



Prof. Dr. Ing. Peter Pelz
Wintersemester 2010/11
Technische Fluidsysteme
Vorlesung 4

$$\frac{Dm}{Dt} = \sigma \Leftrightarrow \frac{\partial}{\partial t} \int_V \rho dV + \oint_{S'} \rho \vec{u} \cdot \vec{n} dS' = 0.$$

$$\frac{D\rho}{Dt} + \rho \operatorname{div} \vec{u} = \sigma.$$

$$\frac{D(\rho \vec{u})}{Dt} = \vec{f}. \Leftrightarrow \frac{\partial}{\partial t} \int_V \rho \vec{u} dV + \oint_{S'} \rho \vec{u} \vec{u} \cdot \vec{n} dS' =$$
$$= \underbrace{\oint_{S'} \vec{t} dS'}_{\text{Oberfläche}} + \underbrace{\int_V \rho \vec{k} dV}_{\text{Volumenkraft}}.$$

$$\rho \frac{D\vec{u}}{Dt} = \rho \vec{h} + \nabla \cdot \underline{\underline{T}} \quad \text{Impulsbilanz in}$$

differenzieller Form.

§

$$\equiv \rho \frac{Zahl \text{ der Teilchen}}{Stoßlänge}$$



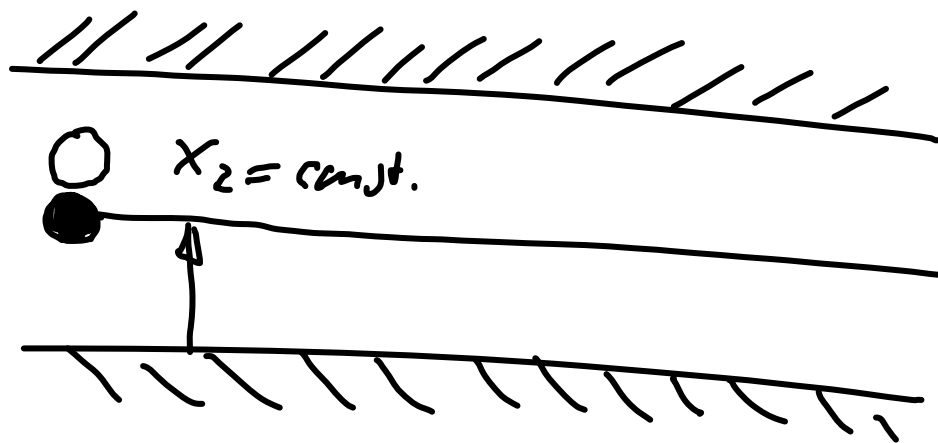
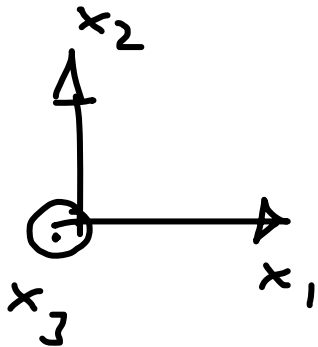
Prof. Dr. Ing. Peter Pelz
Wintersemester 2010/11
Technische Fluidsysteme
Vorlesung 4

Lösen der Transportgleichung für einfache
ebene Schichtströmungen (stationär)

$\mu_3 \equiv 0$

$\frac{\partial}{\partial t} \equiv 0$
 $\frac{D}{Dt} \neq 0$

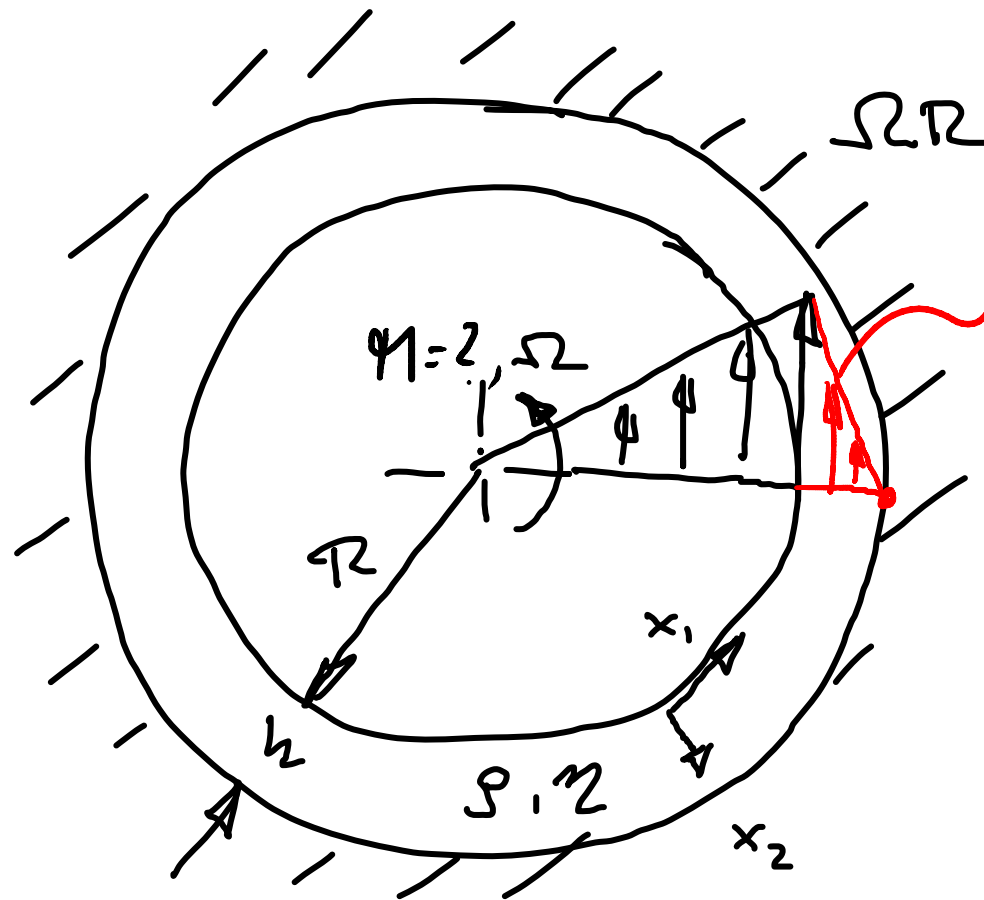
☺ Kinematik
 einfach.



