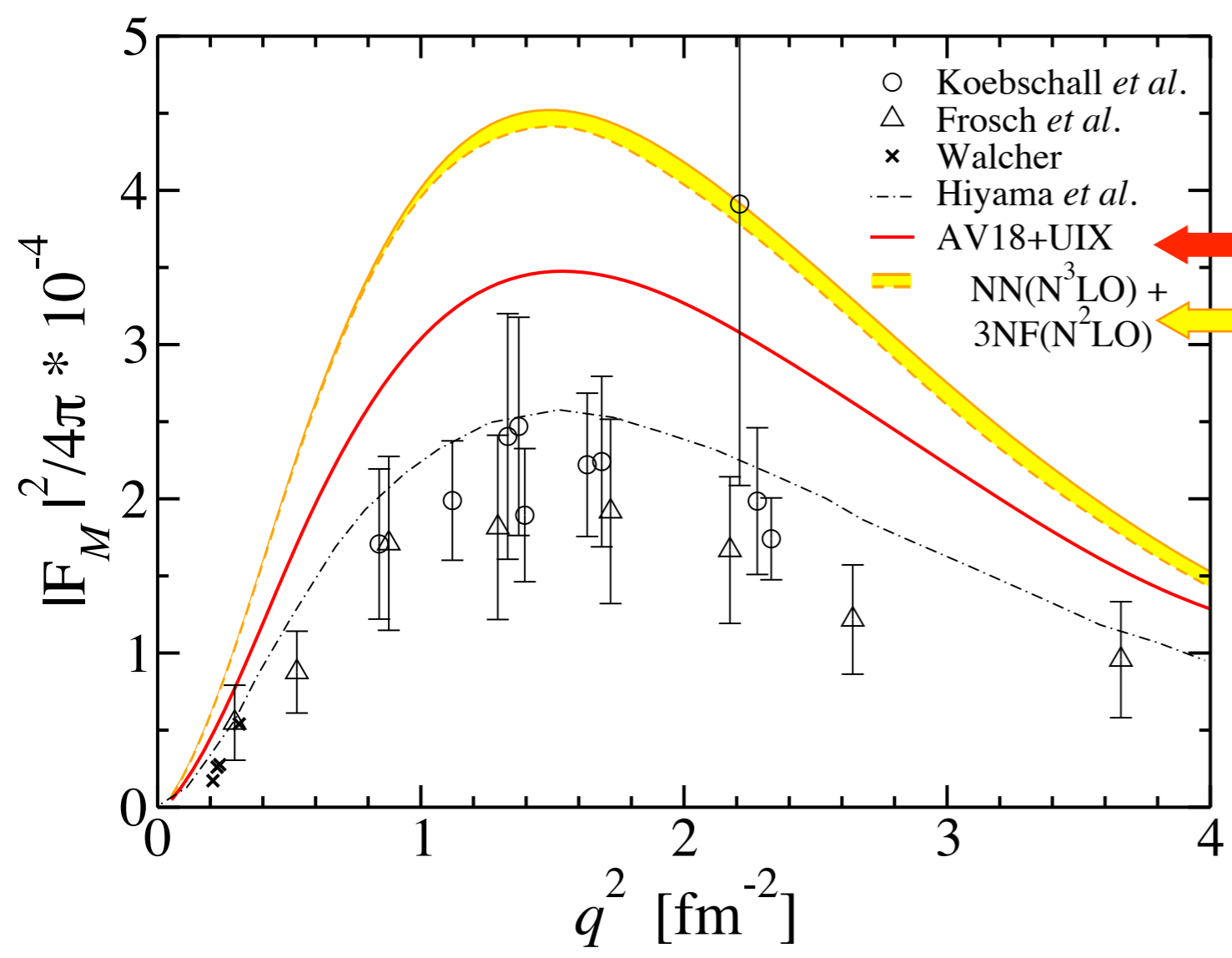


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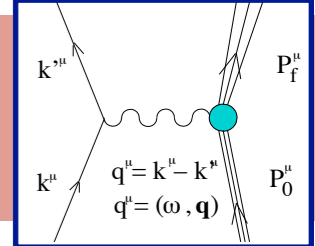
First ab-initio calculation with realistic three-nucleon forces

S.B. *et al.*, PRL **110**, 042503 (2013)



conventional forces
 χ EFT forces

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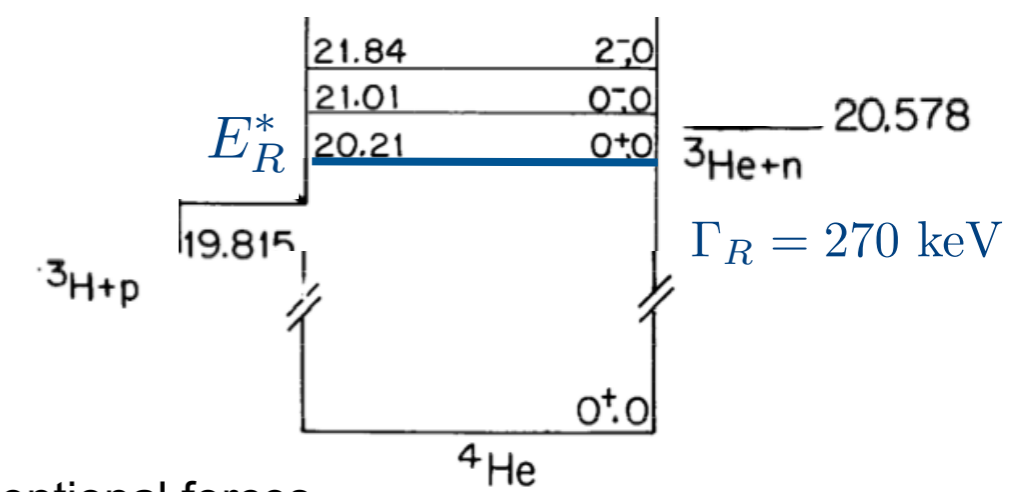
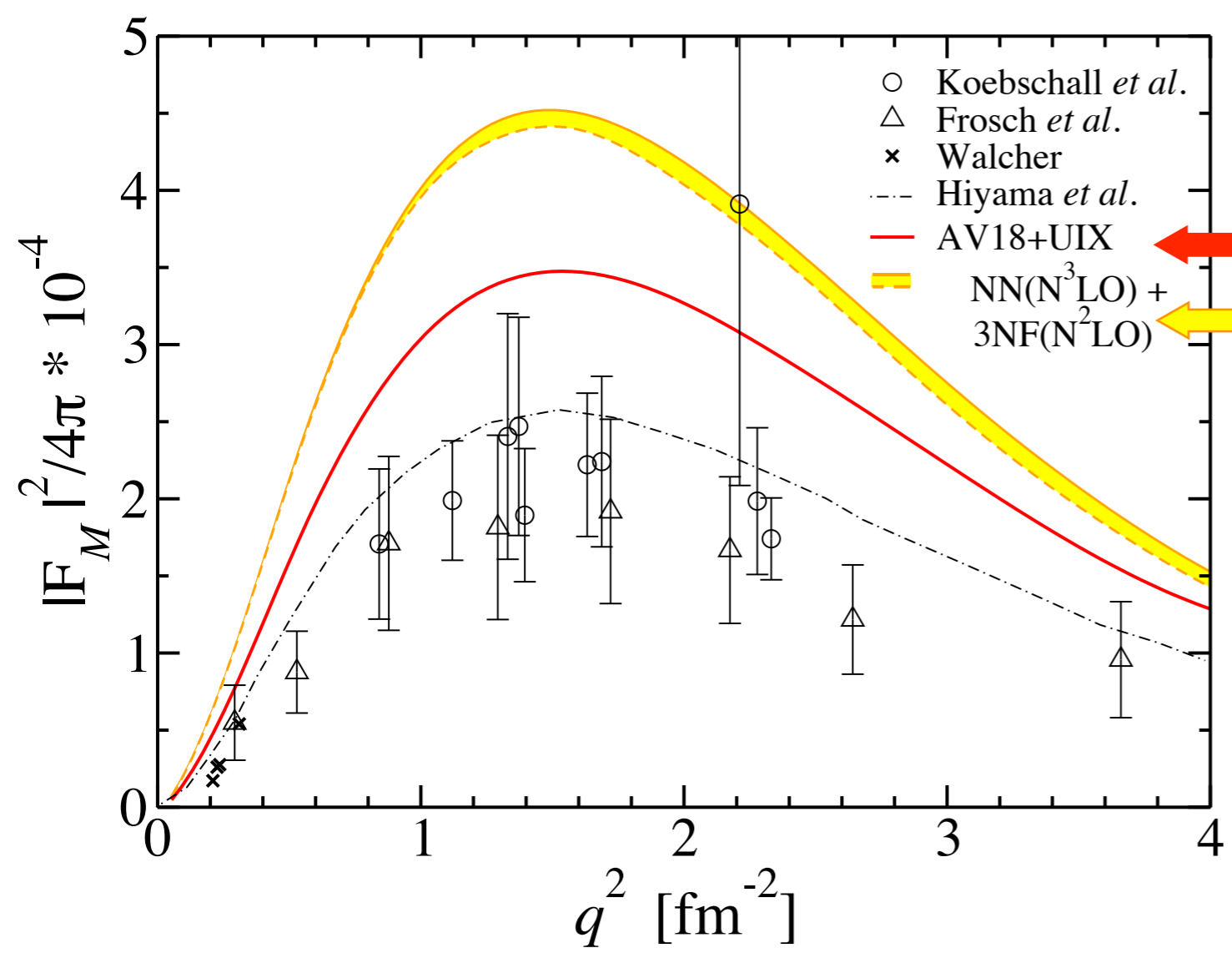


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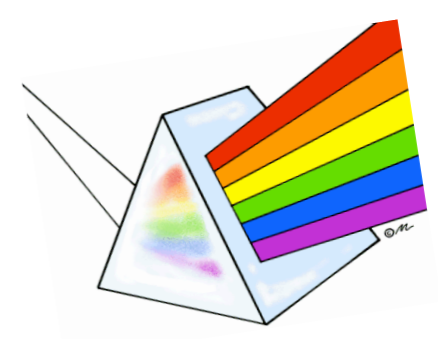
S.B. et al., PRL 110, 042503 (2013)



conventional forces
 χ EFT forces

AV8' + central 3NF	$E_0 = -28.44$ MeV
AV18+UIX	$E_0 = -28.40$ MeV
NN(N ³ LO)+3NF(N ² LO)	$E_0 = -28.357$ MeV
	$E_0^{\text{exp}} = -28.30$ MeV

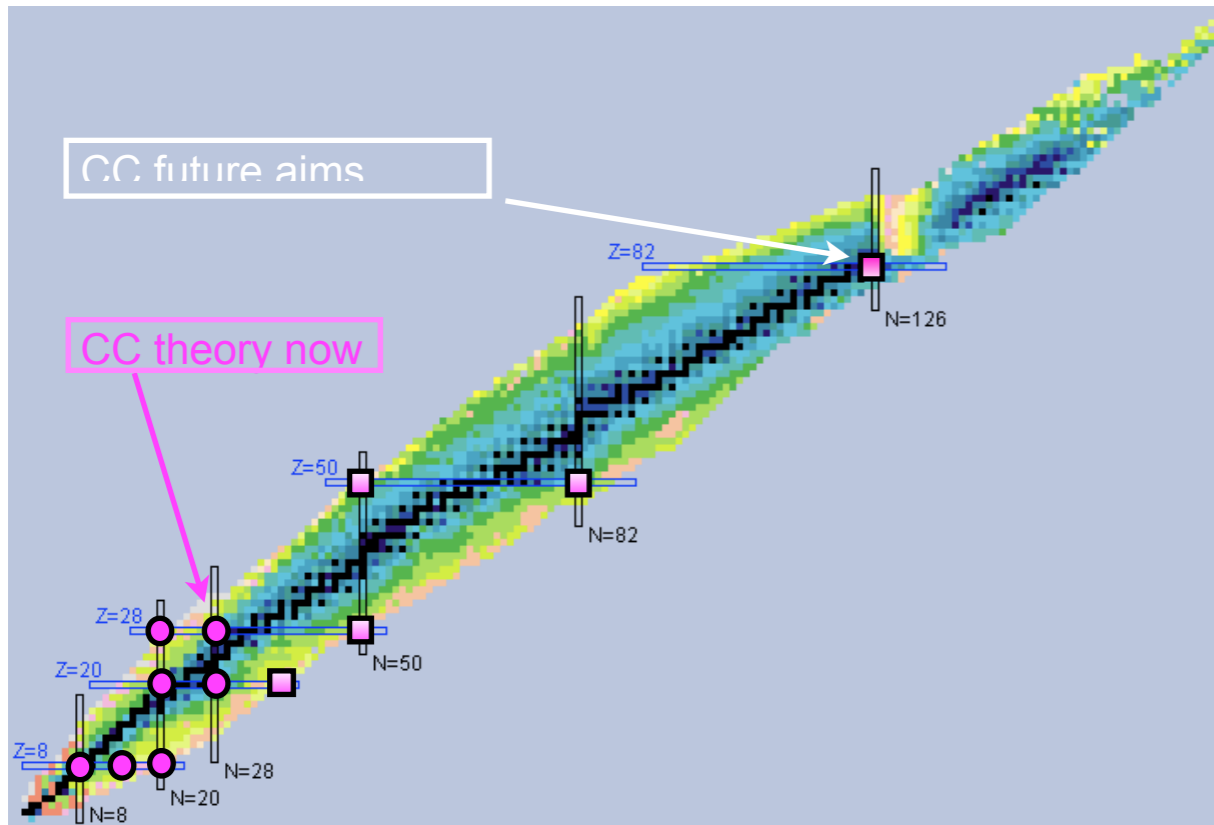
H reproduce
 E_0^{exp}



Pushing the limits in mass number

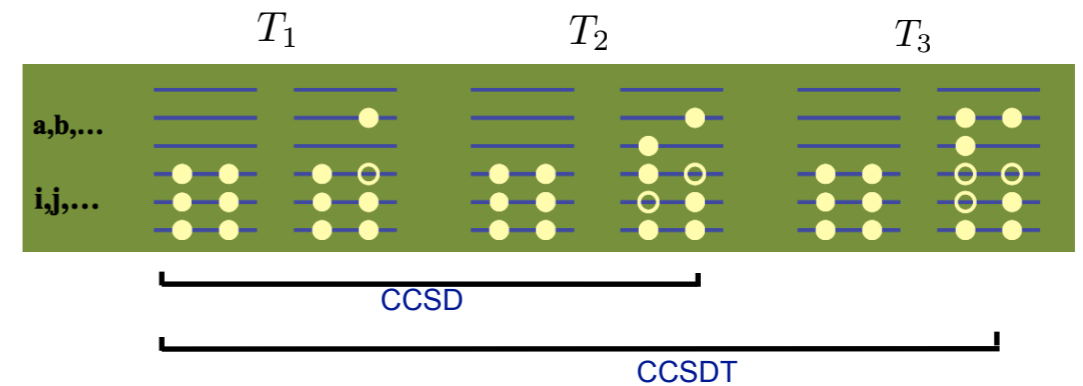
Coupled Cluster Theory

Many-body method that can extend the frontiers of ab-initio calculations to heavier and neutron nuclei



