

# Towards the Equation of State of Neutron Stars

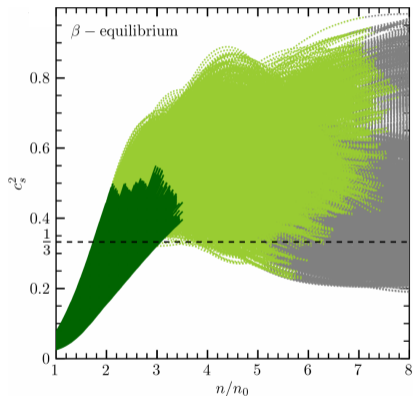
Andreas Geißel

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Technische Universität Darmstadt

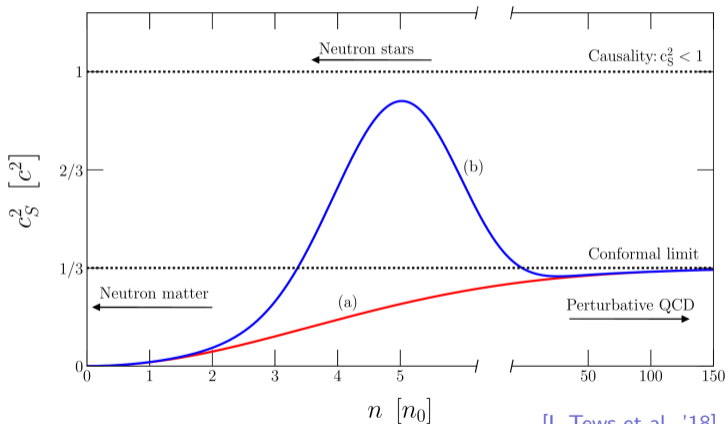
October 6, 2022



# Speed of Sound



[S. Huth, C. Wellenhofer, and  
A. Schwenk '20]

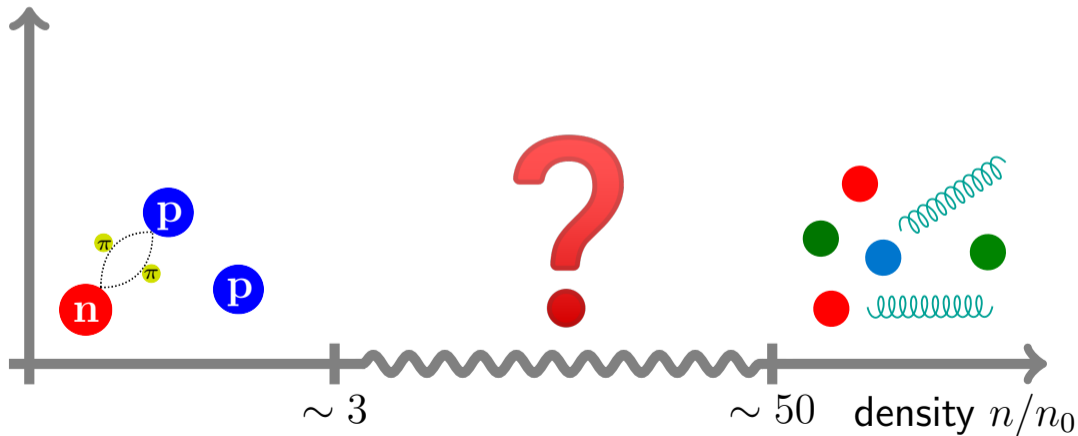


[I. Tews et al. '18]