Prof. Dr. Ing. Peter Pelz
 Wintersemester 2010/11
 Technische Fluidsysteme
 Vorlesung 4



Prof. Dr. Ing. Peter Pelz
Wintersemester 2010/11
Technische Fluidsysteme
Vorlesung 4

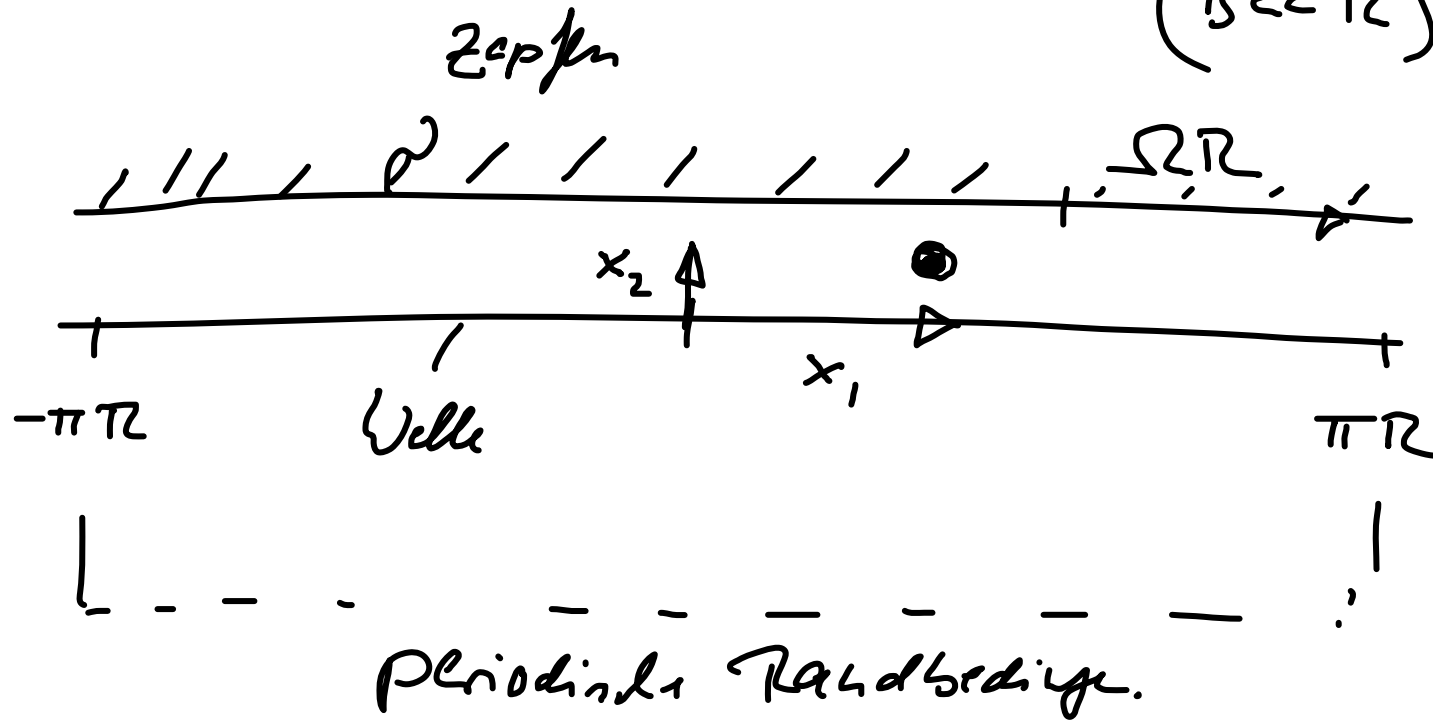


linearen Abfall
für $\frac{h}{R} \ll 1$.

η dynamische Viskosität der Flüssigkeit
im Spalt.

ebene Strömung, wenn die Zepfkreis $\beta \gg \tau$.