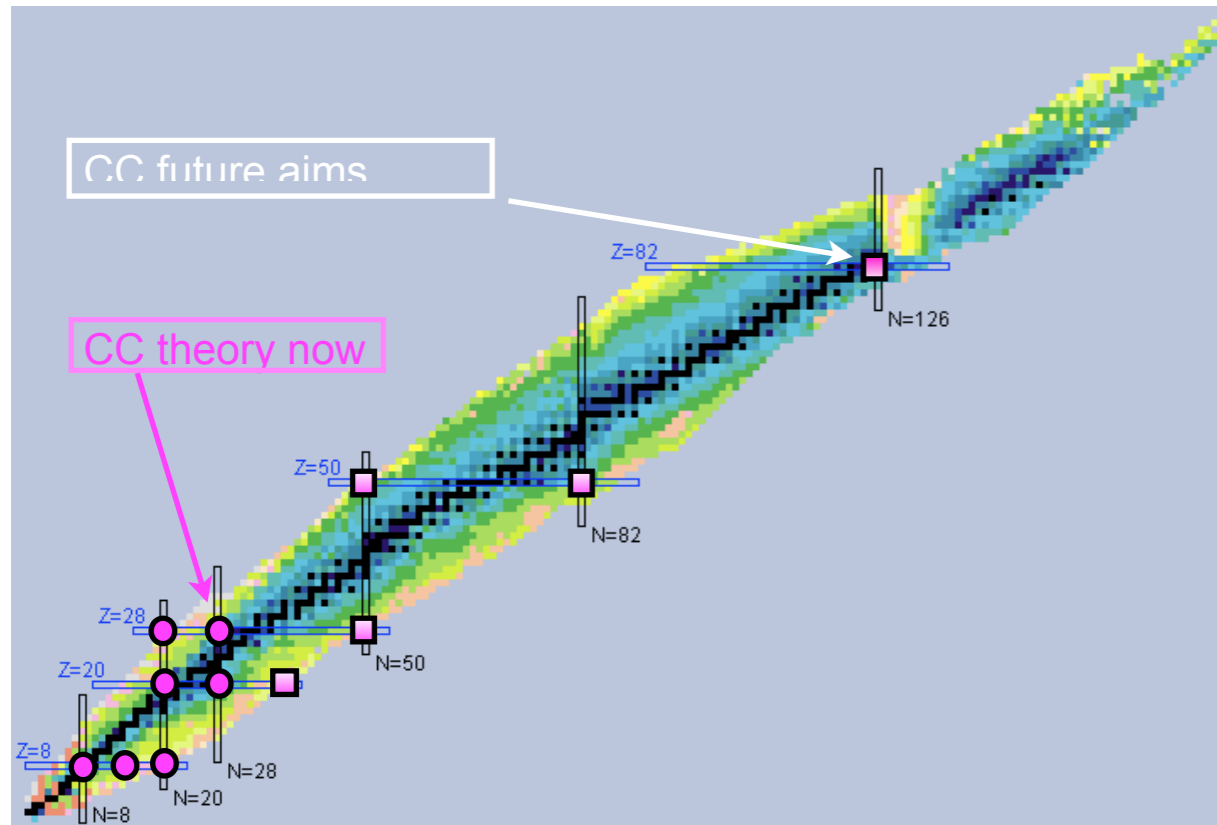
$$|\psi_0(\vec{r}_1, \vec{r}_2, \dots, \vec{r}_A)\rangle = e^T |\phi_0(\vec{r}_1, \vec{r}_2, \dots, \vec{r}_A)\rangle$$

$$T = \sum T_{(A)} \text{ cluster expansion}$$



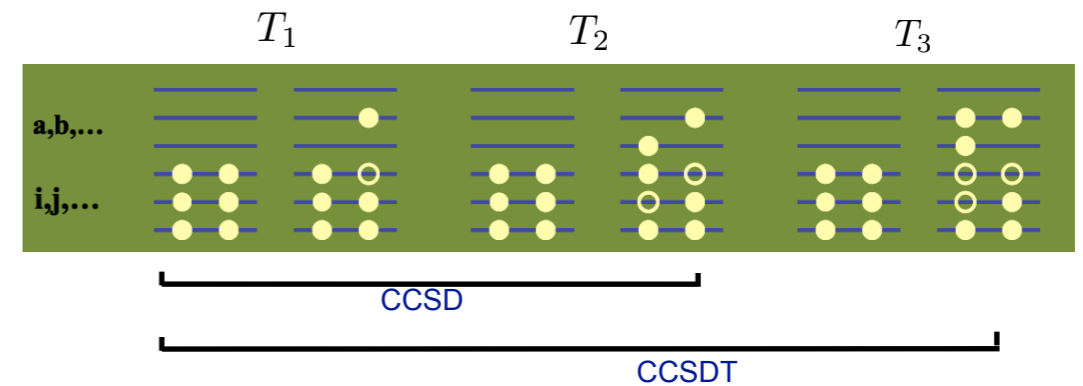
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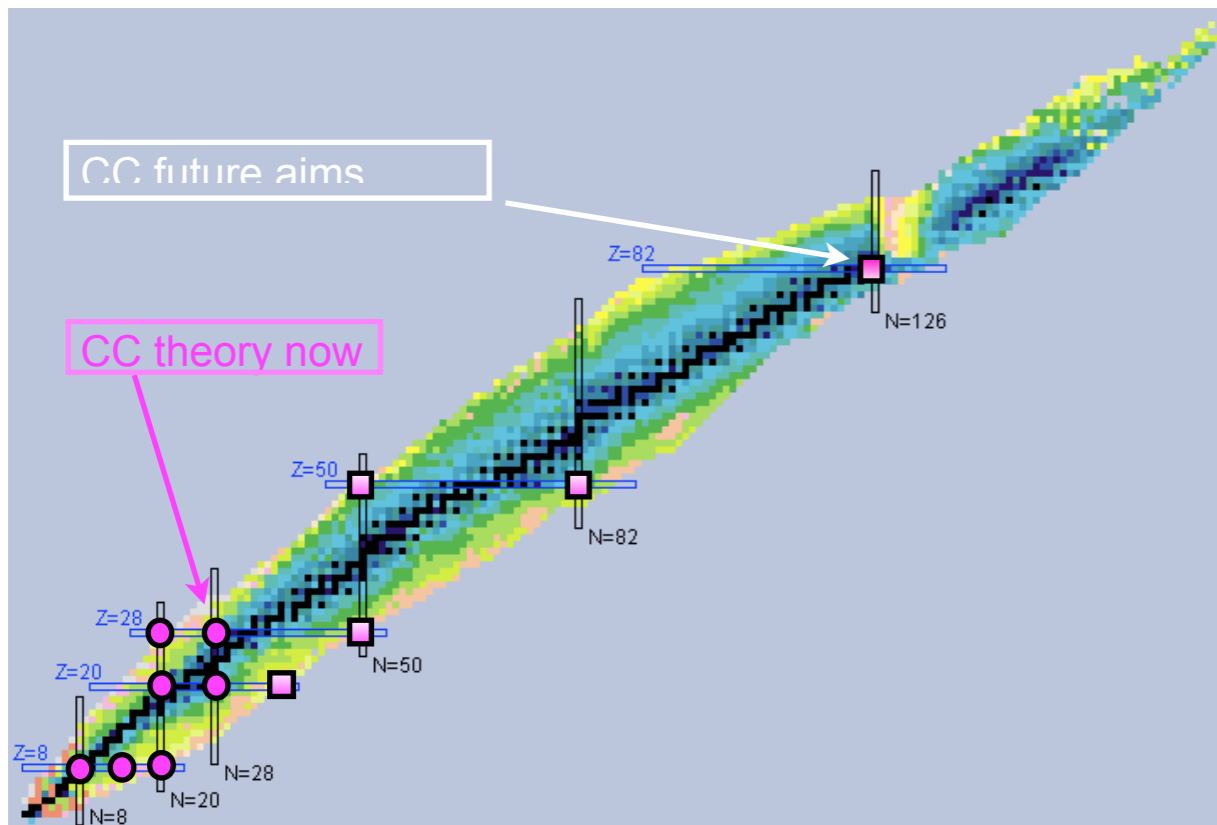
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Mature theory for bound states, but what about electromagnetic reactions?

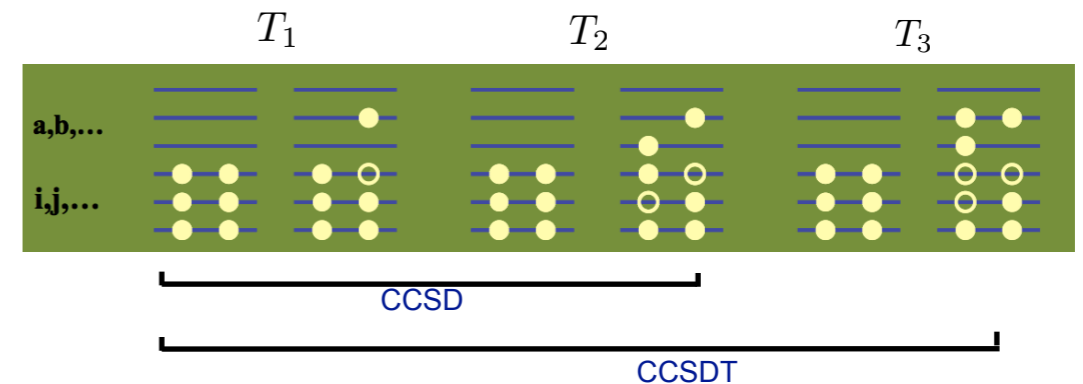
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Mature theory for bound states, but what about electromagnetic reactions?

Merge Lorentz integral transform method with coupled-cluster theory

$$(\bar{H} - E_0 - \sigma + i\Gamma)|\tilde{\Psi}_R\rangle = \bar{\Theta}|\Phi_0\rangle$$

$$\bar{H} = e^{-T} H e^T$$

$$\bar{\Theta} = e^{-T} \Theta e^T$$

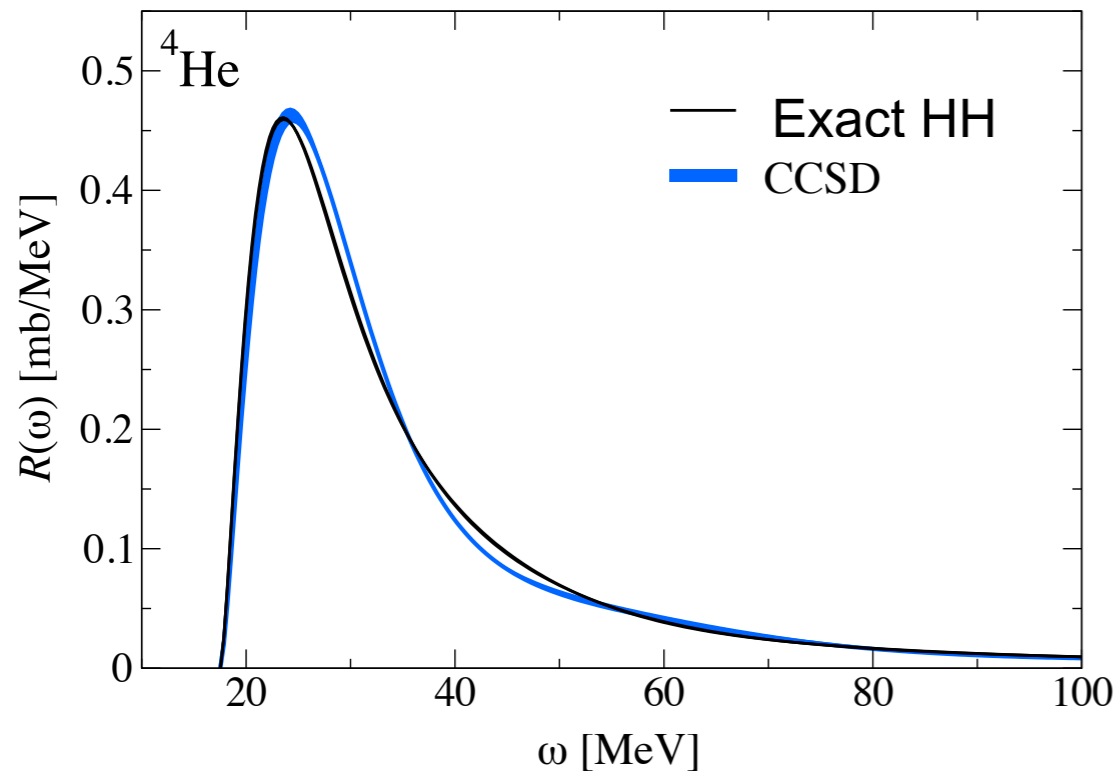
$$|\tilde{\Psi}_R\rangle = \hat{R}|\Phi_0\rangle$$

I will show mostly results at CCSD truncation scheme

S.B. *et al.*, Phys. Rev. Lett. **111**, 122502 (2013)

Dipole Response Functions with NN forces from χ EFT (N³LO)