# Degrees of Freedom



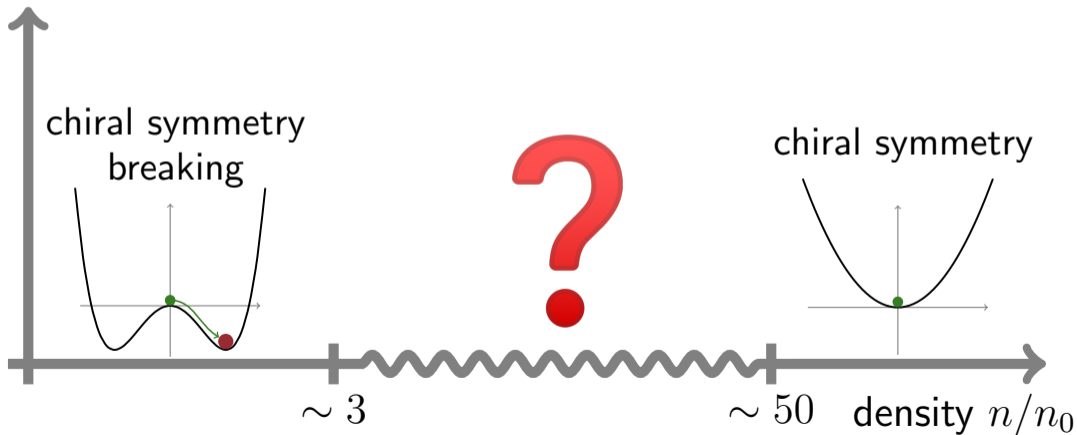
temperature  $T$



# Degrees of Freedom



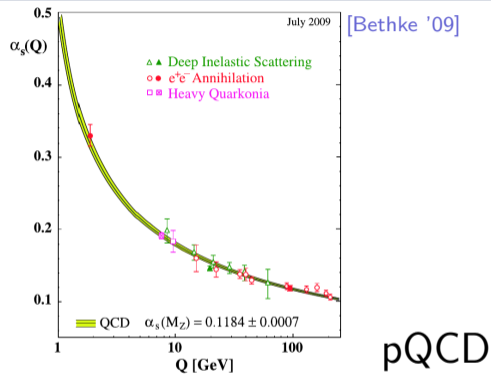
temperature  $T$



# Degrees of Freedom



temperature  $T$



$\chi$ EFT

[K. Hebeler '20;  
E. Epelbaum, H.W. Hammer,  
and Ulf-G. Meißner '08; ...]

Functional  
Renormalization Group

pQCD

[A. Kurkela et al. '14;  
T. Gorda et al. '18; ...]

$\sim 3$

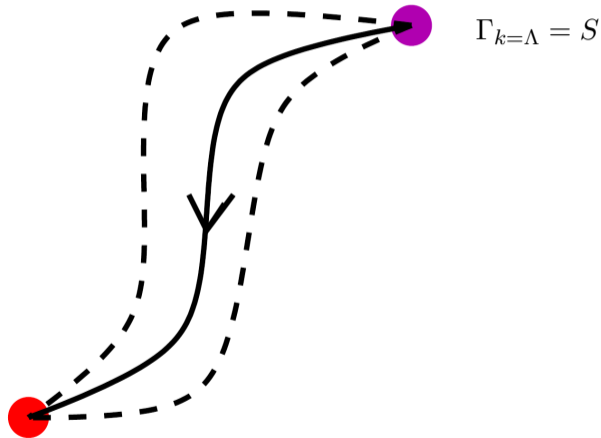
$\sim 50$

density  $n/n_0$

# Functional Renormalization Group



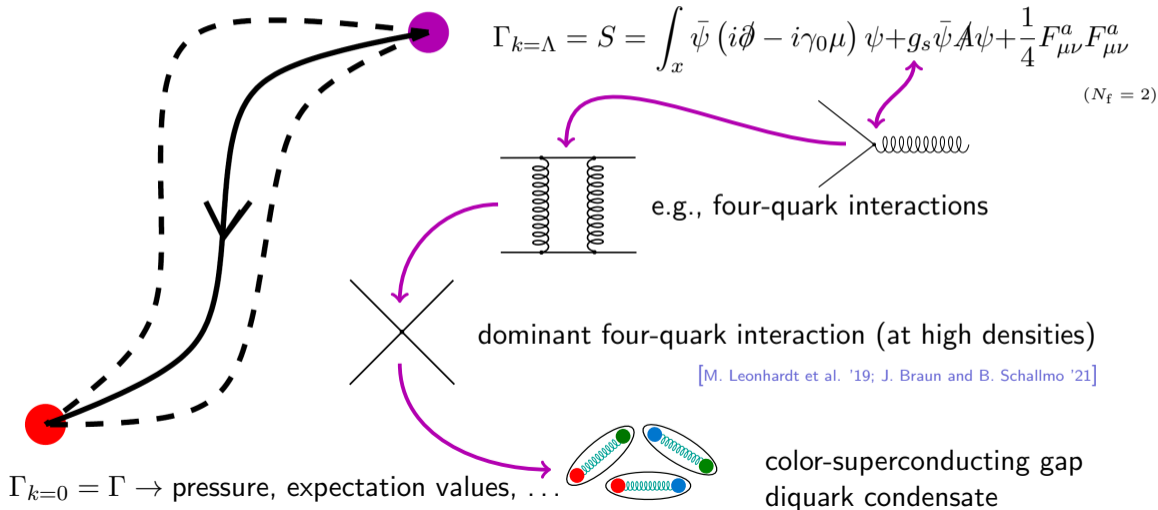
[Wetterich '93; (figure taken from) Gies '06]



$\Gamma_{k=0} = \Gamma \rightarrow$  pressure, expectation values, ...