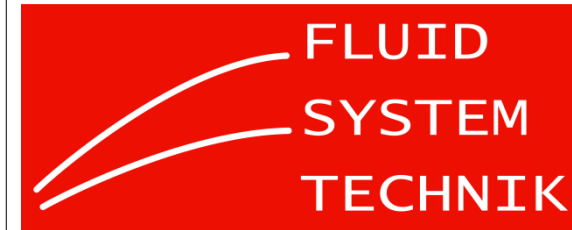
($\beta \ll \tau$).



$$\int \frac{D\vec{u}}{Dt} = \int \vec{h} + \nabla \cdot \vec{T}$$



TECHNISCHE
UNIVERSITÄT
DARMSTADT



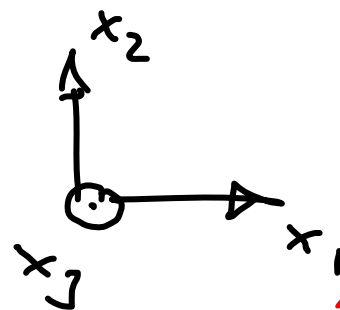
Prof. Dr. Ing. Peter Pelz
Wintersemester 2010/11
Technische Fluidsysteme
Vorlesung 4

allgemein Komp.

$$\vec{M} = \mu_1 \vec{e}_1 + \mu_2 \vec{e}_2 + \mu_3 \vec{e}_3 = \sum_i \mu_i \vec{e}_i \quad \mu_i$$

$$\vec{x} = x_1 \vec{e}_1 + x_2 \vec{e}_2 + x_3 \vec{e}_3 = \sum_i x_i \vec{e}_i \quad x_i$$

$$\begin{aligned} \tau &= \tau_{11} \vec{e}_1 \vec{e}_1 + \tau_{12} \vec{e}_1 \vec{e}_2 + \tau_{13} \vec{e}_1 \vec{e}_3 + \\ &\quad + \tau_{21} \vec{e}_2 \vec{e}_1 + \tau_{22} \vec{e}_2 \vec{e}_2 + \tau_{23} \vec{e}_2 \vec{e}_3 + \\ &\quad + \tau_{31} \vec{e}_3 \vec{e}_1 + \tau_{32} \vec{e}_3 \vec{e}_2 + \tau_{33} \vec{e}_3 \vec{e}_3 = \sum_{i,j} \tau_{ij} \vec{e}_i \vec{e}_j \quad \tau_{ij} \end{aligned}$$



Cartesisches Koordinatensystem.
 sind Ortsvektoren $\vec{e}_i \neq f_i(x_{ij})$

Zylinder $\vec{e}_r = f_r(\varphi)$



TECHNISCHE
UNIVERSITÄT
DARMSTADT

FLUID
SYSTEM
TECHNIK



Prof. Dr. Ing. Peter Pelz
Wintersemester 2010/11
Technische Fluidsysteme
Vorlesung 4

Skalarprodukt

$$\vec{x} \cdot \vec{y} = \sum_{i,j} x_i \vec{e}_i \cdot y_j \vec{e}_j = \sum_{i,j} x_i y_j \vec{e}_i \cdot \vec{e}_j$$

Recht: Immer dann, wenn Indizes doppelt vorkommen, ~~ist~~ über die Indize wird summiert!

$$\vec{x} \cdot \vec{y} = x_i y_i \vec{e}_i \cdot \vec{e}_i = x_i y_i \delta_{ii} = x_i y_i$$

$\vec{e}_1 \cdot \vec{e}_1 = 1$	$\vec{e}_2 \cdot \vec{e}_1 = 0$	$\hat{=} \hat{=} \hat{=} \delta_{ij}$
$\vec{e}_1 \cdot \vec{e}_2 = 0$	$\vec{e}_2 \cdot \vec{e}_2 = 1$	

$\hat{=}$ Einheitskenn. Kronecker Symbol.



Prof. Dr. Ing. Peter Pelz
Wintersemester 2010/11
Technische Fluidsysteme
Vorlesung 4

$$\rho \frac{D\vec{u}}{Dt} = \rho \vec{k} + \nabla \cdot \underline{\underline{\tau}}$$

Symbolische Schreibweise.

$$\rho \frac{D\mu_i}{Dt} = \rho k_i + \frac{\partial}{\partial x_j} \tau_{ij}$$

Indexnotation.

$i=1$

$$\rho \frac{D\mu_1}{Dt} = \rho k_1 + \frac{\partial}{\partial x_1} \tau_{11} + \frac{\partial}{\partial x_2} \tau_{12}$$

$i=2$

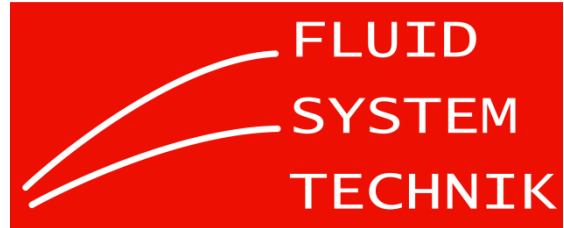
$$\rho \frac{D\mu_2}{Dt} = \rho k_2 + \frac{\partial}{\partial x_1} \tau_{21} + \frac{\partial}{\partial x_2} \tau_{22}$$

30.11.2010

$i=3$ ist für das oben Problem nicht relevant!