Validation ⁴He

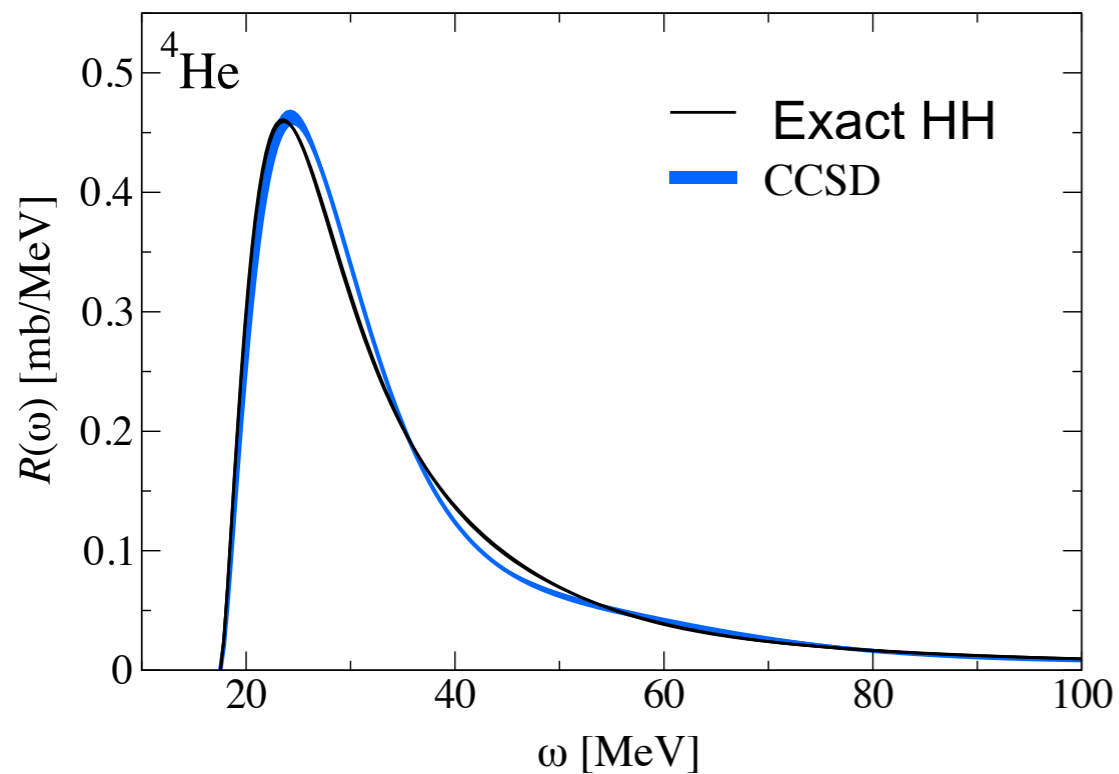


Photoexcitation of stable nuclei

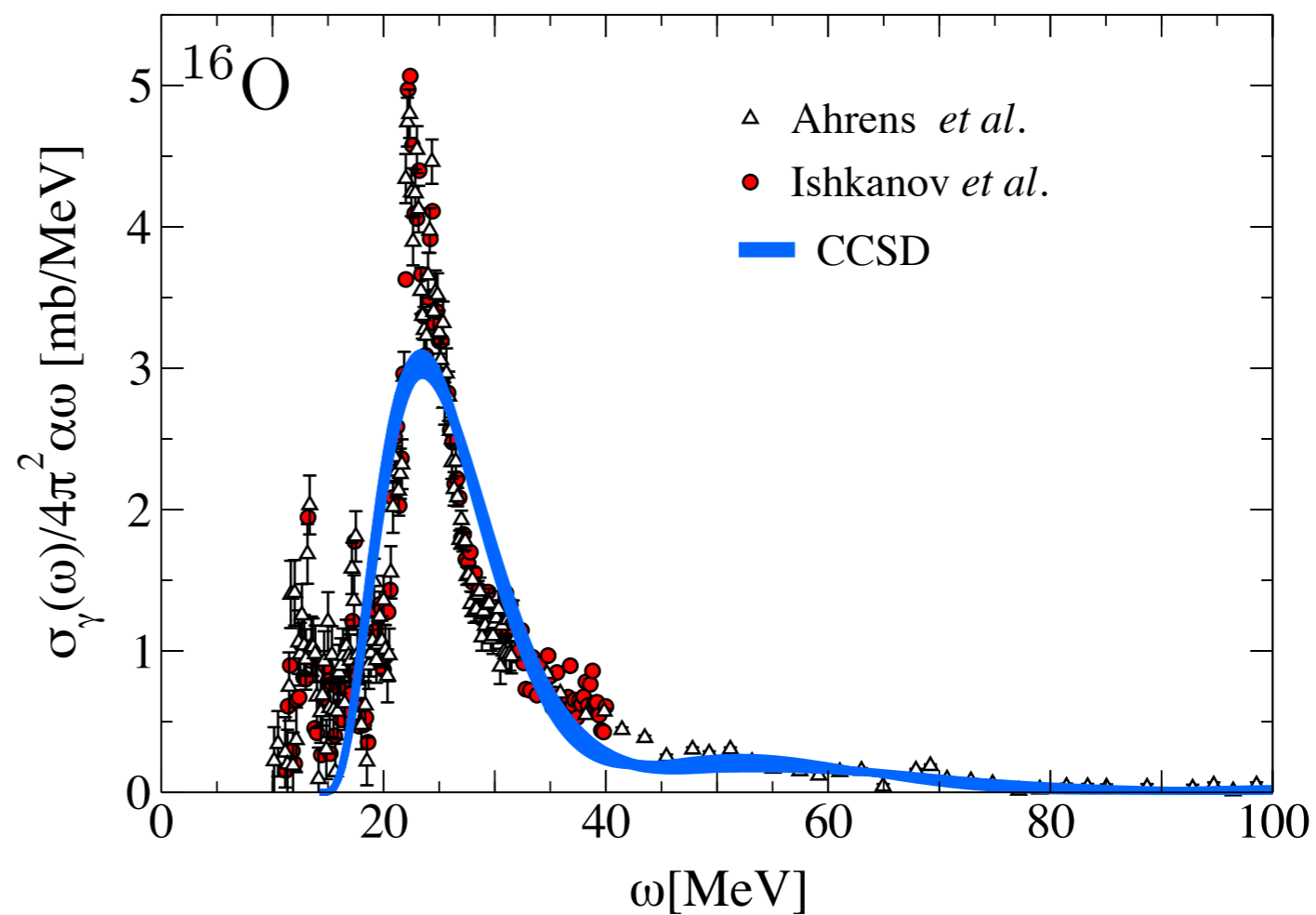
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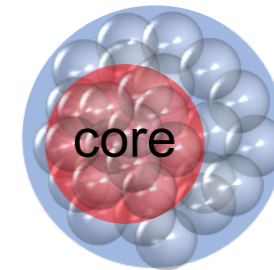
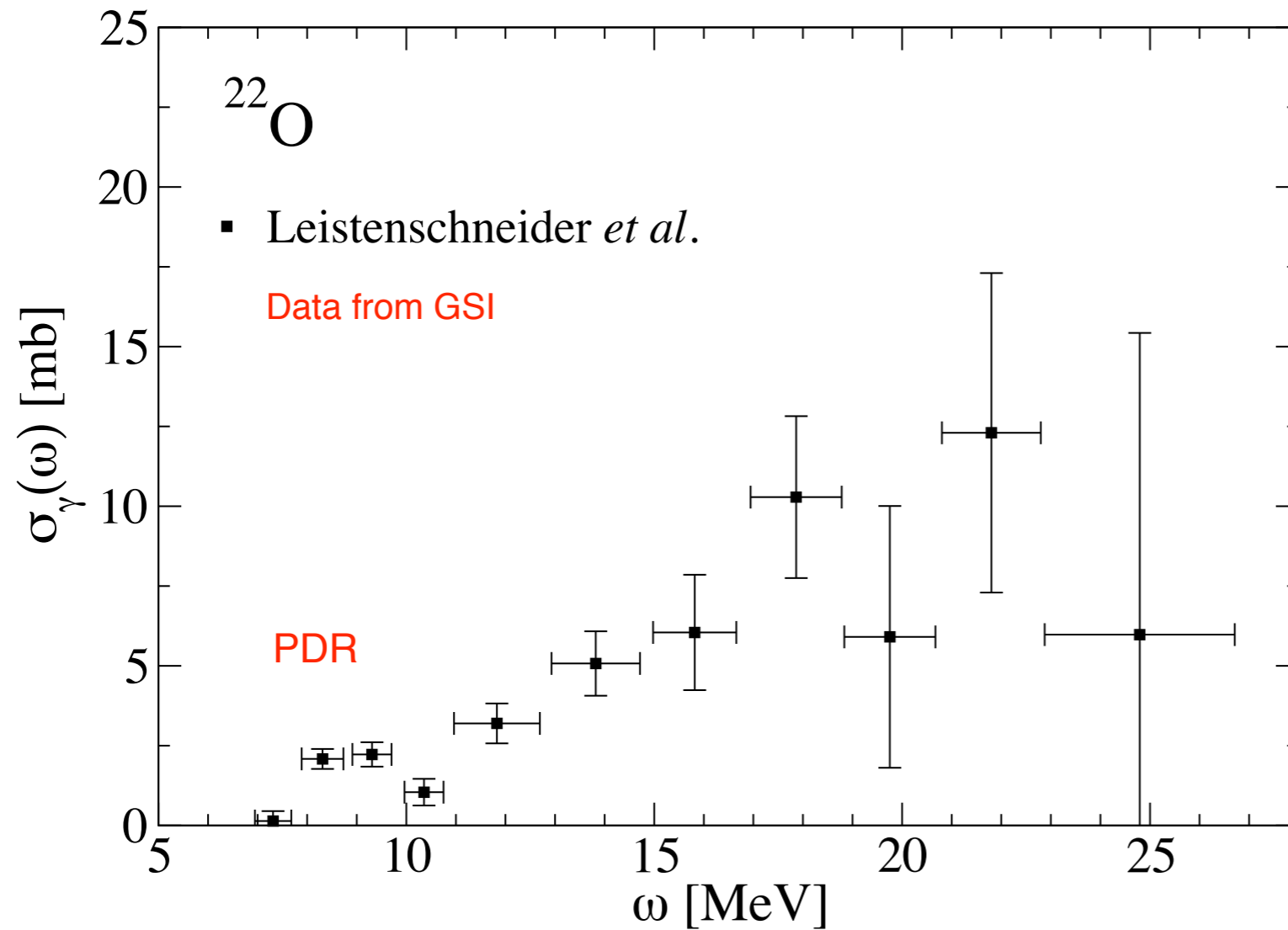
Validation ⁴He



Extension to ¹⁶O



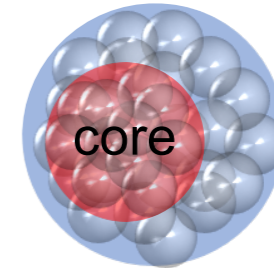
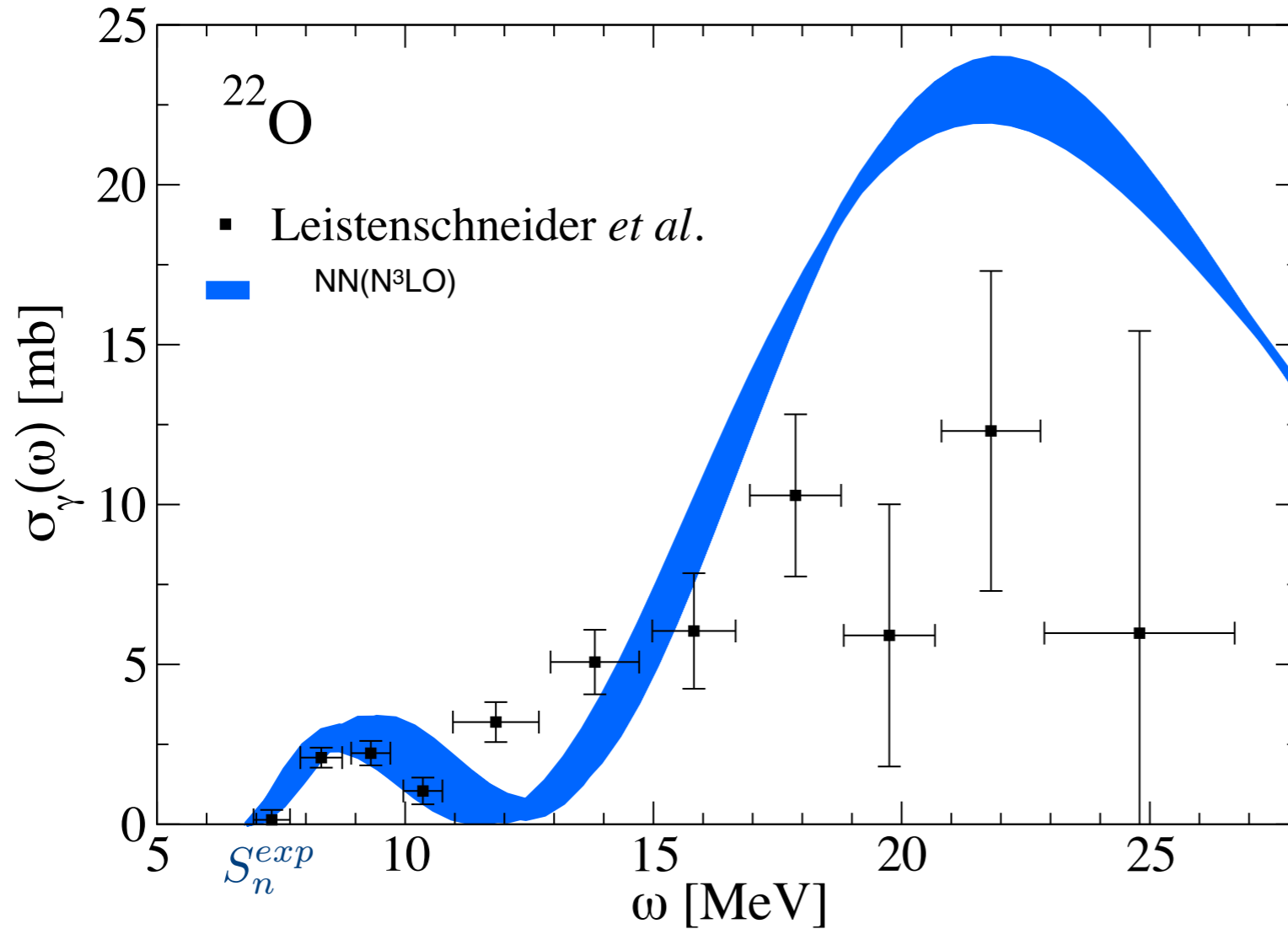
Photoexcitation of neutron-rich nuclei



Pigmy Dipole Resonance (PDR)

Photoexcitation of neutron-rich nuclei

S.B. *et al.*, PRC **90**, 064619 (2014)

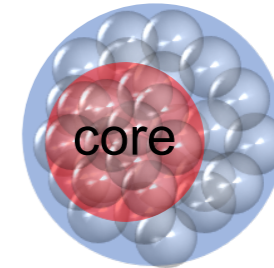
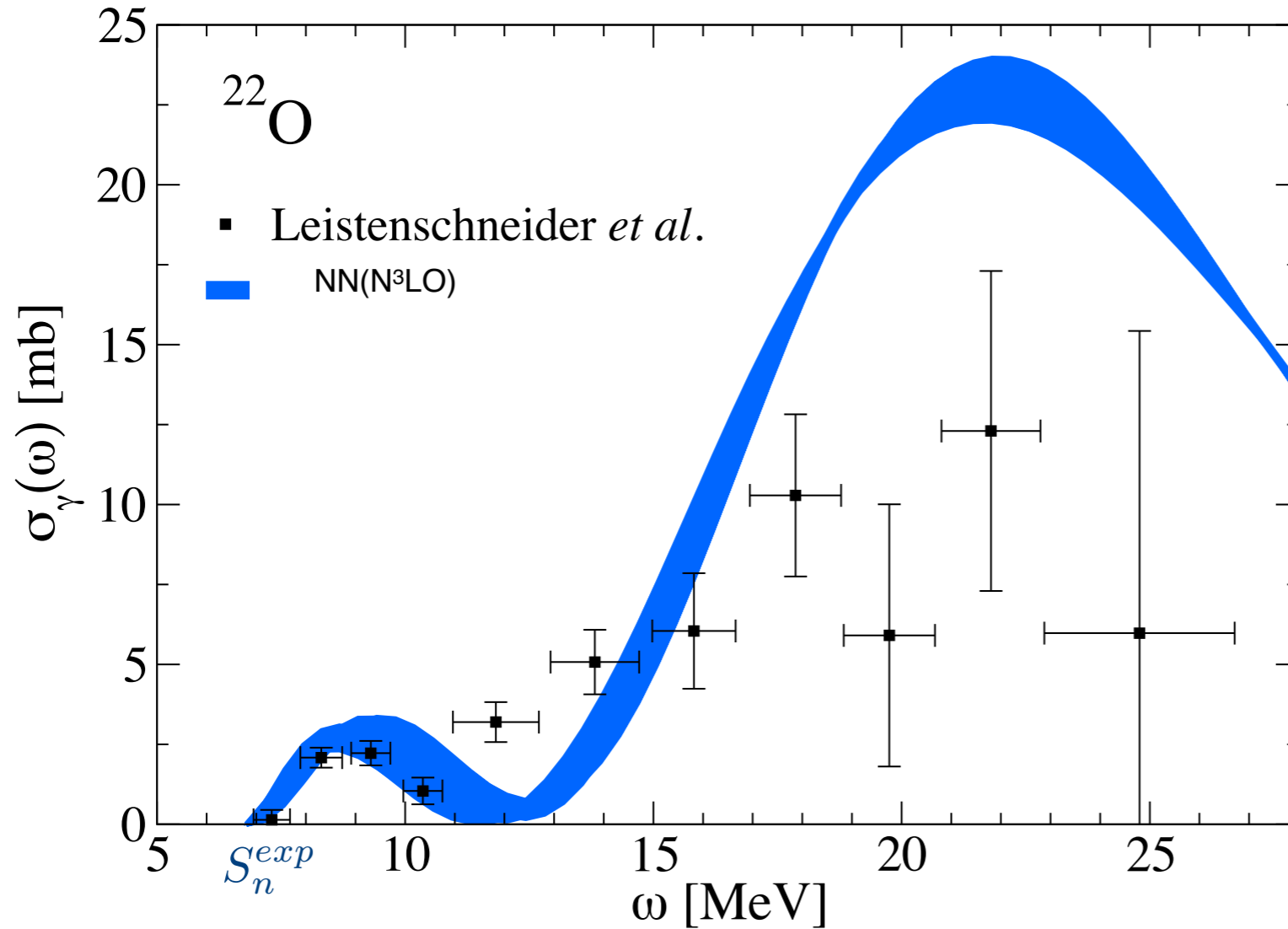


Pigmy Dipole Resonance (PDR)