# Functional Renormalization Group

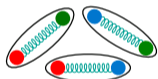
[Wetterich '93; (figure taken from) Gies '06]



# Color Superconductivity

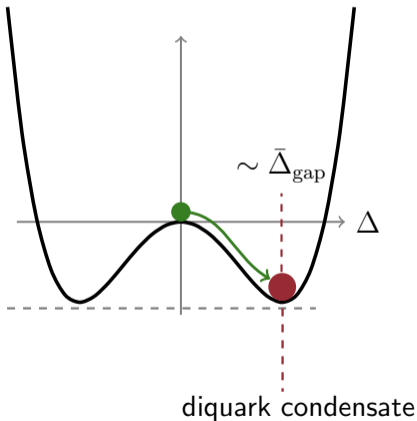


quark cooper pairs  
~ diquarks



ground state lowered

$U(1)_V$  symmetry breaking

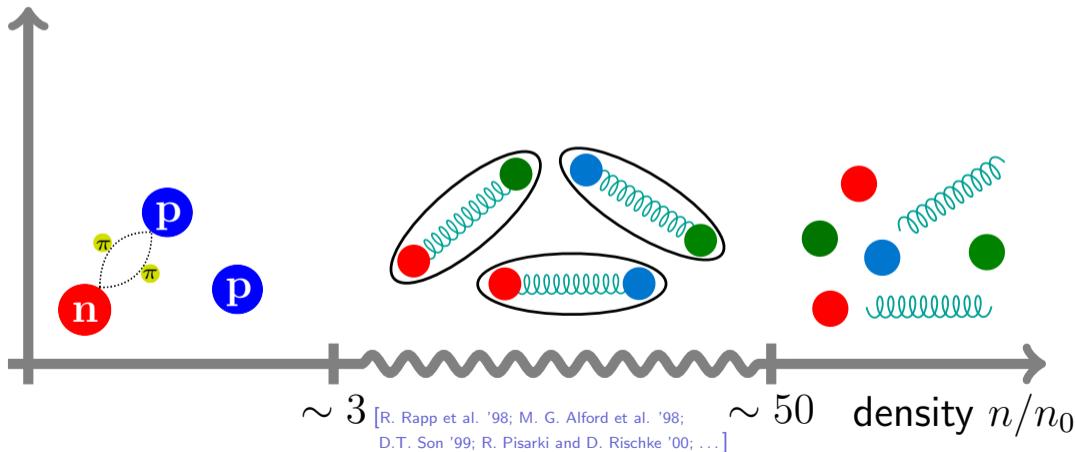


2-flavor diquark  
condensate is  
chirally symmetric



# Degrees of Freedom

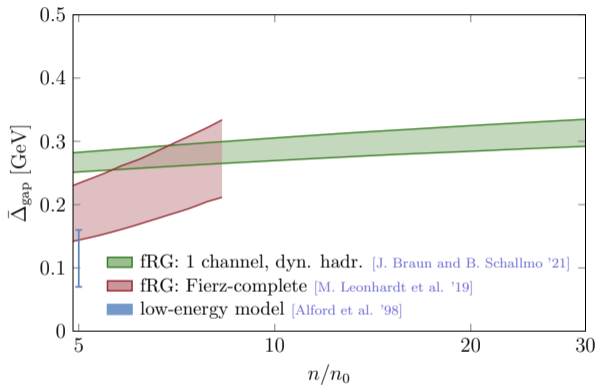
temperature  $T$



# Diquark Gap



no Taylor expansion  
about  $g_s = 0$



[M. Leonhardt et al. '19;  
J. Braun and B. Schallmo '21;  
J. Braun and B. Schallmo '22]