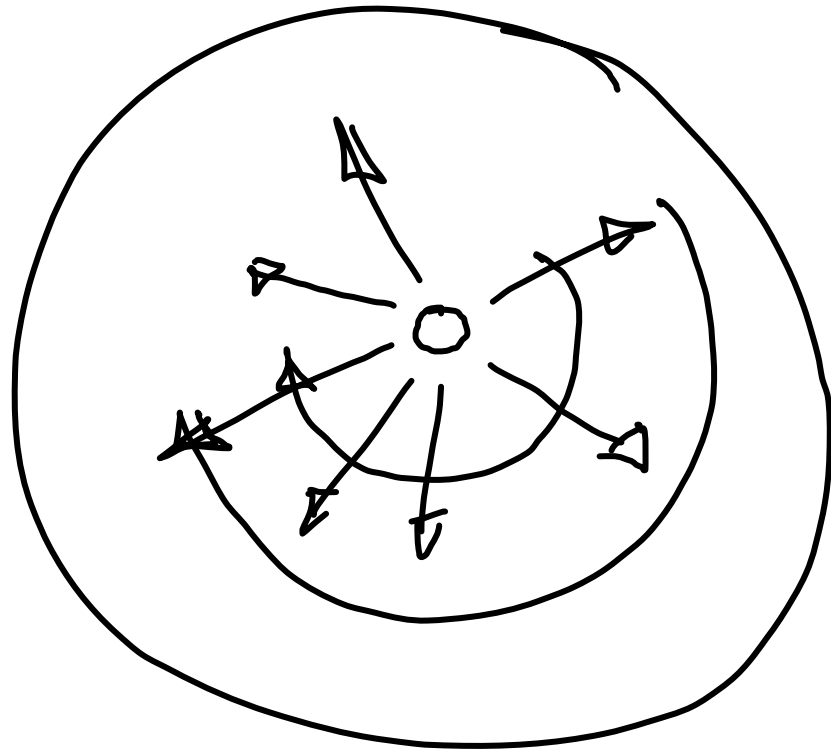


TECHNISCHE
UNIVERSITÄT
DARMSTADT



Prof. Dr. Ing. Peter Pelz
Wintersemester 2010/11
Technische Fluidsysteme
Vorlesung 4

Zur Volumenkraft $\rho \mathbf{k}$:



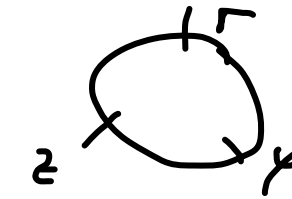
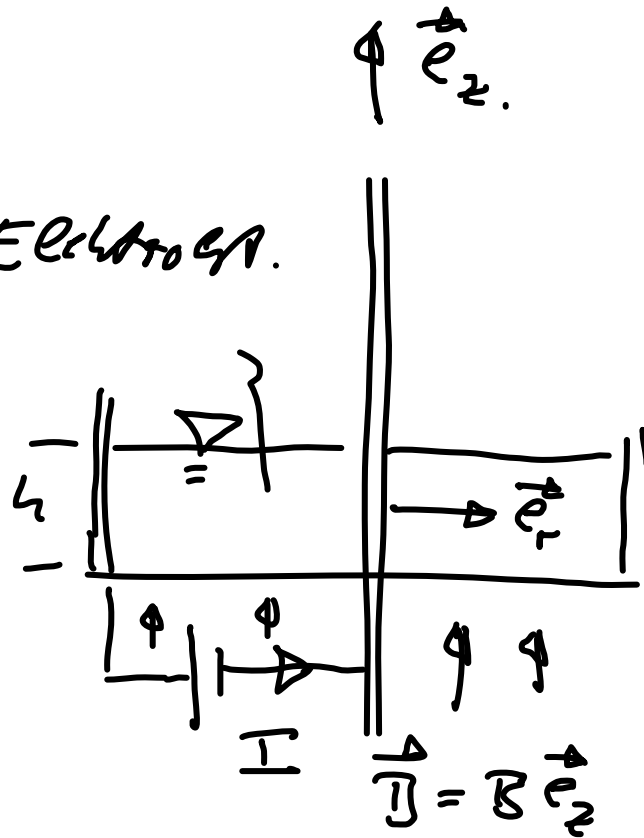
$$\vec{c} = \frac{I}{2\pi r h} \vec{e}_r$$

Gesetz von Ampère

$$\rho \mathbf{k} = \vec{c} \times \vec{B}$$

$$= \frac{I}{2\pi r h} \vec{e}_r \times B \vec{e}_z = -\frac{I B}{2\pi r h} \vec{e}_\varphi$$

Elektronen



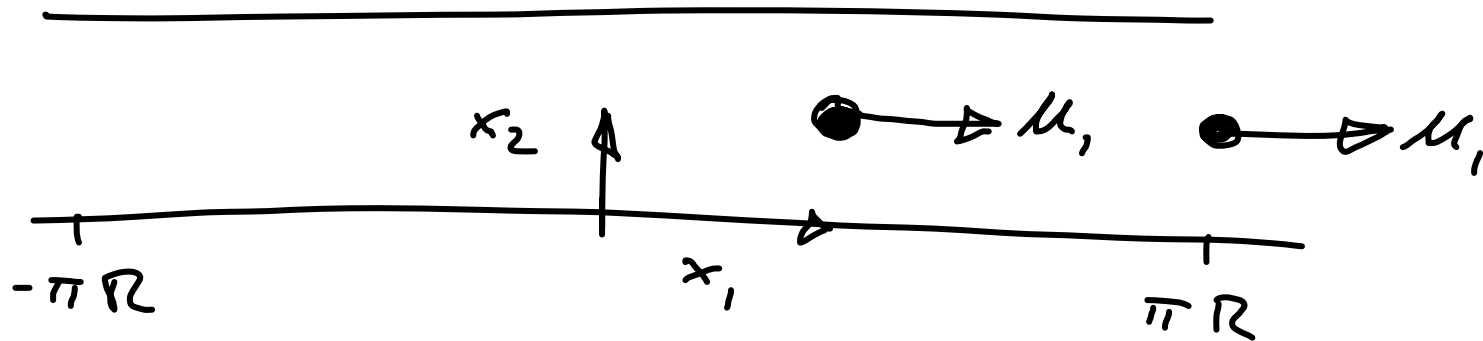
TECHNISCHE
UNIVERSITÄT
DARMSTADT

FLUID
SYSTEM
TECHNIK



Prof. Dr. Ing. Peter Pelz
Wintersemester 2010/11
Technische Fluidsysteme
Vorlesung 4