Nicely described by a first principle calculation

Photoexcitation of neutron-rich nuclei

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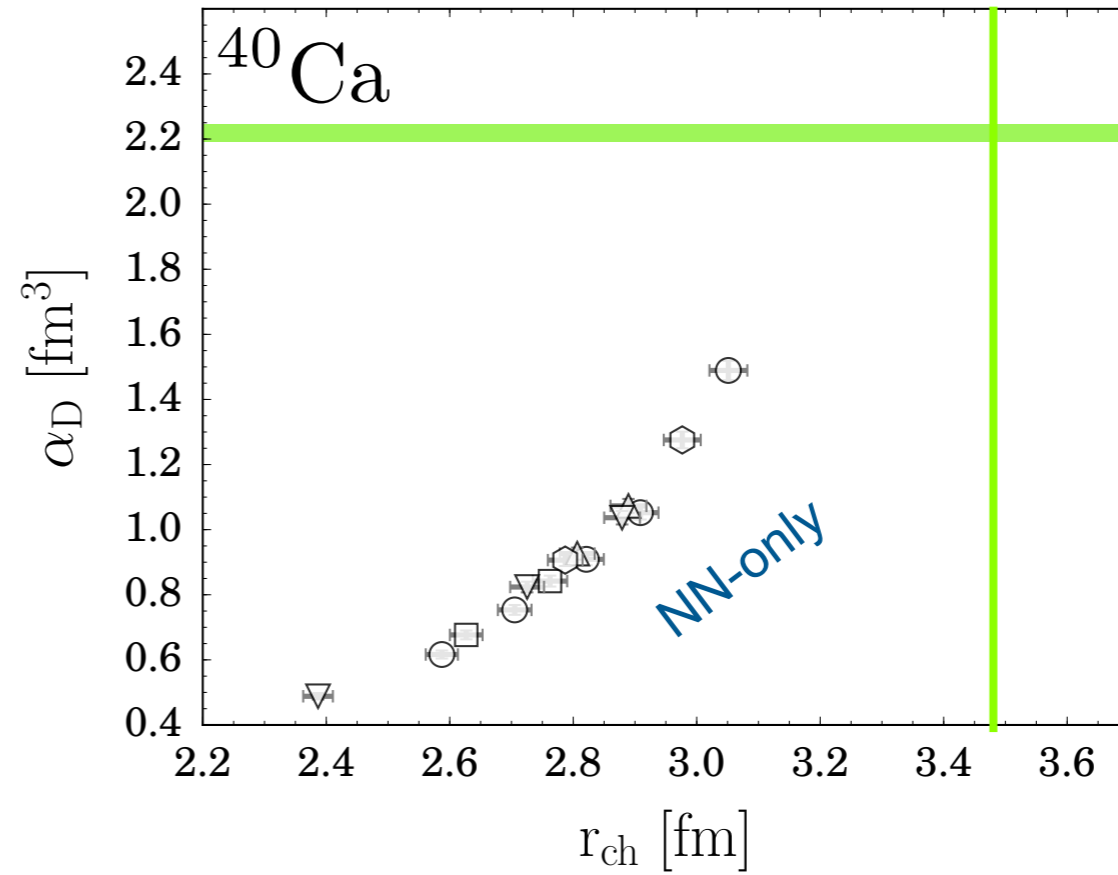
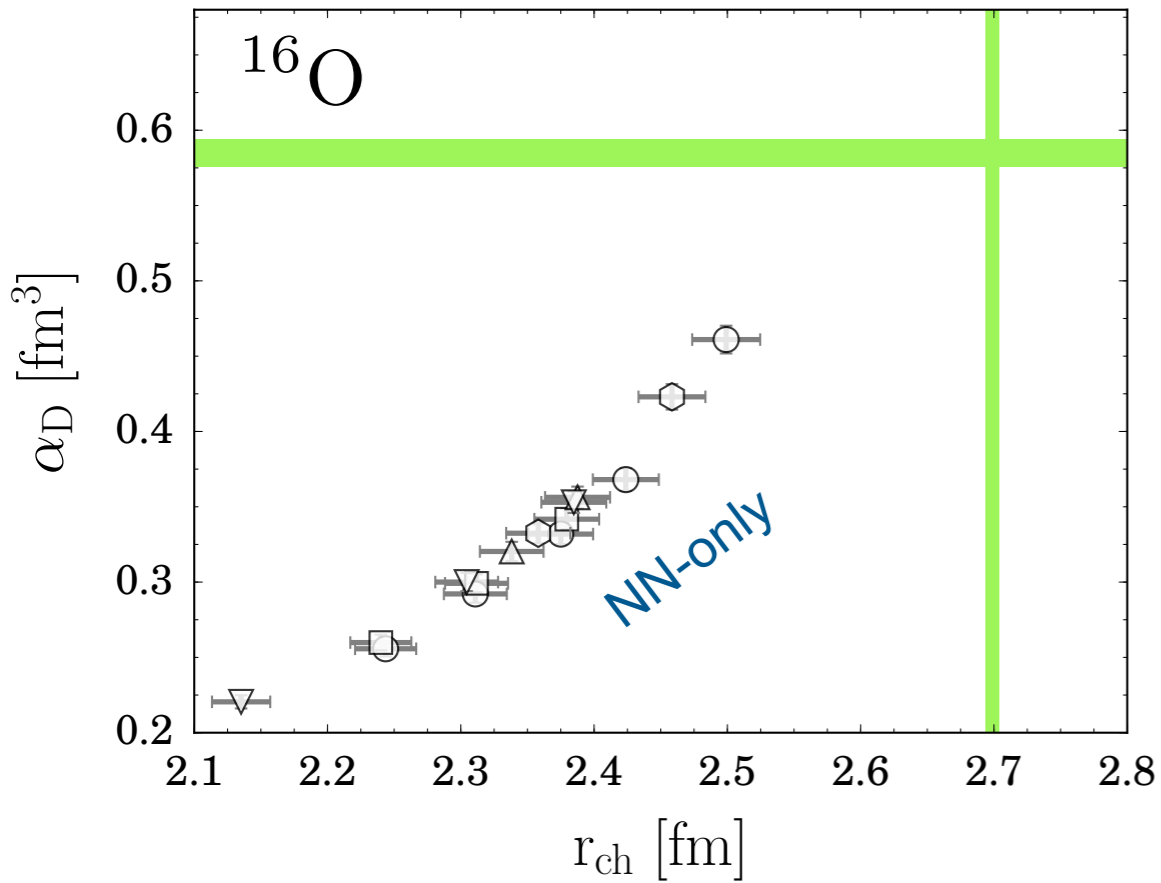
Theory provides a deeper understanding: microscopic interpretation of collective phenomena

Theory motivates new experiments: e.g. ^8He will be measured in RIKEN by T. Aumann

Electric Dipole Polarizability

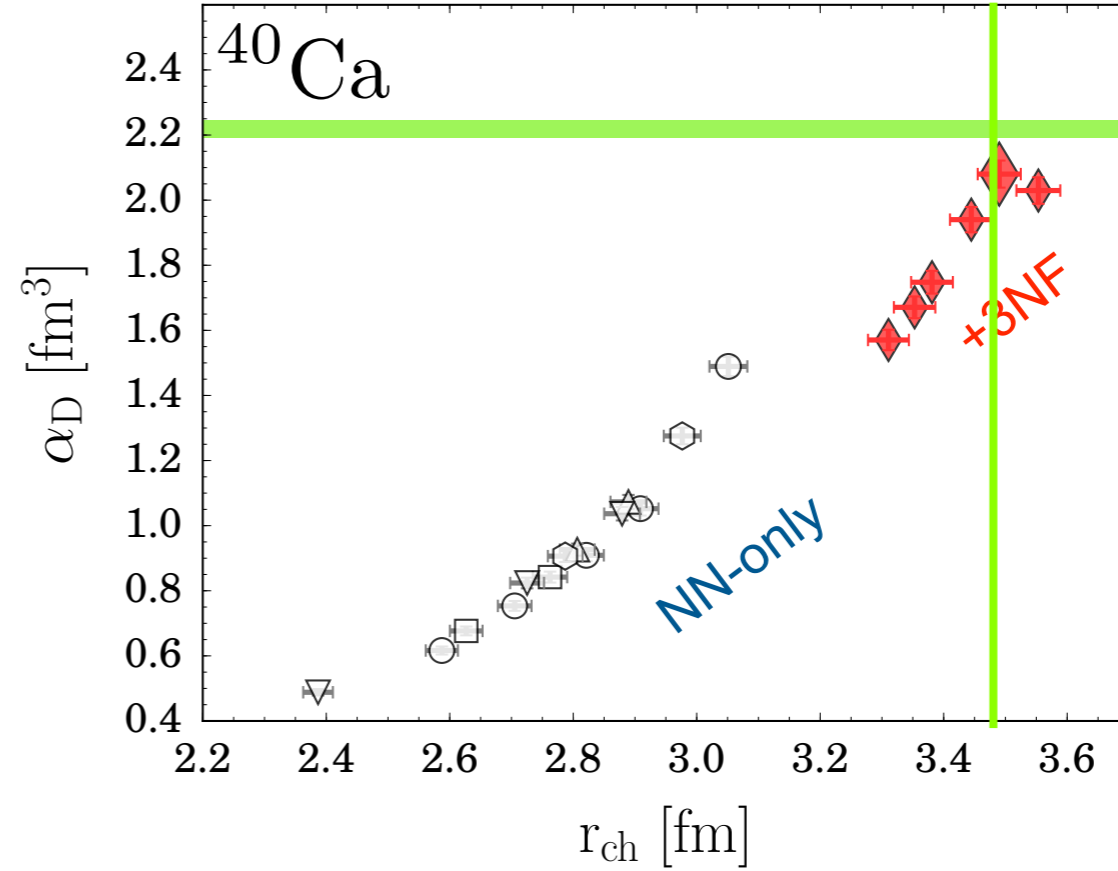
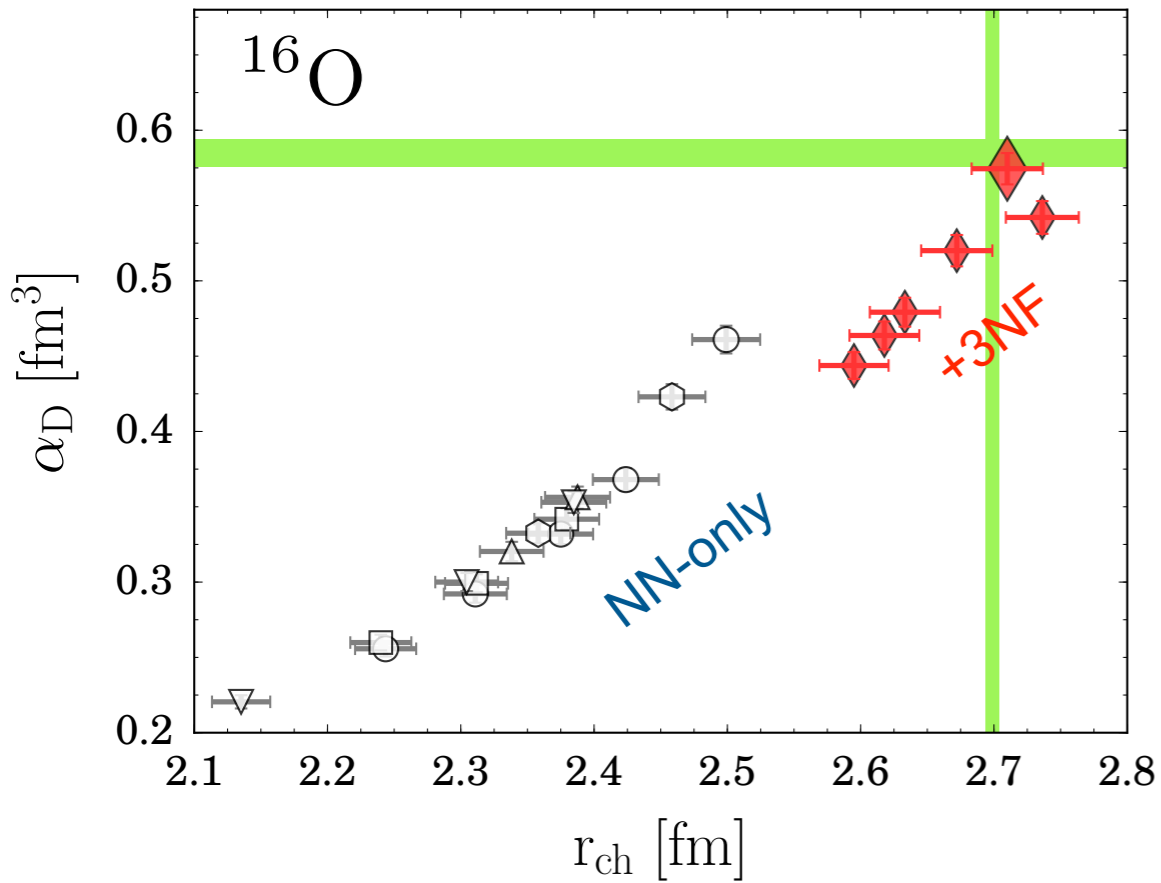
Medium-mass nuclei with NN + 3NF interactions

M. Miorelli *et al.*, PRC **94** 034317 (2016)



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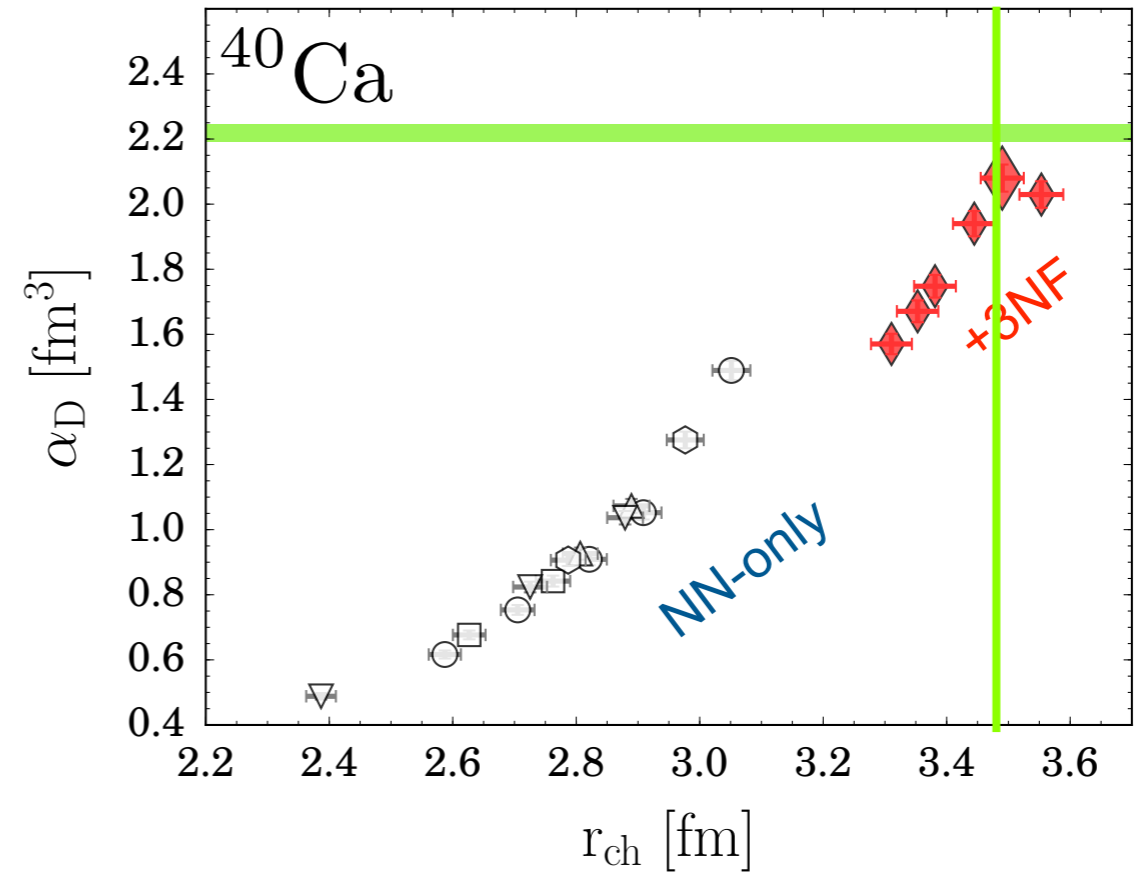
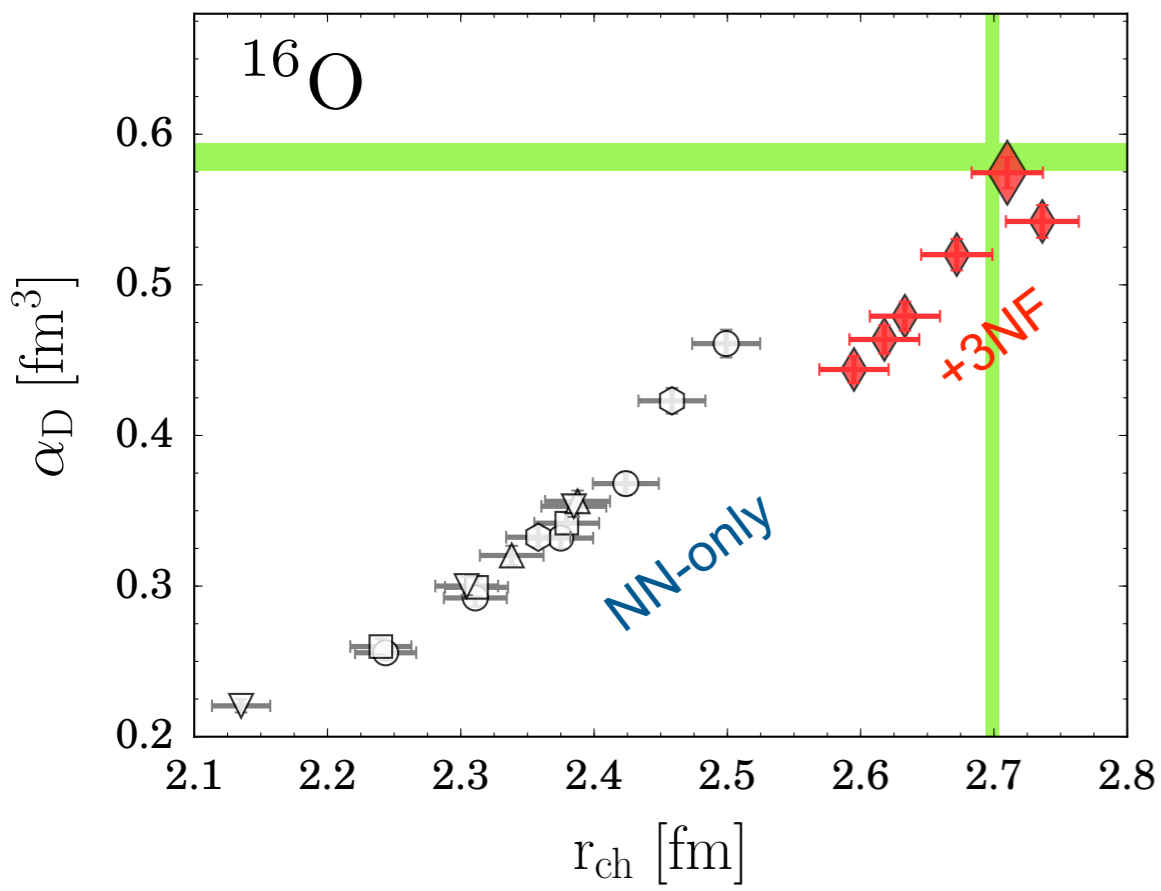