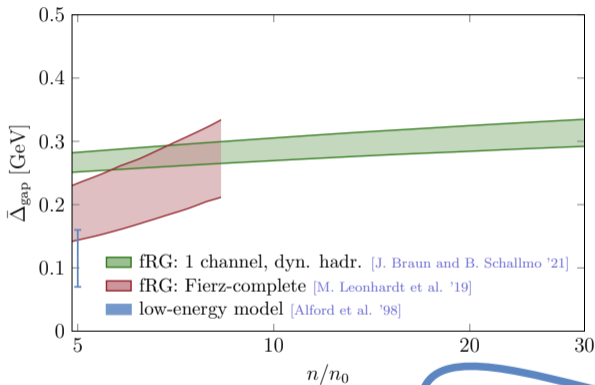
$$\bar{\Delta}_{\text{gap}} \sim \exp\left(-\frac{c}{g_s^4 n^{2/3}}\right), \quad c > 0$$

# Diquark Gap



[M. Leonhardt et al. '19;  
J. Braun and B. Schallmo '21;  
J. Braun and B. Schallmo '22]

no Taylor expansion  
about  $g_s = 0$



very high densities

$$\bar{\Delta}_{\text{gap}} \sim \exp\left(-\frac{c}{g_s^4 n^{2/3}}\right), \quad c > 0$$

$$\bar{\Delta}_{\text{gap}} \sim \frac{1}{g_s^5} \exp\left(-\frac{c'}{g_s}\right), \quad c' > 0$$

# Expansion of the Pressure



[J. Braun, AG, and B. Schallmo '22]

expansion about vanishing gap

$$P = P_{\text{SB}} \left( \gamma_0(g_s) + \gamma_1(g_s) \left( \frac{|\bar{\Delta}_{\text{gap}}|}{\mu} \right)^2 + \dots \right)$$

Stefan-Boltzmann  
pressure

$$\gamma_0 = 1 - \frac{g_s^2}{2\pi^2}$$

[T. Toimela '85]

$$\gamma_1 = 2$$

[K. Rajagopal and F. Wilczek '01]

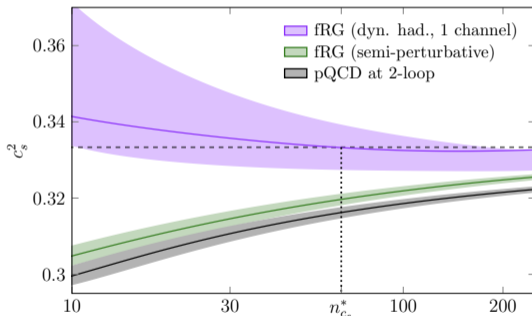
$\bar{\Delta}_{\text{gap}}(g_s)$

# Speed of Sound



[J. Braun, AG, and B. Schallmo '22; J. Braun, AG, and B. Schallmo (in prep.); J. Braun and AG (in prep.)]

$$c_s^2 = \frac{\partial P}{\partial \epsilon} = \frac{1}{\mu} \left( \frac{\partial P}{\partial \mu} \right) / \left( \frac{\partial^2 P}{\partial \mu \partial \mu} \right)$$



“crossing” density

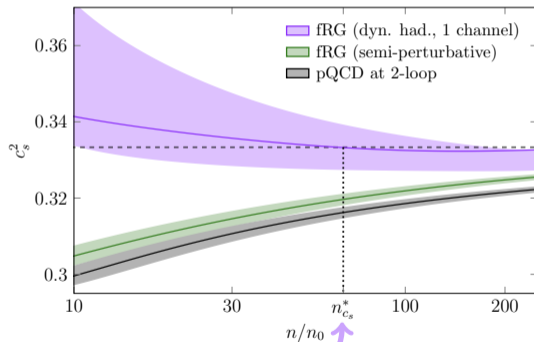
# Speed of Sound



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$$c_s^2 = \frac{\partial P}{\partial \epsilon} = \frac{1}{\mu} \left( \frac{\partial P}{\partial \mu} \right) / \left( \frac{\partial^2 P}{\partial \mu \partial \mu} \right)$$

- $n \gg n_{c_s}^*$ : gap contributions can be neglected
- $n > n_{c_s}^*$ : gap contributions are sizable
- $n \gtrsim n_{c_s}^*$ : gap contributions become significant



“crossing” density

# Speed of Sound



[J. Braun, AG, and B. Schallmo '22; J. Braun, AG, and B. Schallmo (in prep.); J. Braun and AG (in prep.)]