$\rho k_i \equiv \sigma$ keine Volumenkr. ρ



$u_2 \equiv 0$, da Schrittström.

$$\frac{\partial}{\partial x_1} \equiv 0$$

$$\frac{\partial}{\partial t} \equiv 0 \quad \text{stationär, Ström.}$$

kinematische
Vergleich bei
dem Schrittström.

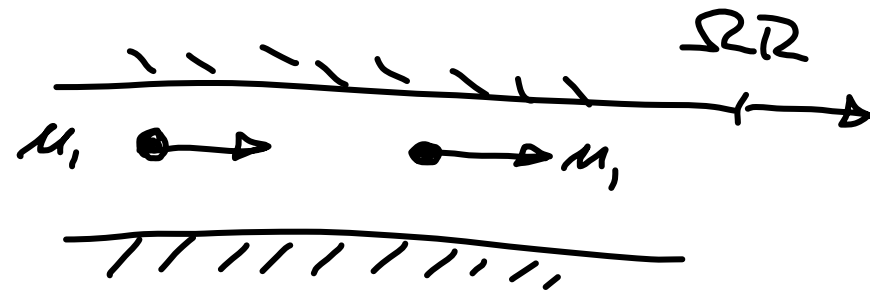


Prof. Dr. Ing. Peter Pelz
Wintersemester 2010/11
Technische Fluidsysteme
Vorlesung 4



Prof. Dr. Ing. Peter Pelz
Wintersemester 2010/11
Technische Fluidsysteme
Vorlesung 4

$$\underbrace{\rho \frac{D\mu_1}{Dt}}_{\equiv \sigma} = \frac{\partial}{\partial x_1} \tau_{11} + \frac{\partial}{\partial x_2} \tau_{21}$$



$$\sigma = \frac{\partial}{\partial x_1} \tau_{11} + \frac{\partial}{\partial x_2} \tau_{21}$$

$$\tau_{ij} = \underbrace{-p \delta_{ij}}_{\text{Hydrost. Druck}} + \underbrace{\tau_{ij}}_{\text{Reibspannung}}$$

Abspalten des hydro-
statischen Drucks

Hydrost. Druck Reibspannung.

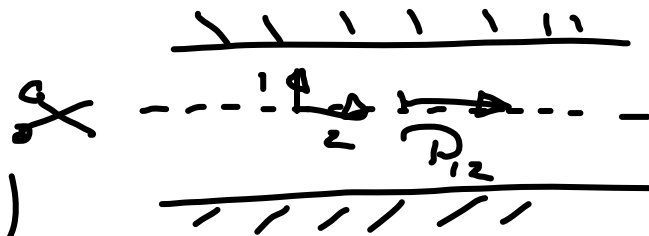


$$\underline{\underline{T}} = -p \underline{\underline{T}} + \underline{\underline{P}}$$

$$\tau_{ij} = -p \delta_{ij} + P_{ij}$$

Materialgesetz für den Reibspannungszustand.

$$P_{12} = \int \mu (\text{Deformationsgeschw.})$$



$$P_{12} = 2\eta e_{12}$$

$$e_{12} = \frac{1}{2} \left(\frac{\partial u_1}{\partial x_2} + \frac{\partial u_2}{\partial x_1} \right)$$

$$P_{ij} = 2\eta e_{ij}$$

$$e_{ij} = \frac{1}{2} \left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right)$$

30.11.2010 Newtonsche Fluide.



Prof. Dr. Ing. Peter Pelz
Wintersemester 2010/11
Technische Fluidsysteme
Vorlesung 4



Prof. Dr. Ing. Peter Pelz
Wintersemester 2010/11
Technische Fluidsysteme
Vorlesung 4

$$0 = -\frac{\partial p}{\partial x_1} + \frac{\partial \tau_{11}}{\partial x_1} + \frac{\partial \tau_{12}}{\partial x_2} + \frac{\partial \tau_{22}}{\partial x_2}$$

$$= -\frac{\partial p}{\partial x_1} + \underbrace{2\eta \frac{\partial^2 u_1}{\partial x_1^2}}_{\equiv 0} + \eta \frac{\partial^2 u_1}{\partial x_2^2} + \underbrace{2\eta \frac{\partial^2 u_2}{\partial x_2^2}}_{\equiv 0}$$

$$\frac{\partial p}{\partial x_1} = \eta \frac{\partial^2 u_1}{\partial x_2^2}$$

⊙ lineare Gleichg.
liefert zu löse.