3NF A. Ekström *et al.*, Phys. Rev. C91, 051301 (2015)
 K. Hebeler *et al.*, Phys. Rev. C83, 031301 (2011)

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M. Miorelli *et al.*, PRC **94** 034317 (2016)

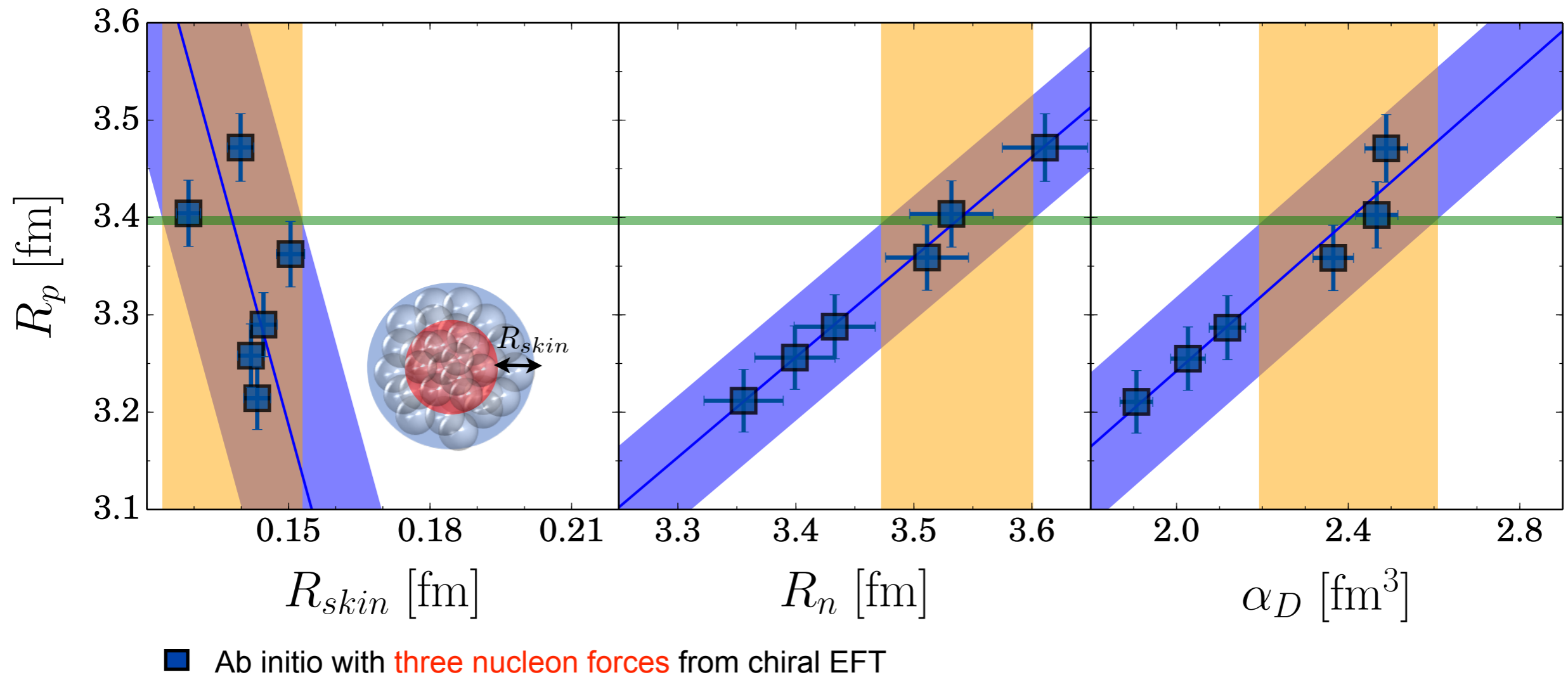


3NF A. Ekström *et al.*, Phys. Rev. C91, 051301 (2015)
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Much better agreement with experimental data
 Variation of Hamiltonian can be used to assess the theoretical error bar

^{48}Ca from first principles

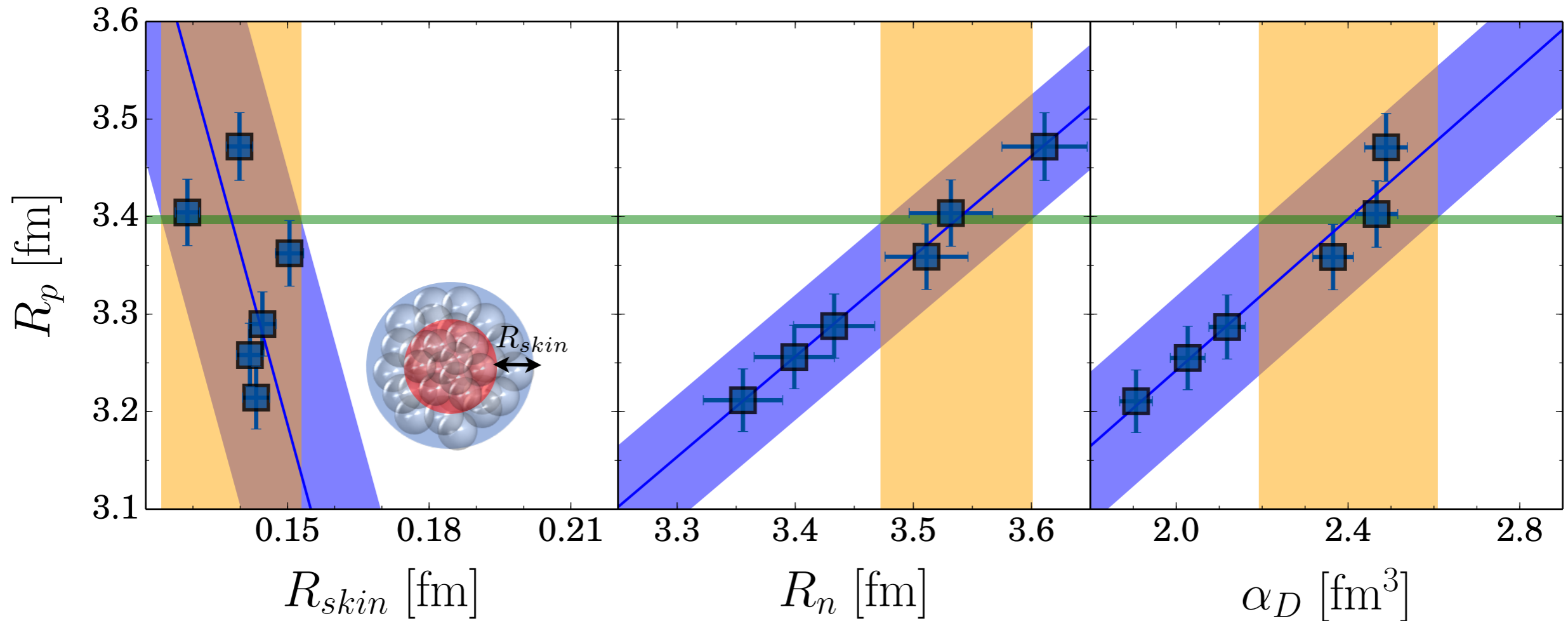
International collaboration (USA/Canada/Europe/Israel) using coupled-cluster theory
 Hagen *et al.*, *Nature Physics* **12**, 186 (2016)



^{48}Ca from first principles

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■ Ab initio with **three nucleon forces** from chiral EFT

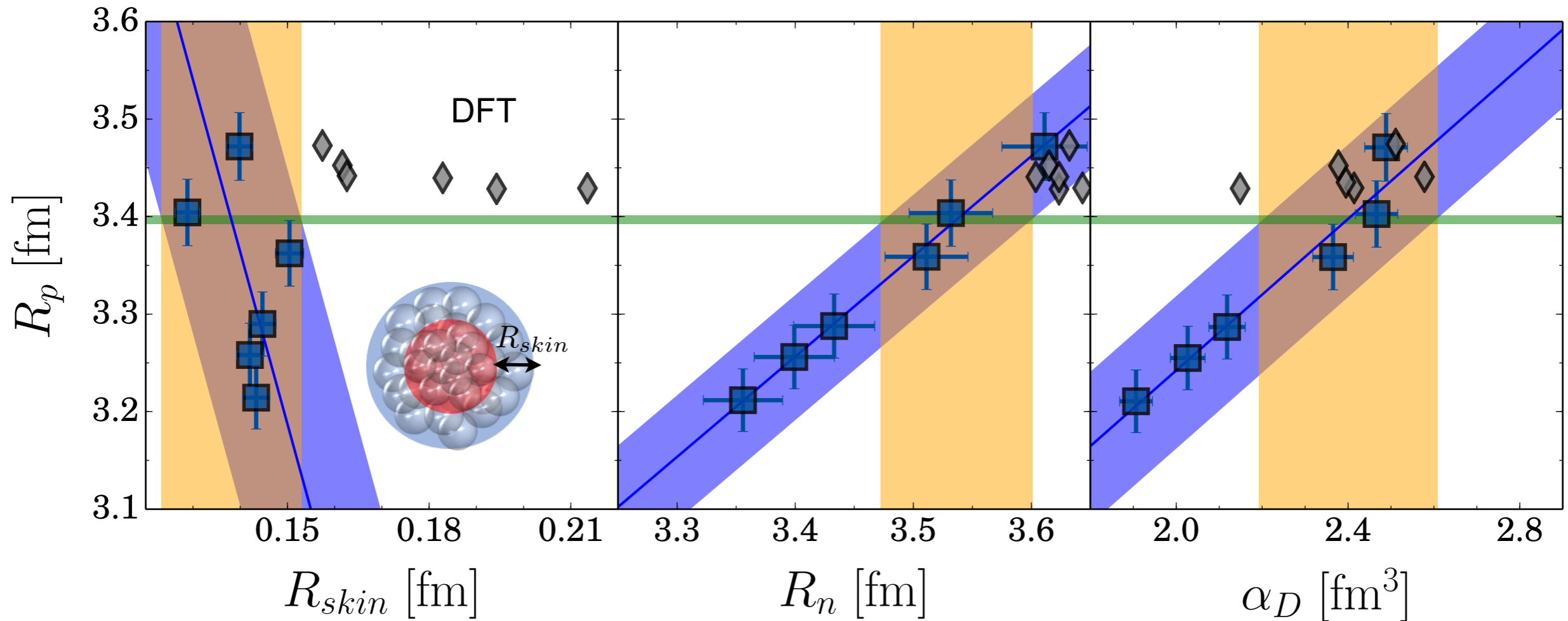
Strong correlations with R_p allow to put narrow constraints to R_{skin} and α_D

Ab-initio predictions: $0.12 \leq R_{skin} \leq 0.15$ fm
 $2.19 \leq \alpha_D \leq 2.60$ fm³

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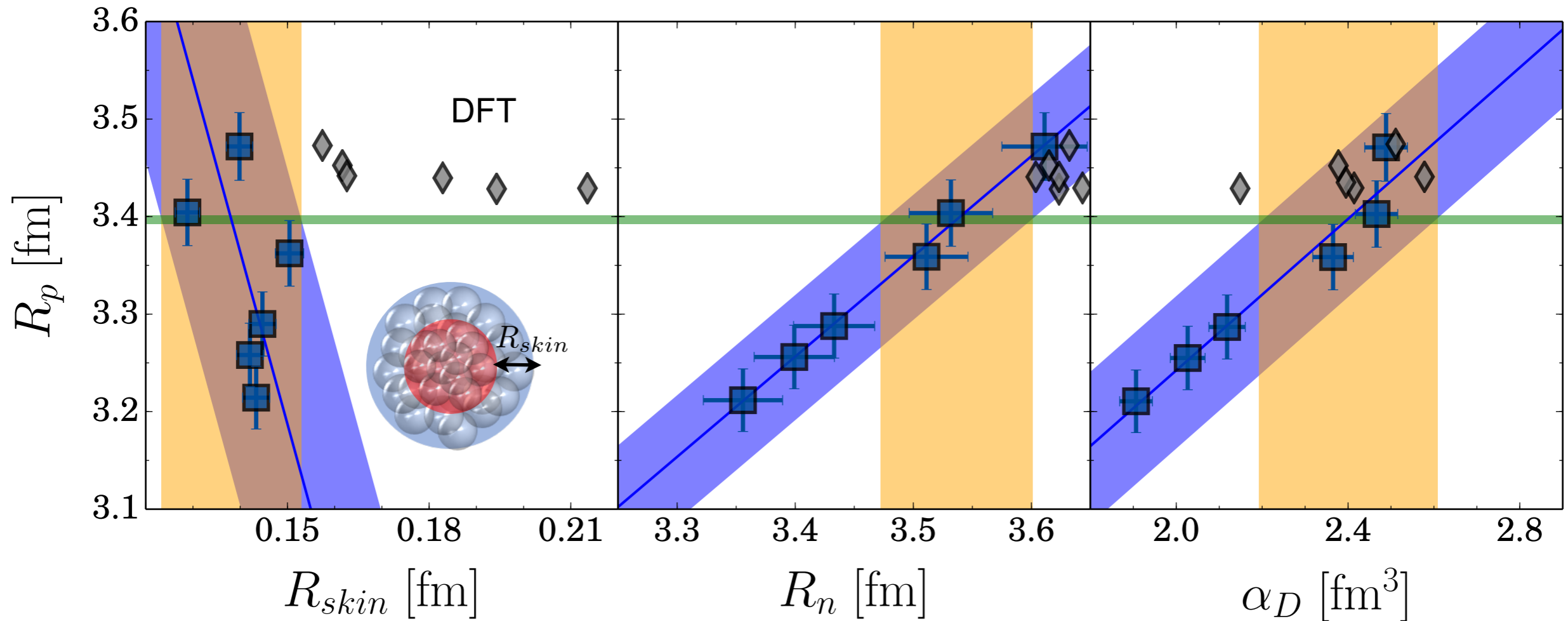
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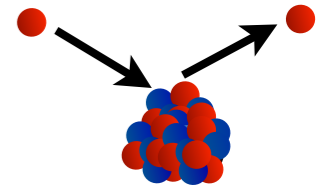
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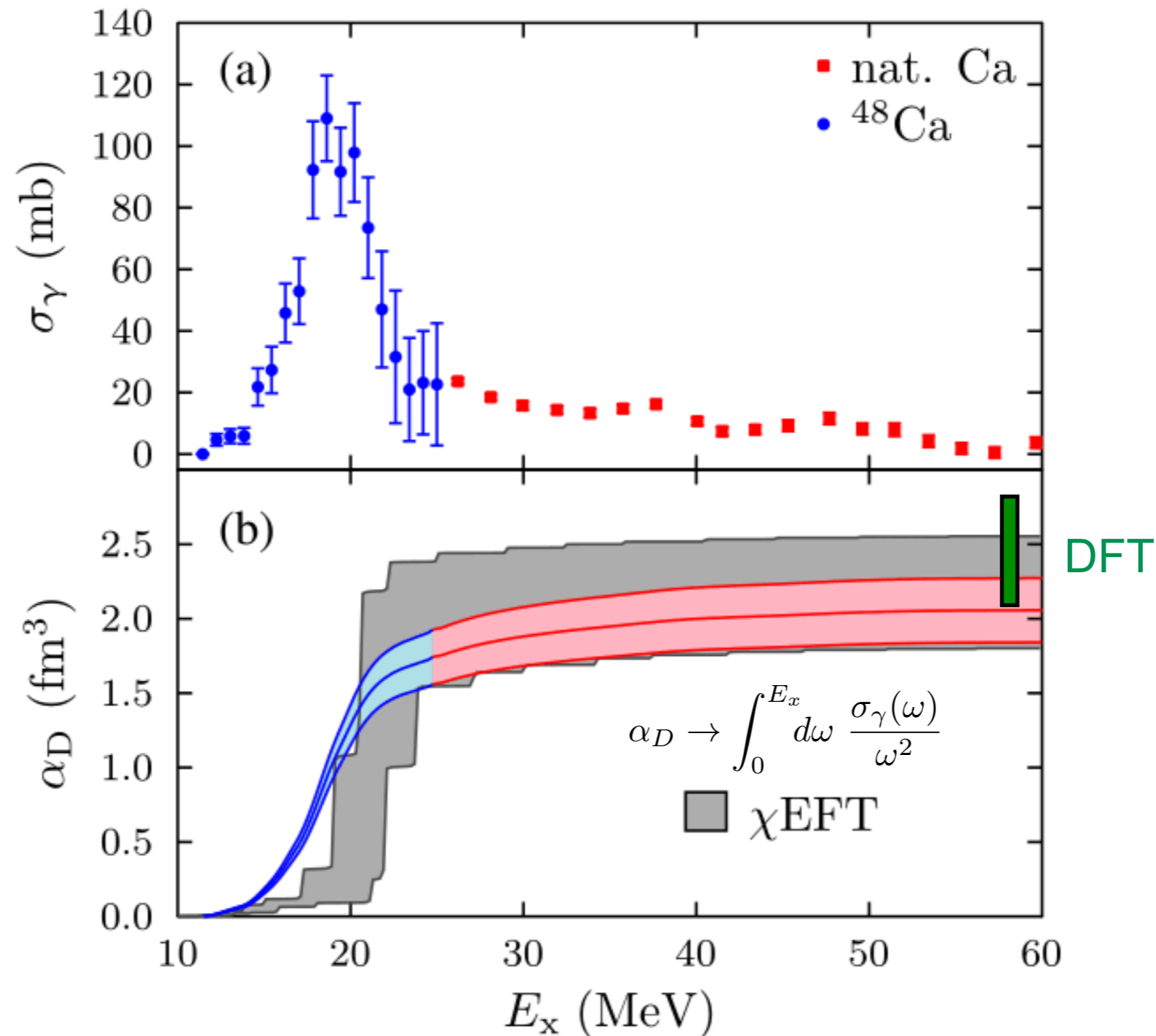
R_{skin} will be measured with **Parity violating electron scattering** CREX MREX