## scaling law

$$c_s^2 \approx \frac{1}{3} + c_0 (1 + \sigma) n^{-\frac{2(1-\sigma)}{3}} - \frac{c_1}{\ln(c_2 n^{1/3})}$$

free quark gas

gap contribution

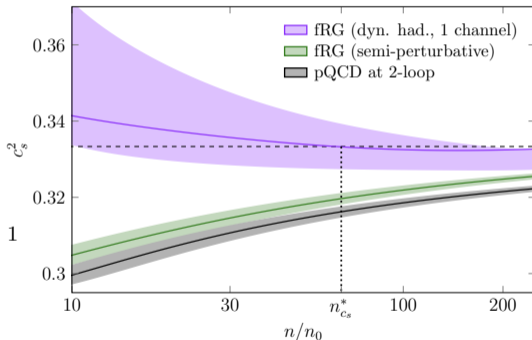
( $\sigma$  determined by scaling of gap)

$c_0, c_1 > 0$

$|\sigma| < 1$

$c_2 n^{1/3} > 1$

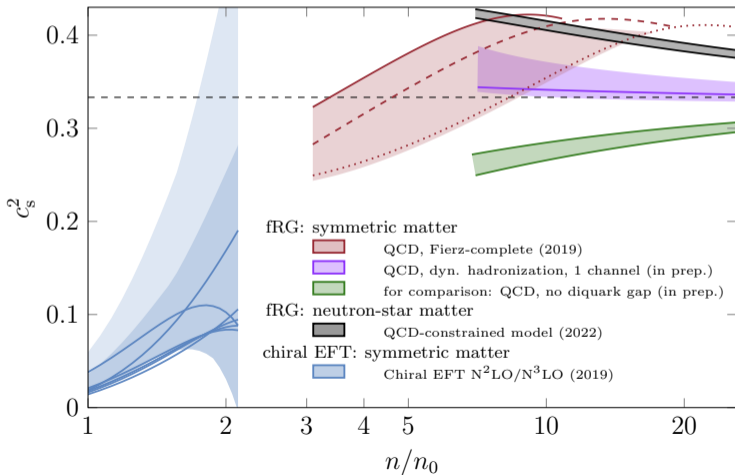
perturbative contribution



# Speed of Sound



[Leonhardt et al. '19; J. Braun and B. Schallmo '22; J. Braun, AG, and B. Schallmo '22; J. Braun, AG, and B. Schallmo (in prep.)]

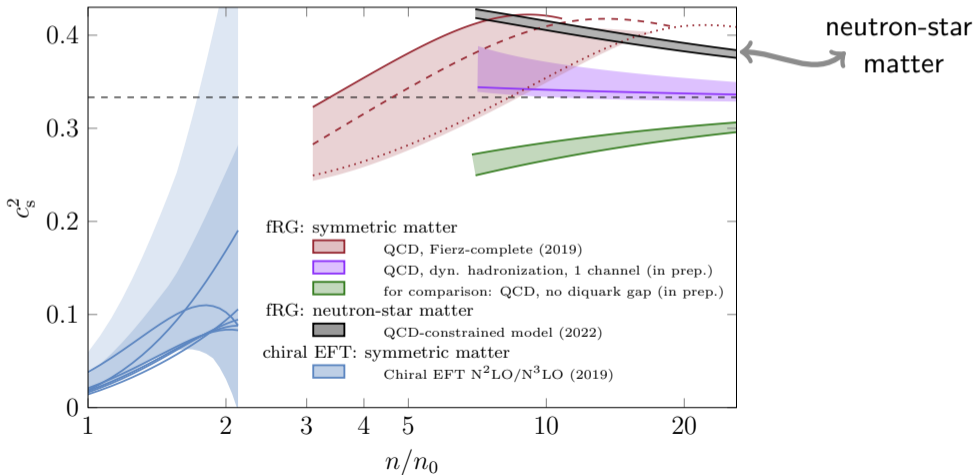




# Speed of Sound



[Leonhardt et al. '19; J. Braun and B. Schallmo '22; J. Braun, AG, and B. Schallmo '22; J. Braun, AG, and B. Schallmo (in prep.)]



# Take Home Message

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A color-superconducting **gap** suggests a **maximum** in the speed of sound at supranuclear densities.

At even higher densities, the speed of sound again **crosses** the conformal limit and approaches it from **below**.

# Thank you for your attention!