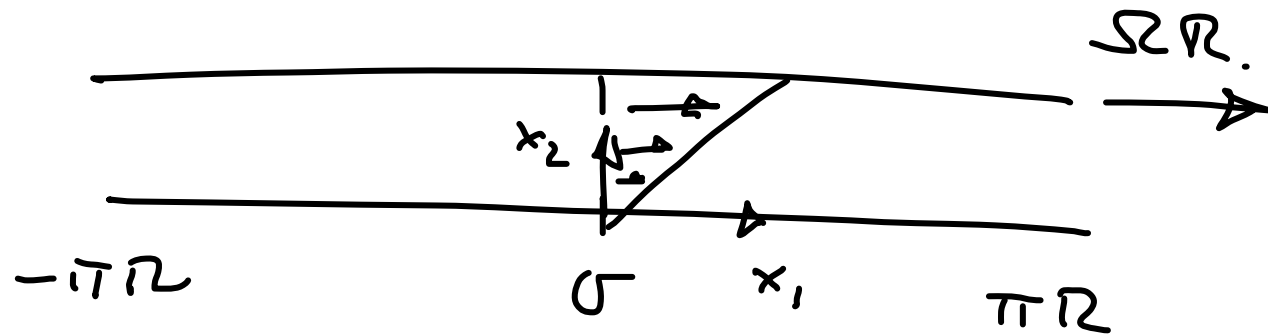
1. Lösung für die lin. f. Schleppström.



Prof. Dr. Ing. Peter Pelz
Wintersemester 2010/11
Technische Fluidsysteme
Vorlesung 4

Randbeding $\mu_1(0) = 0 \rightsquigarrow C_2 = 0$

$$\mu_1(h) = \Omega R \rightsquigarrow C_1 = \Omega \frac{R}{L}$$



$$\frac{\partial \mu}{\partial x_2} = \rho$$

$$\rho = \frac{\partial^2 \mu_1}{\partial x_2^2}$$

$$C_1 = \frac{\partial \mu_1}{\partial x_2}$$

$$C_1 x_2 + C_2 = \mu_1$$

$$\underline{\underline{\mu_1(x_2) = \Omega R \frac{x_2^2}{L}}}$$

Reibung: Reibmoment pro Tiefen-einheits

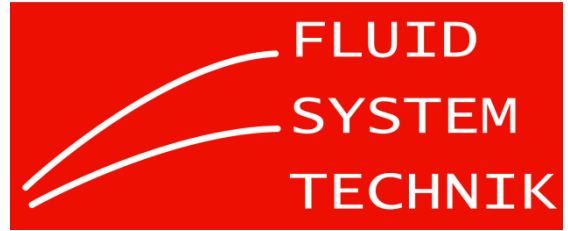
$$M_z = \int_{\text{Zylinder}} \vec{x} \times \vec{t} \, dA \Big| \cdot \vec{e}_z$$

$$= \int_0^{2\pi} R \vec{e}_r \times \tau \cdot \vec{e}_r \, R \, d\varphi$$

$$= 2\pi R^2 \tau_{r\varphi} = 2\pi R^2 \tau_{12}$$

$$= 2\pi R^2 \eta \frac{\partial u_1}{\partial x_2} = \frac{2\pi R^3 \eta \Omega}{h}$$

$$P_v = \vec{T} \cdot \vec{\Omega} = \frac{2\pi R^3 \eta \Omega^2}{h} \quad \text{Verl. Leistung}$$



Prof. Dr. Ing. Peter Pelz
Wintersemester 2010/11
Technische Fluidsysteme
Vorlesung 4



Prof. Dr. Ing. Peter Pelz
Wintersemester 2010/11
Technische Fluidsysteme
Vorlesung 4

Cauchy-Gl. + Newtonsche
Ableitung = Navier-Stokes-Gl.

$$\rho \frac{D u_i}{D t} = \rho k_i + \frac{\partial \tau_{ij}}{\partial x_j}$$

$$\tau_{ij} = -p \delta_{ij} + 2\eta e_{ij}$$

$$e_{ij} = \frac{1}{2} \left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right)$$

$$\rho \frac{D u_i}{D t} = \rho k_i - \frac{\partial p}{\partial x_i} + \eta \frac{\partial^2 u_i}{\partial x_j^2}$$