^{48}Ca electric dipole polarizability

New measurements from the Osaka-Darmstadt collaboration using inelastic proton scattering

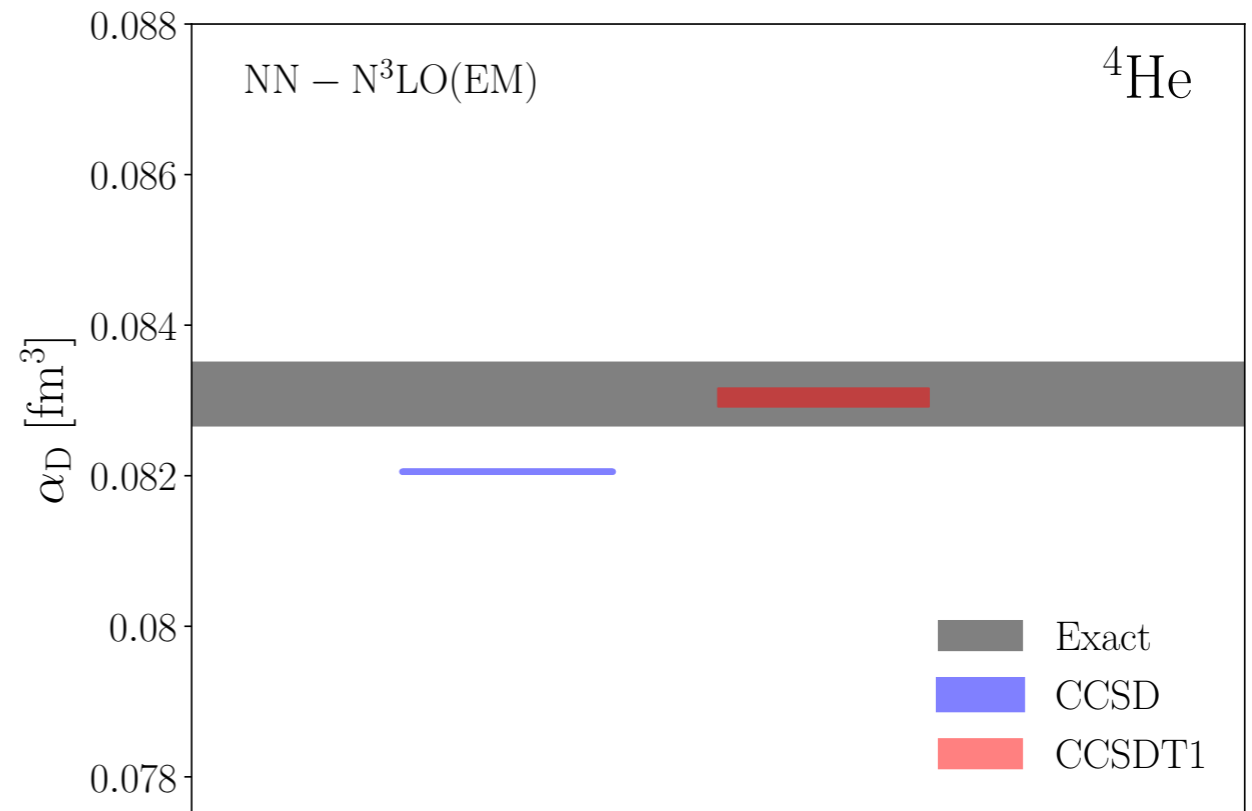
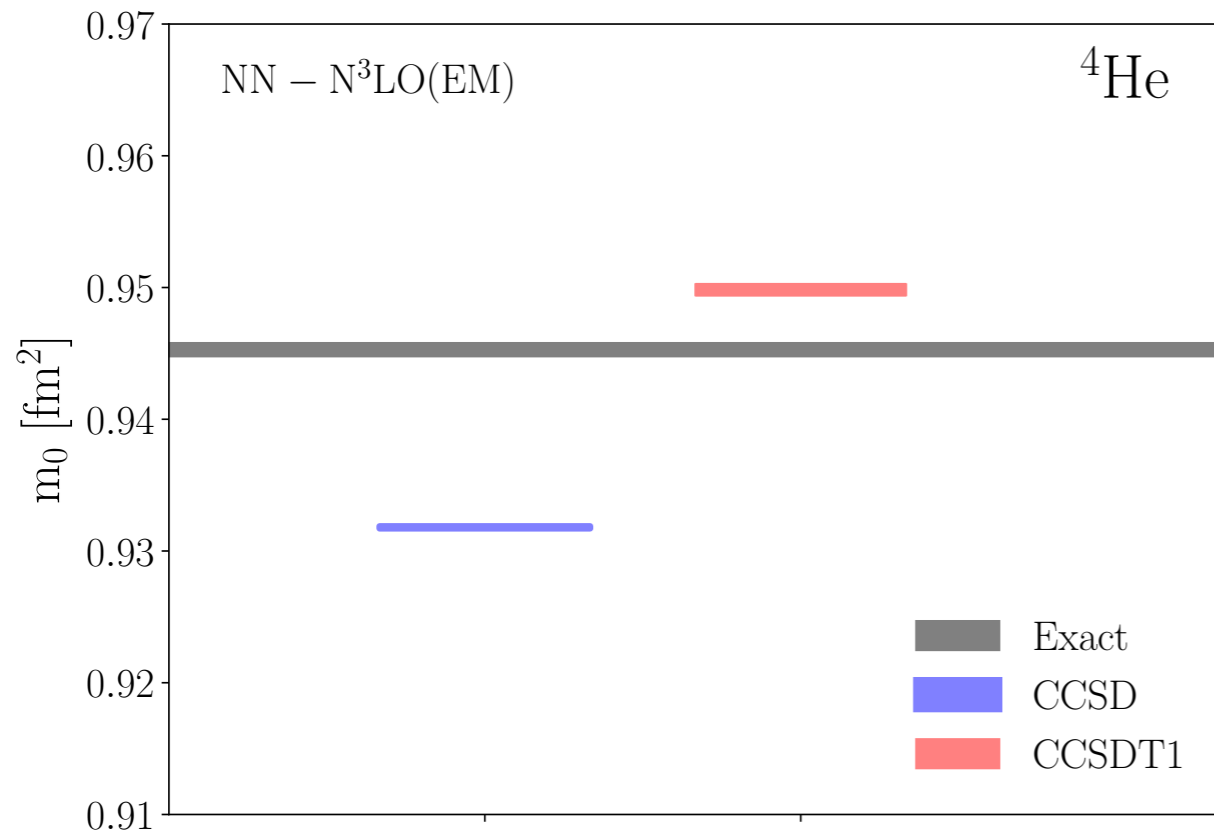


J.Birkhan, *et al.*, Phys. Rev. Lett. **118**, 252501 (2017)



How to improve our calculations

M. Miorelli *et al.*, in preparation (2017)



CCSD scheme $e^T = e^{T_1+T_2}$

$$R = R_0 + R_1 + R_2$$

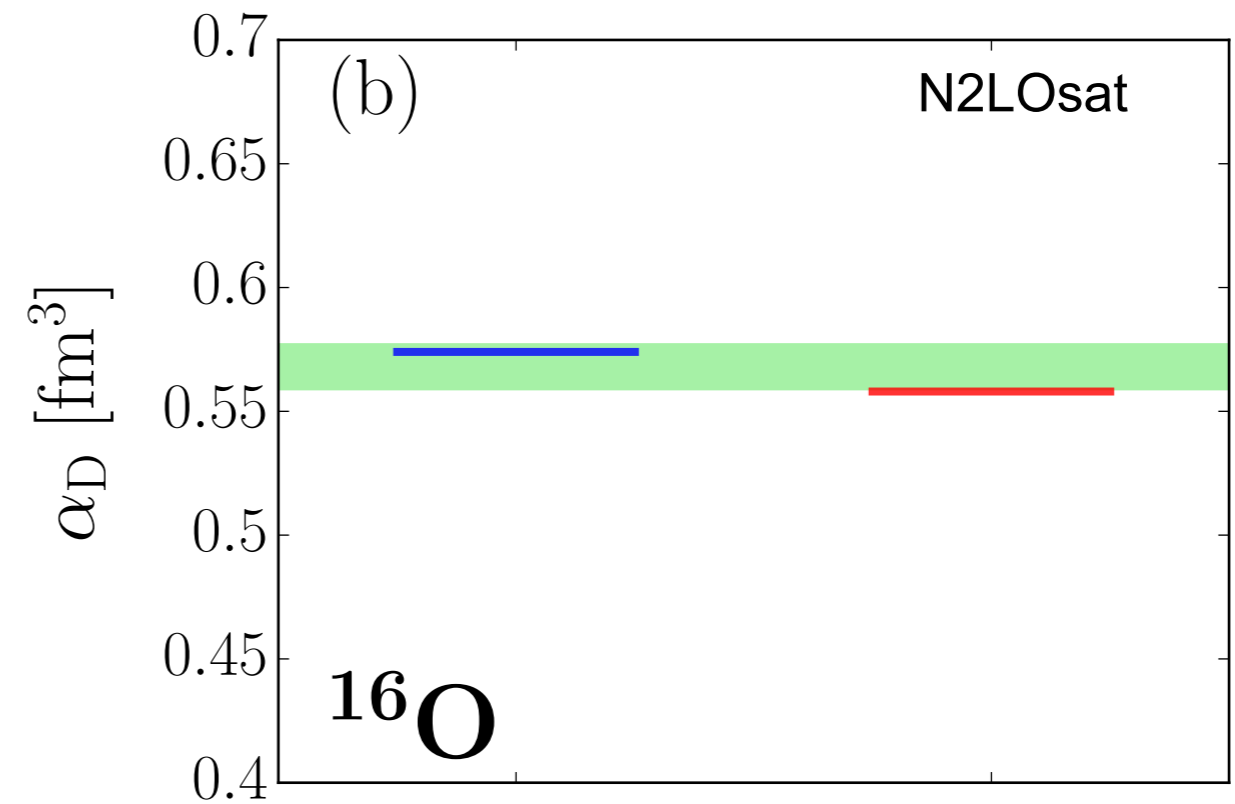
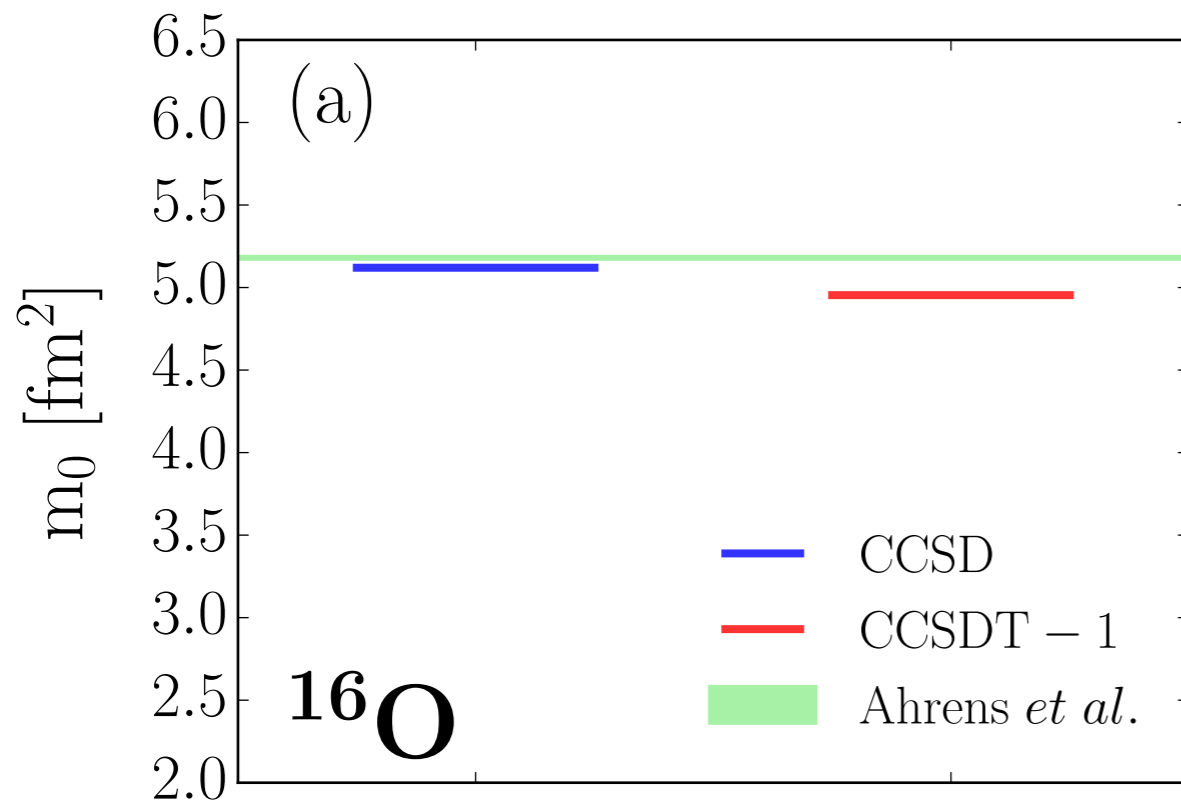
CCSDT1 scheme
(linearized triples) $e^T = e^{T_1+T_2} + T_3$

$$R = R_0 + R_1 + R_2 + R_3$$

Exact \Rightarrow hyperspherical harmonics, all correlations included (up to quadruples)

