

PC III

1.7.13

### 8.4 Wärme und Arbeit

$$U = \langle E \rangle = \sum_j E_j(N, V) \cdot p_j(N, V, T)$$

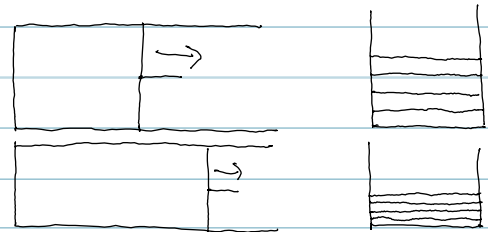
$$\begin{aligned} du &= \sum_j d[E_j(N, V) p_j(N, V, T)] \\ &= \sum_j E_j dp_j + \sum_j p_j dE_j \\ &= \sum_j E_j dp_j + \sum_j p_j \left( \frac{\partial E_j}{\partial V} \right)_N dV \end{aligned}$$

makroskopische Beschreibung von  $du$ :

$$du = dq_{\text{rev}} + dw_{\text{rev}}$$

$$dw_{\text{rev}} = -p dV$$

$$dw_{\text{rev}} = \sum_j p_j \left( \frac{\partial E_j}{\partial V} \right)_N dV$$



reversibel ausgetauschte Arbeit resultiert aus einer infinitesimalen Änderung der Energien des Systems

$$p = - \sum_j p_j \left( \frac{\partial E_j}{\partial V} \right)_N$$

$$p = - \left\langle \left( \frac{\partial E}{\partial V} \right)_N \right\rangle$$

$$dq_{\text{rev}} = \sum_j E_j dp_j$$

reversibel ausgetauschte Wärme resultiert aus einer Änderung der Wahrscheinlichkeitsverteilung der Zustände des Systems

### 8.5 Druck

$$\langle p \rangle = \sum_j p_j(N, V, T) \overset{\text{Druck}}{p_j(N, V)} = \sum_j p_j \left( -\frac{\partial E}{\partial V} \right)_N = \sum_j \left( \frac{\partial E_j}{\partial V} \right)_N \frac{e^{-E_j/kT}}{Q}$$

Druck in Abhängigkeit von der Zustandssumme!

$$\left( \frac{\partial Q}{\partial V} \right)_{N, T} = -\frac{1}{kT} \sum_j \left( \frac{\partial E_j}{\partial V} \right)_N e^{-E_j/kT}$$

$$\frac{1}{Q} \left( \frac{\partial Q}{\partial V} \right)_{N, T} = \left( \frac{\partial \ln Q}{\partial V} \right)_{N, T} = -\frac{1}{kT} \sum_j \left( \frac{\partial E_j}{\partial V} \right)_N \frac{e^{-E_j/kT}}{Q}$$

$$\boxed{\langle p \rangle = kT \left( \frac{\partial \ln Q}{\partial V} \right)_{N, T}} \quad \text{Druck in Abhängigkeit der Zustandssumme}$$

z.B. 1-atomiges ideales Gas

$$Q = \frac{q^N}{N!} \quad \ln Q = N \ln q - \ln N!$$

$$q = \left( \frac{2\pi m k T}{h^2} \right)^{3/2} V g_{\text{rel}}$$

$$\ln Q = N \cdot \frac{3}{2} \ln \left( \frac{2\pi m k T}{h^2} \right) + N \ln V + N \ln g_{\text{rel}} - \ln N!$$

$$\langle p \rangle = kT \frac{1}{V} N = \frac{NkT}{V} = \frac{nRT}{V} = p \quad \underline{\text{ideales Gasgesetz}}$$

Warum ist  $\beta = \frac{1}{kT}$ ?