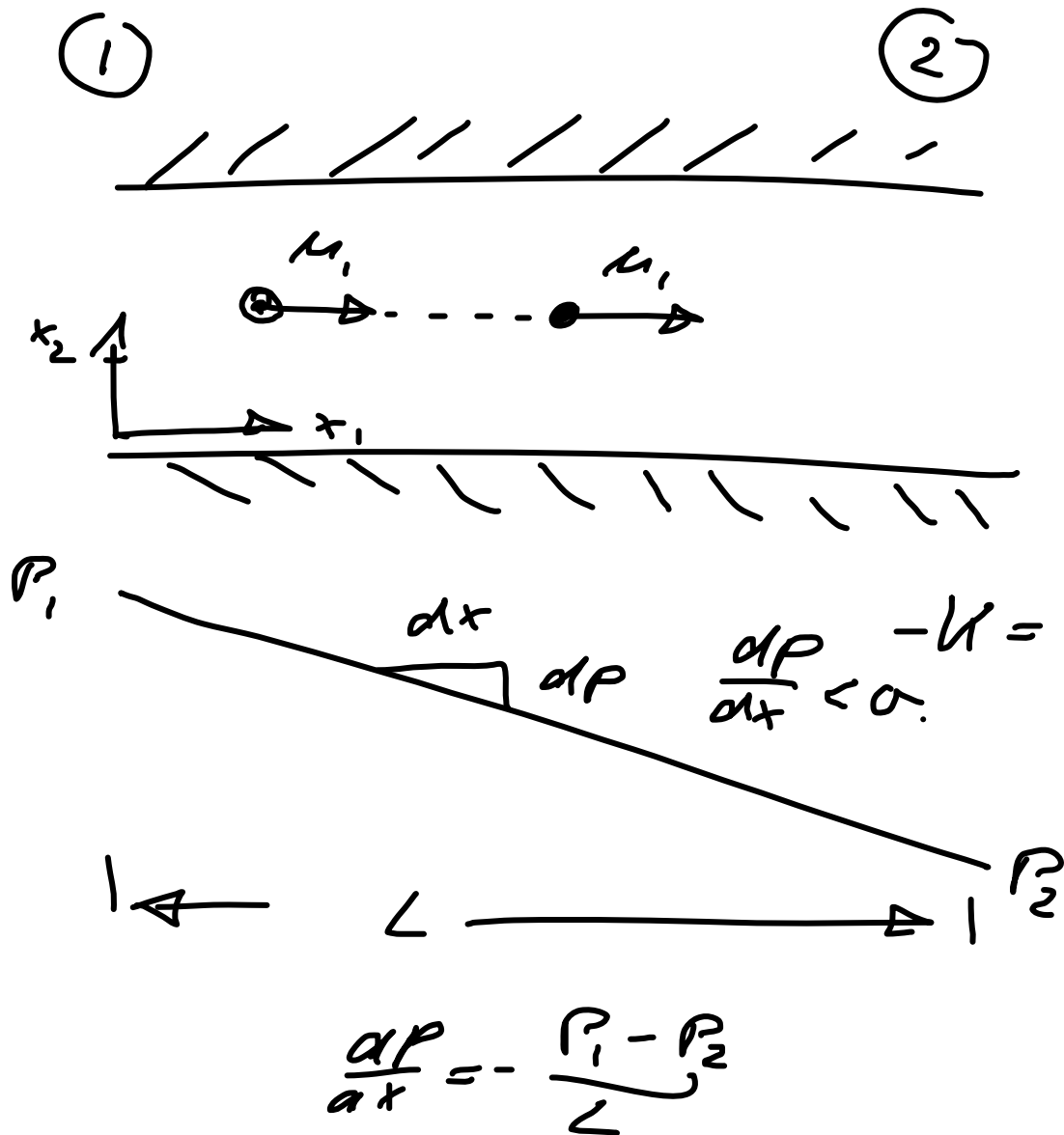
$\eta = \text{const.}$

Für $\eta = \text{const.}$

$$\rho \frac{D u_i}{D t} = \rho k_i - \frac{\partial p}{\partial x_i} + \eta \frac{\partial^2 u_i}{\partial x_j^2}$$



Prof. Dr. Ing. Peter Pelz
Wintersemester 2010/11
Technische Fluidsysteme
Vorlesung 4



2D $\mu_3 \equiv 0$

$$\rho \frac{D u_1}{D t} \equiv 0$$

$$\rho \bar{h}_1 \equiv 0$$

$$\frac{\partial P}{\partial x_1} = \gamma \frac{\partial^2 u_1}{\partial x_2^2}$$

1. I.d.V.

$$C_1 - \frac{\gamma}{2} x_2 = \frac{\partial M_1}{\partial x_2}$$

$$\mu := - \frac{dP}{dx_1}$$

2. I.d.V.

$$C_1 x_2 - \frac{\gamma}{2} x_2^2 + C_2 = M_1$$

$$\text{R.B. } M_1(0) = 0$$

$$C_2 = 0$$

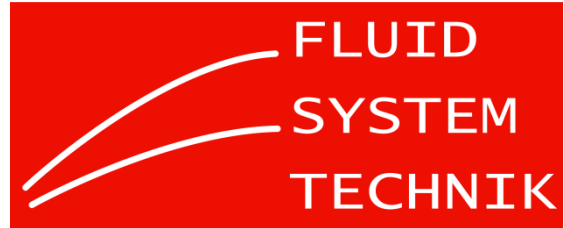
$$M_1(h) = 0$$

$$C_1 = \frac{\gamma}{2} h$$

$$M_1 = \frac{\gamma h^2}{2} \left(\frac{x_2}{h} - \left(\frac{x_2}{h} \right)^2 \right) = \frac{\gamma h^2}{2} \frac{x_2}{h} \left(1 - \frac{x_2}{h} \right)$$



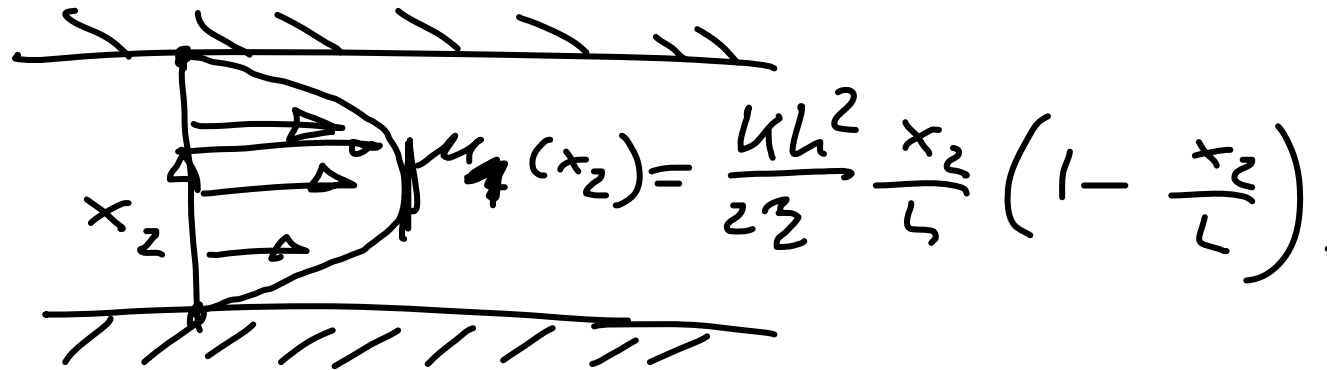
TECHNISCHE
UNIVERSITÄT
DARMSTADT



Prof. Dr. Ing. Peter Pelz
Wintersemester 2010/11
Technische Fluidsysteme
Vorlesung 4