How to improve our calculations

M. Miorelli *et al.*, in preparation (2017)



CCSD scheme $e^T = e^{T_1+T_2}$

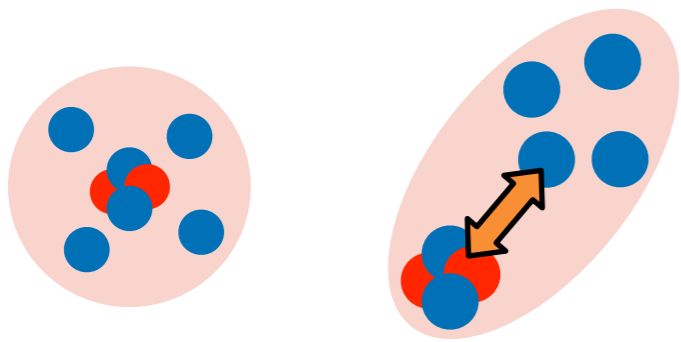
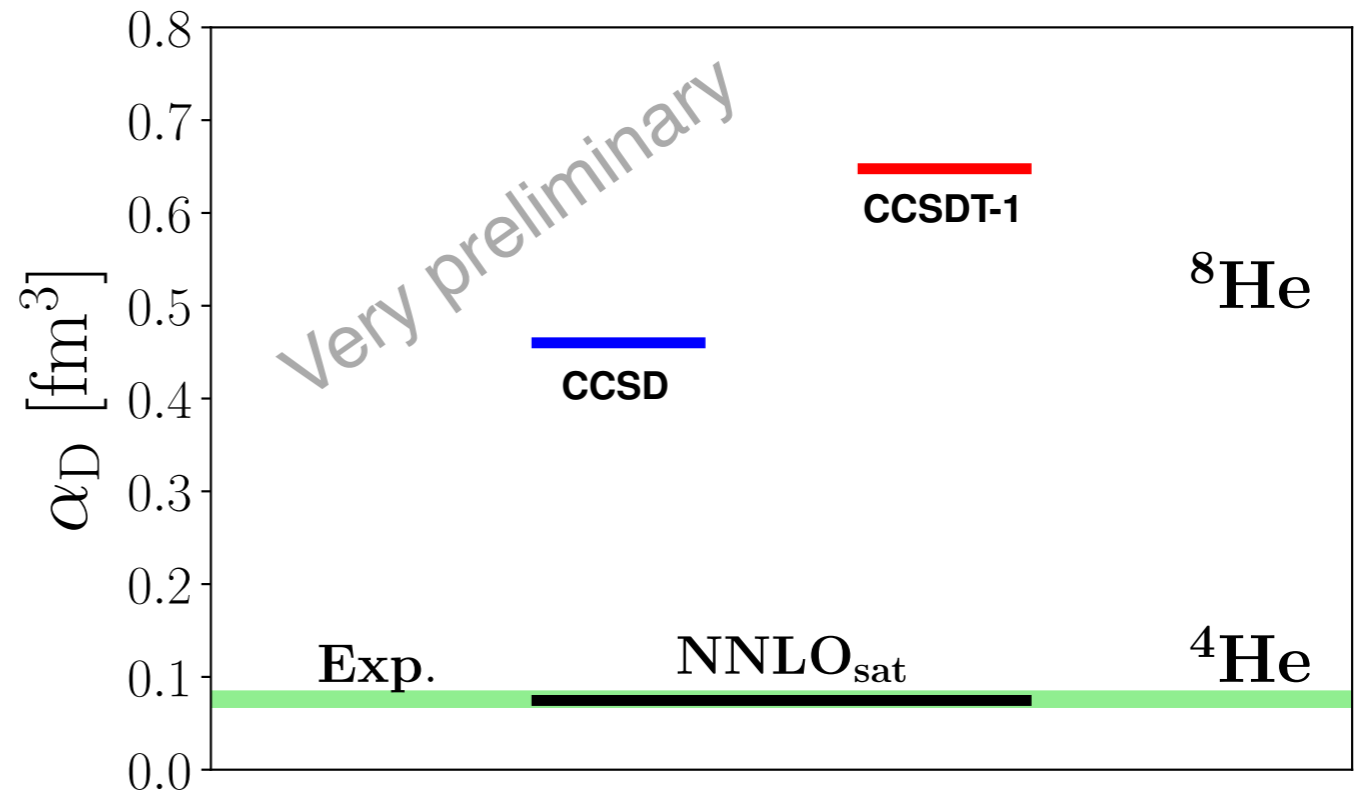
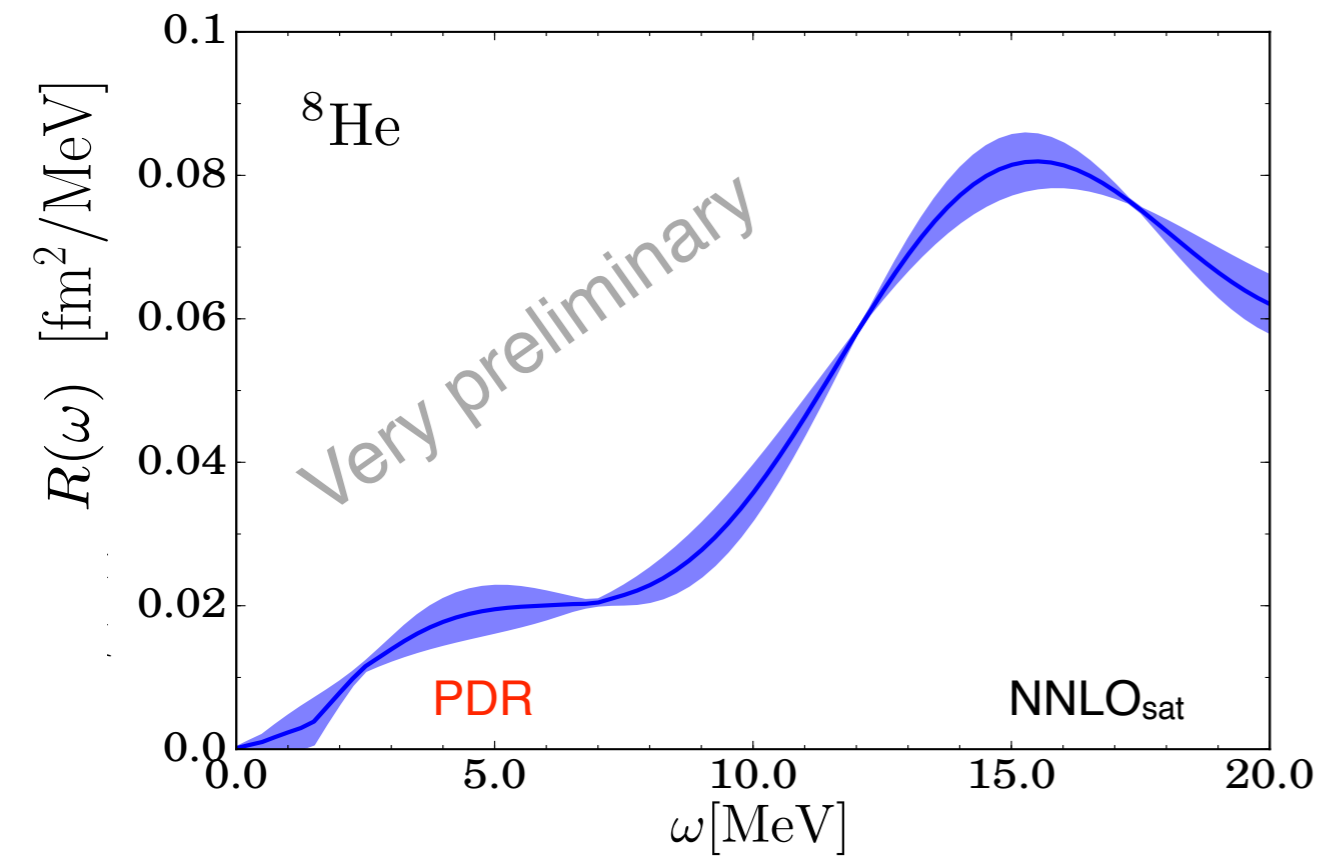
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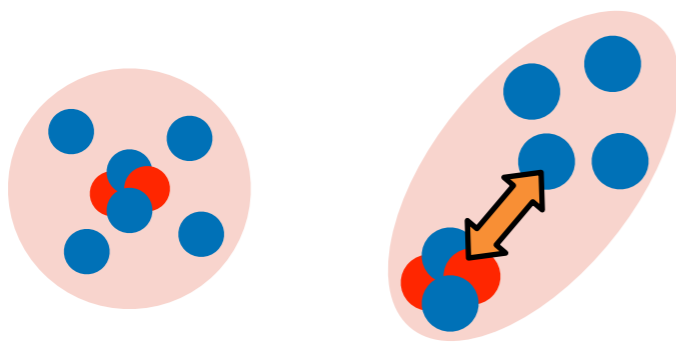
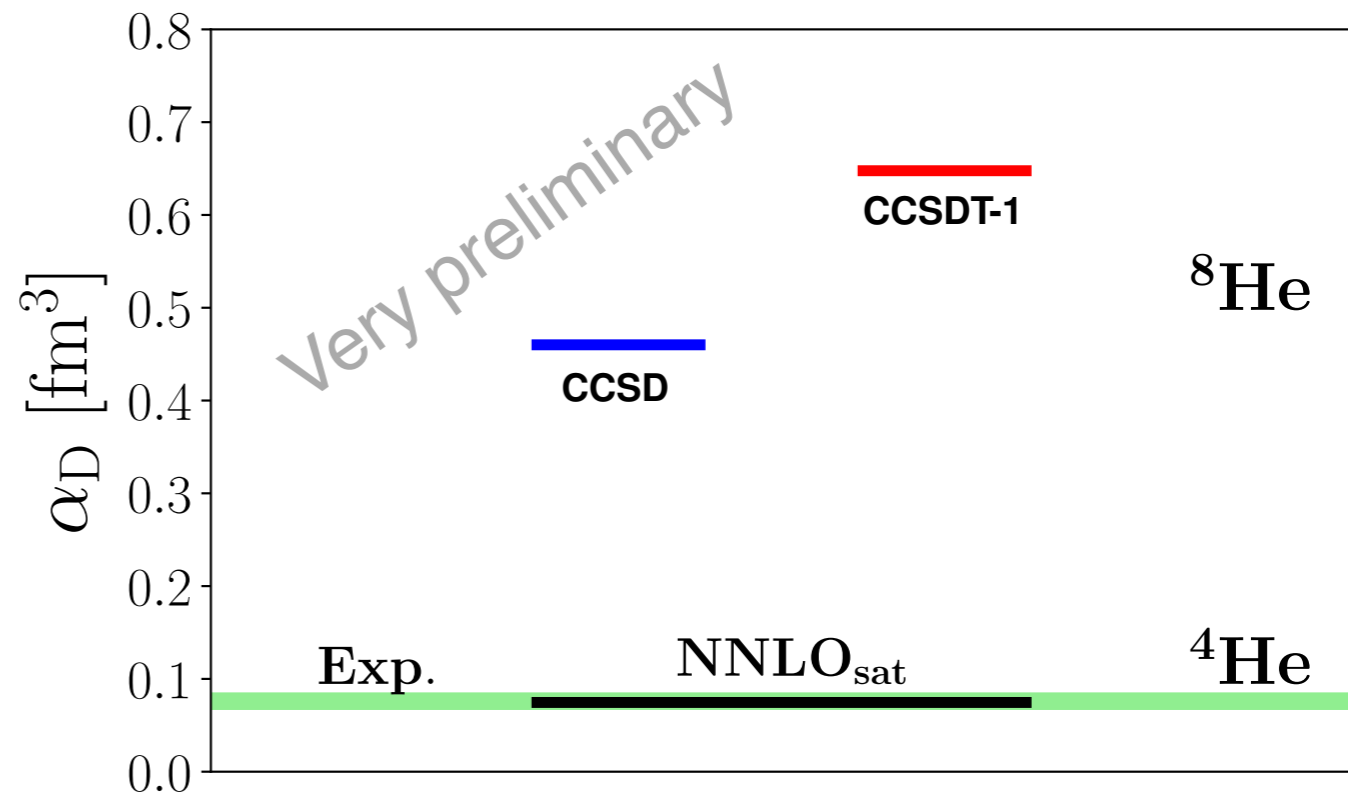
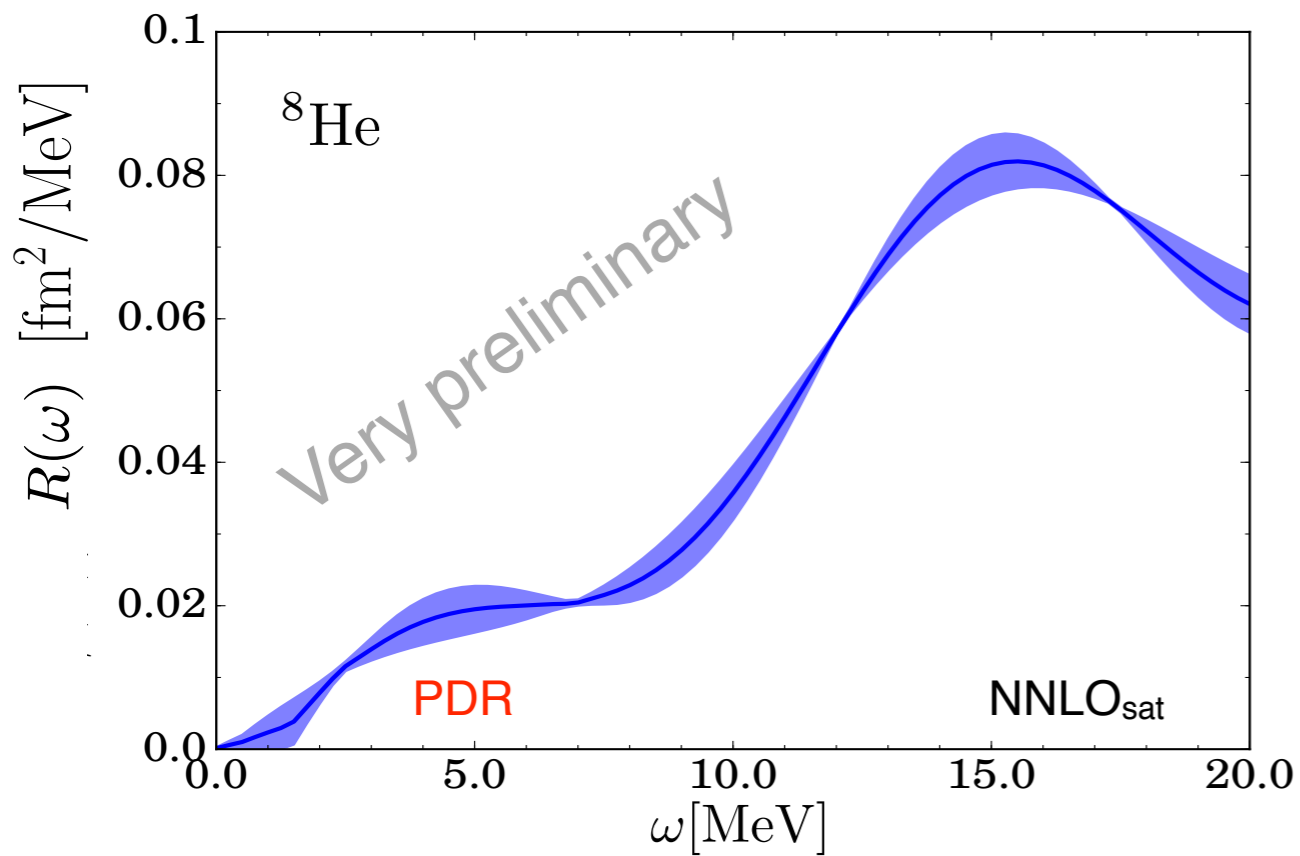
Applications to ^8He



$$\alpha_D = 2\alpha \int_{\omega_{th}}^{\infty} d\omega \frac{R(\omega)}{\omega}$$

$$\alpha_D(^8\text{He}) \gg \alpha_D(^4\text{He})$$

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Theory motivates new experiments
Will be measured by T. Aumann

- The ab-initio pathway is arguably the best way develop a strong predictive theory and connect to experiment
- Such connection can be exploited in several areas of physics
- In the future we will address electron-nucleus and neutrino-nucleus scattering

Thanks to all my collaborators

Thanks for your attention!