$$S_{\text{ensemble}} = k \ln W = k \ln \frac{A!}{\prod_j a_j} = k \ln A! - k \sum_j \ln a_j!$$

$$\text{Stirling: } \ln x! \approx x \ln x - x$$

$$S_{\text{ensemble}} = k \ln A - kA - k \sum_j a_j \ln a_j + k \sum_j a_j \quad a_j = p_j A$$

$$= kA \ln A - k \sum_j p_j A \ln p_j A$$

$$= kA \ln A - k \sum_j p_j A \ln p_j - k \sum_j p_j A \ln A$$

$$S_{\text{ensemble}} = -k \sum_j p_j A \ln p_j$$

$$S_{\text{System}} = \frac{S_{\text{Ensemble}}}{A} = -k \sum_j p_j \ln p_j \quad \rightarrow 0$$

$$dS = -k \sum_j d(p_j \ln p_j) = -k \sum_j (\ln p_j dp_j + dp_j) \quad p_j = \frac{e^{-E_j \beta}}{Q}$$

$$= -k \sum_j \ln p_j dp_j = -k \sum_j (-E_j \beta - \ln Q) dp_j$$

$$= -k \sum_j -E_j \beta dp_j + k \sum_j \ln Q dp_j$$

$$dS = k\beta \sum_j E_j dp_j$$

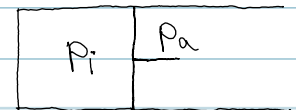
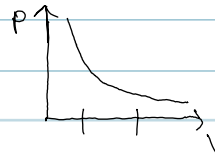
$$dS = k\beta \cdot dq_{\text{rev}} = \frac{dq_{\text{rev}}}{T}$$

$$\boxed{\beta = \frac{1}{kT}}$$

### isotherme Expansion eines Gases (1-atomig, ideal)

Volumenarbeit

$$W_{\text{rev}} = - \int_{V_A}^{V_E} p dV = - \int_{V_A}^{V_E} \frac{R n T}{V} dV$$



$$= -R n T \int_{V_A}^{V_E} \frac{dV}{V} = -R n T \ln \frac{V_E}{V_A} = -N k T \ln \frac{V_E}{V_A} \quad R = k N_A$$

$$n = \frac{N}{N_A}$$

statistisch:

$$W_{\text{rev}} = - \int_{V_A}^{V_E} p dV = - \int_{V_A}^{V_E} k T \left( \frac{\partial \ln Q}{\partial V} \right)_{N, T} dV = - k T \int_{V_A}^{V_E} \left( \frac{\partial \ln Q}{\partial V} \right)_{N, T} dV$$

$$Q = \frac{q^N}{N!} \quad \ln Q = N \ln q - \ln N!$$

$$q = \left( \frac{2\pi m k T}{h^2} \right)^{\frac{3}{2}} V g_{\text{el}} \quad \left( \frac{\partial \ln Q}{\partial V} \right) = \frac{1}{V}$$

$$W_{\text{rev}} = -k T \int_{V_A}^{V_E} \frac{1}{V} dV = -N k T \ln \frac{V_E}{V_A}$$

## Entropieänderung

isotherm:  $du = 0$

$$dq_{\text{rev}} = -dw_{\text{rev}}$$

$$ds = \frac{dq_{\text{rev}}}{T} = -\frac{dw_{\text{rev}}}{T}$$

$$\Delta S = nR \ln \frac{V_E}{V_A}$$

statistisch:

$$S = kT \left( \frac{\partial \ln Q}{\partial T} \right)_{N,V} + k \ln Q$$

$$q = \left( \frac{2\pi mkT}{h^2} \right)^{\frac{3}{2}} V_{\text{gel}}$$

$$S_A = kT \frac{3}{2} N \frac{1}{T} + Nk \ln(\alpha V_A)$$

$$S_E = kT \frac{3}{2} N \frac{1}{T} + Nk \ln(\alpha V_E)$$

$$\Delta S = S_E - S_A = Nk \ln \left( \frac{V_E}{V_A} \right)$$

Bsp.: Enthalpie in Abhängigkeit von der Zustandssumme

$$H = U + pV = kT^2 \left( \frac{\partial \ln Q}{\partial T} \right)_{N,V} + kT \left( \frac{\partial \ln Q}{\partial V} \right)_{N,T} \cdot V$$
$$= U + nRT$$

z.B.: 1-atomiges ideales Gas

$$H = kT^2 \frac{3}{2} N \frac{1}{T} + kT \frac{N}{V} V = \frac{3}{2} NkT + NkT = \frac{5}{2} NkT$$

$$1 \text{ Mol: } H_m = \frac{5}{2} RT$$