Prof. Dr. Ing. Peter Pelz
Wintersemester 2010/11
Technische Fluidsysteme
Vorlesung 4



$$u_{\max} = u_1\left(\frac{L}{2}\right) = \frac{4L^2}{8\zeta}$$

$$\frac{\partial p}{\partial x_1} = \zeta \frac{\partial^2 u_1}{\partial x_2^2} \quad \text{linear profile.}$$

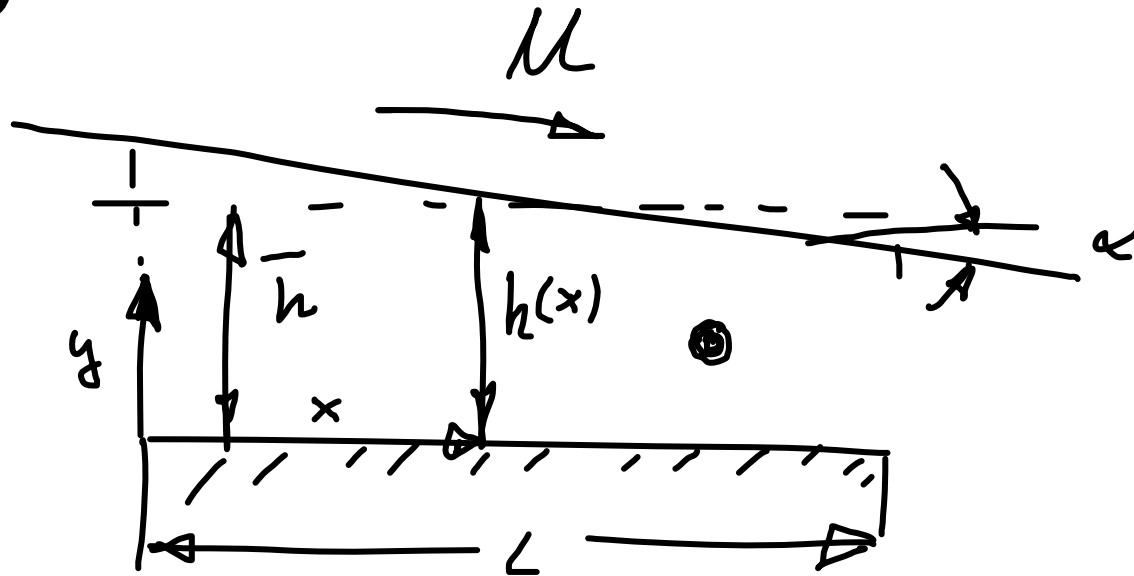
no superposition in motion.



Hydrodynamische Schminz

\bar{h} ist bei einem
Zerfließen $\mu =$
viskoses Spiel.

Keillay.



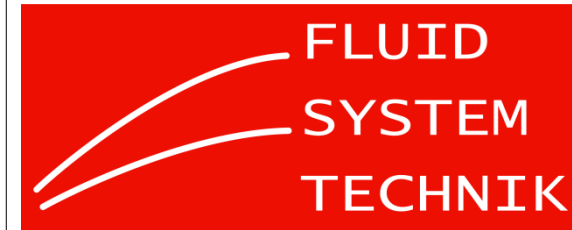
$$\left\{ \frac{\partial p}{\partial x} = \mu \frac{\partial^2 u}{\partial y^2} \right.$$

gilt für $\alpha \text{Re} \ll 1$.

Reynoldszahl $\text{Re} = \frac{\mu \bar{h} \delta}{\eta}$



TECHNISCHE
UNIVERSITÄT
DARMSTADT



Prof. Dr. Ing. Peter Pelz
Wintersemester 2010/11
Technische Fluidsysteme
Vorlesung 4

Wichtig

$$\alpha * Re \ll 1$$

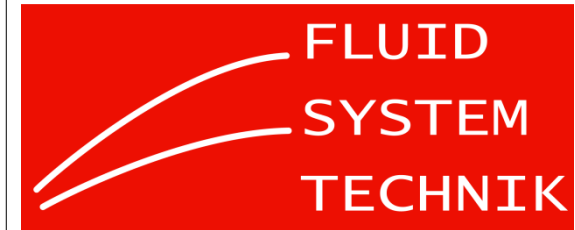
Re kann groß sein!

turbulente Strömung nicht möglich!

$$\mu(x) = \mu \frac{y}{h(x)} - \frac{\partial p}{\partial x} \frac{h^2(x)}{2\eta} \left(1 - \frac{y}{h(x)}\right) \frac{y}{h(x)}$$



TECHNISCHE
UNIVERSITÄT
DARMSTADT



Prof. Dr. Ing. Peter Pelz
Wintersemester 2010/11
Technische Fluidsysteme
Vorlesung 4