Connection to Neutron Stars

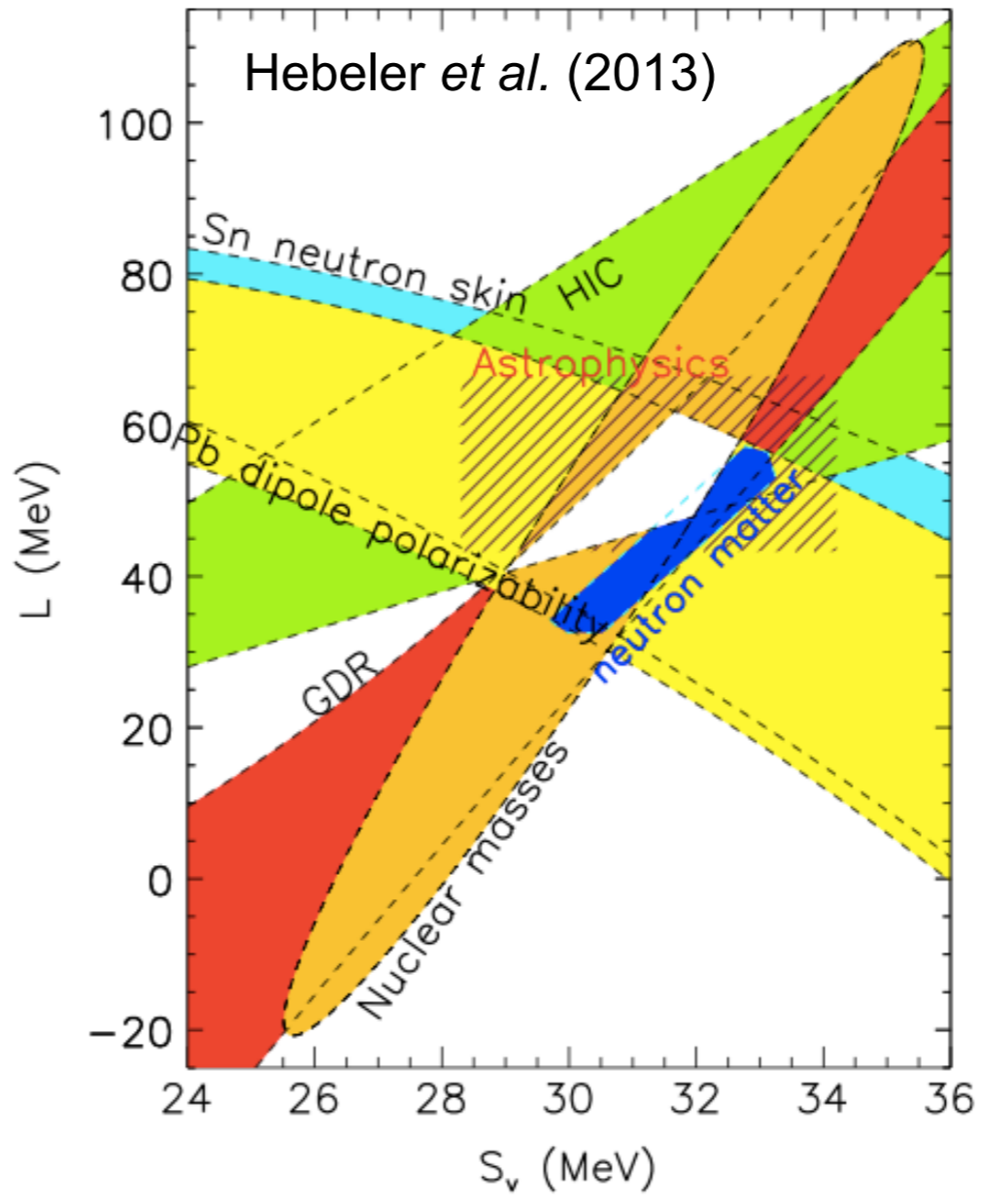
Equation of state

$$E(\rho, \delta) = E(\rho, 0) + S(\rho)\delta^2 + \mathcal{O}(\delta^4)$$

$$S(\rho) = S_v + \frac{L}{3\rho_0}(\rho - \rho_0) + \frac{K_{sym}}{18\rho_0^2}(\rho - \rho_0)^2 + \dots$$

$$\rho = \rho_n + \rho_p, \quad \delta = \frac{\rho_n - \rho_p}{\rho_n + \rho_p}$$

S_v and L can be inferred from heavy ion collisions and are correlated with finite nuclei observables



Connection to Neutron Stars

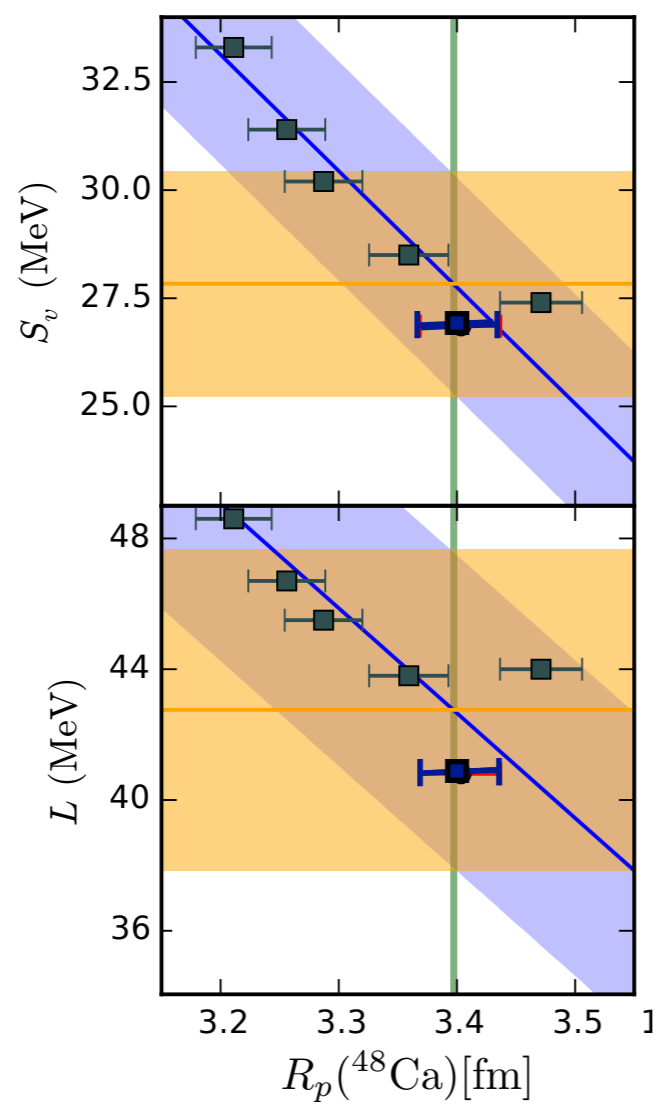
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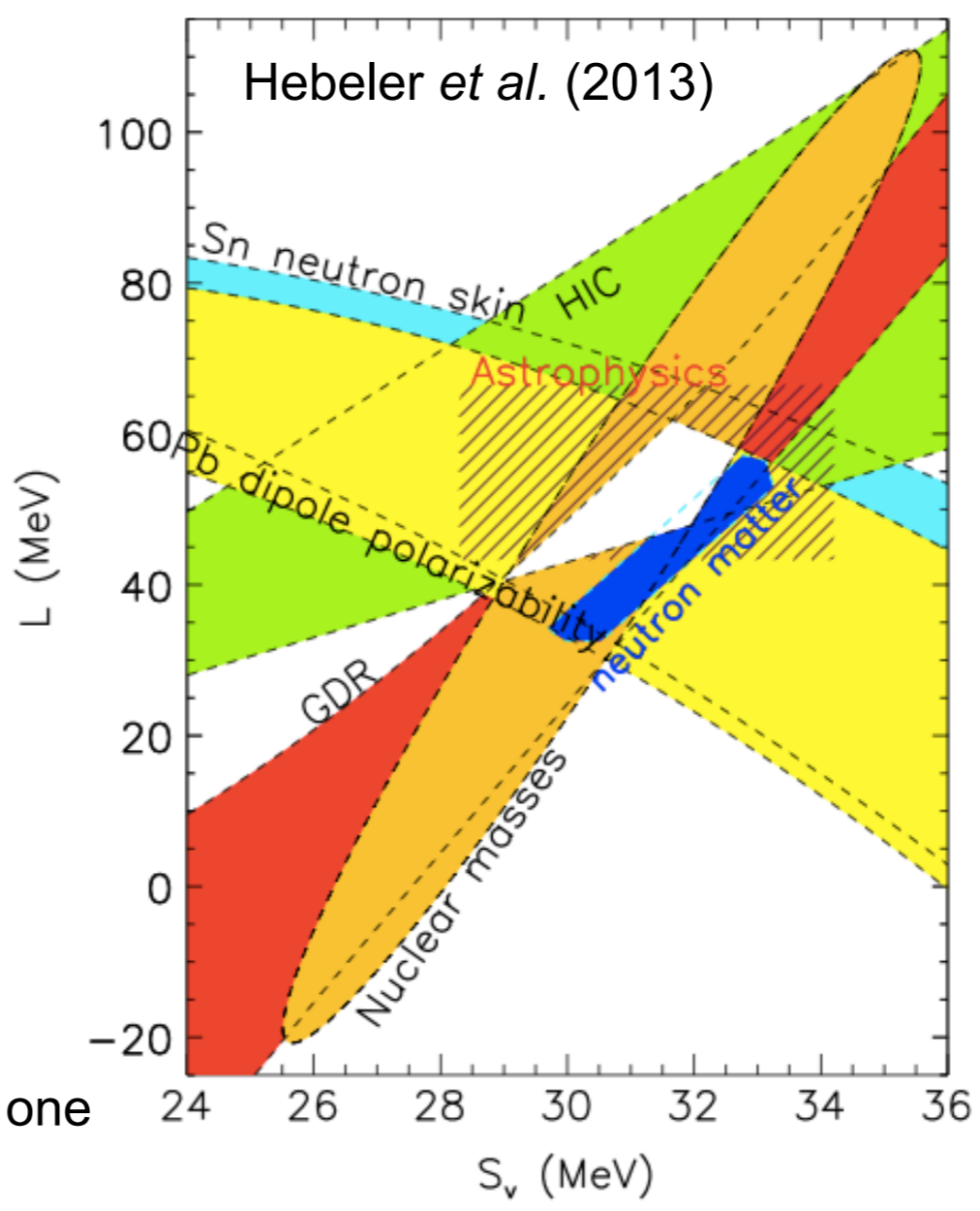
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Using correlations with R_p in ^{48}Ca one can constrain nuclear matter



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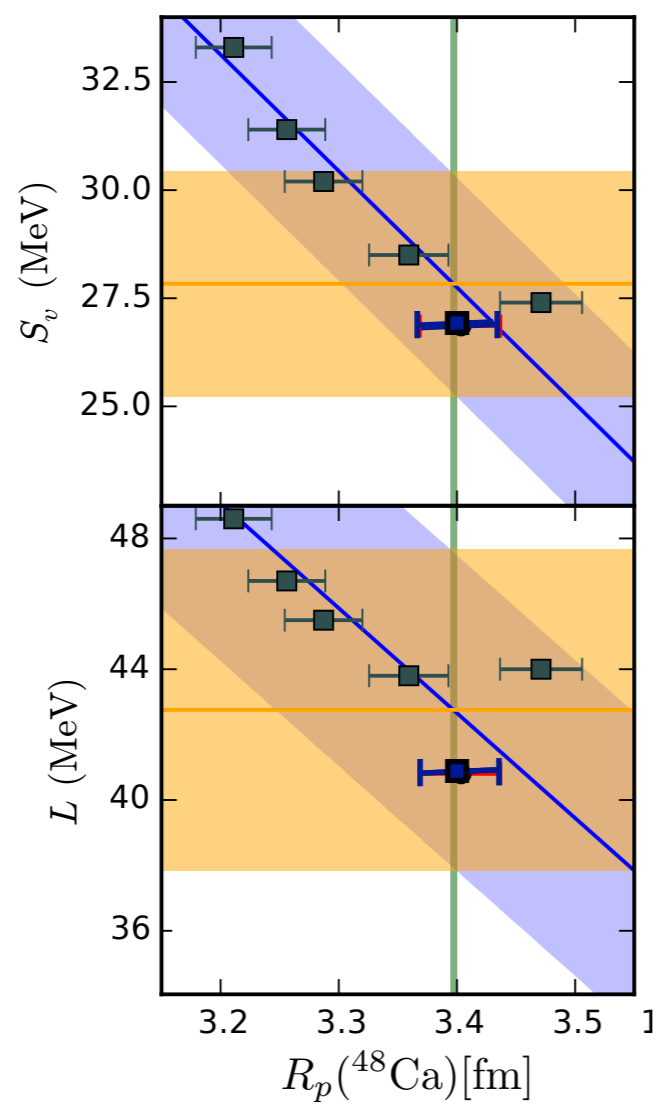
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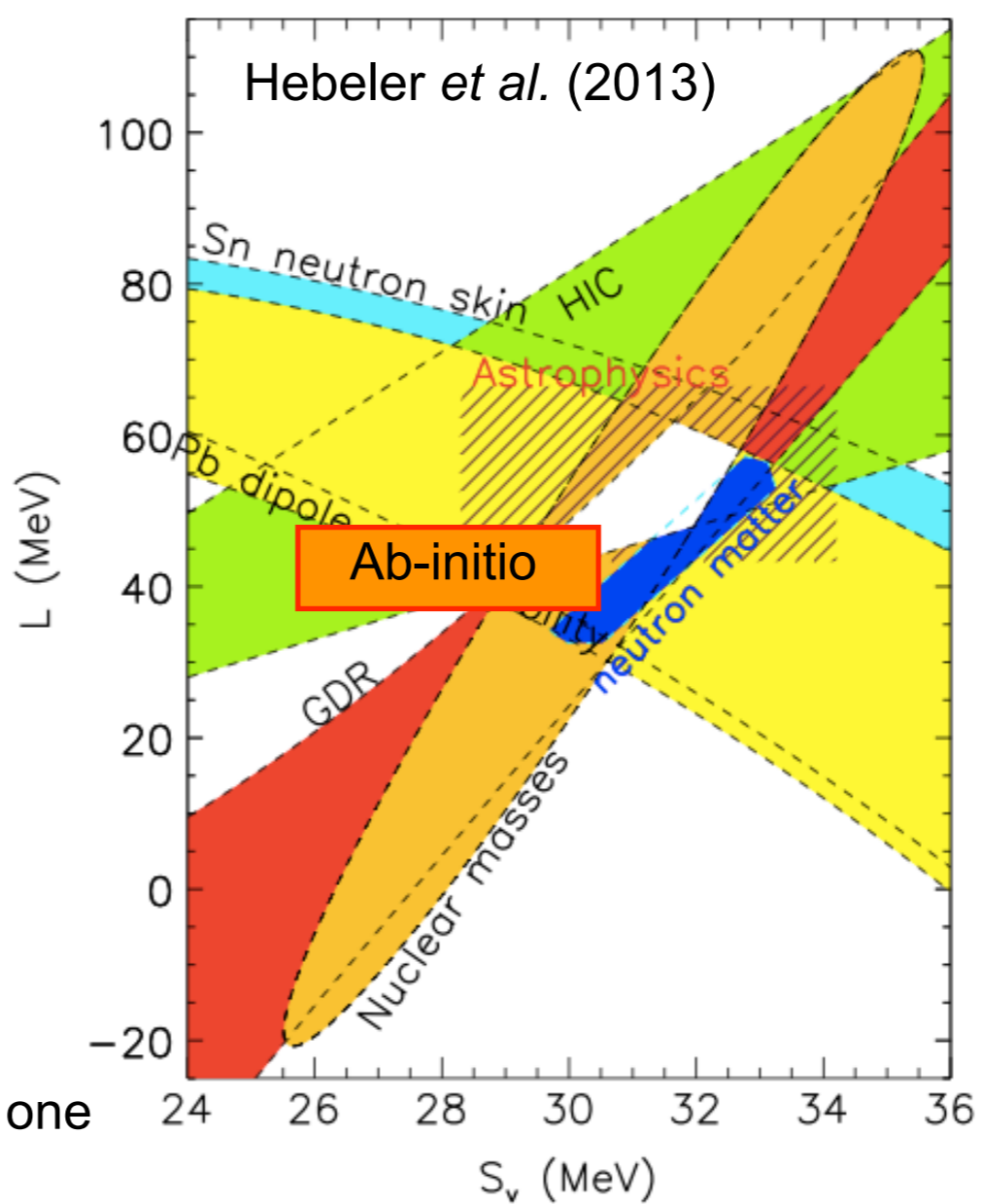
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Using correlations with R_p in ^{48}Ca one can constrain nuclear matter



$25.2 \leq S_v \leq 30.4$ MeV
 $37.8 \leq L \leq 47.7$ MeV